Neutral Higgs Bosons, Searches for

CONTENTS:

MASS LIMITS FOR NEUTRAL HIGGS BOSONS IN SUPERSYMMETRIC MODELS

The minimal supersymmetric model has two complex doublets of Higgs bosons. The resulting physical states are two scalars [H_1^0 and H_2^0, where we define $m_{H_1^0} < m_{H_2^0}$], a pseudoscalar (A^0), and a charged Higgs pair ($H^\pm$). H_1^0 and H_2^0 are also called h and H in the literature. There are two free parameters in the Higgs sector which can be chosen to be m_{A^0} and $\tan\beta = v_2/v_1$, the ratio of vacuum expectation values of the two Higgs doublets. Tree-level Higgs masses are constrained by the model to be $m_{H_1^0} \leq m_Z$, $m_{H_2^0} \geq m_Z$, $m_{A^0} \geq m_{H_1^0}$, and $m_{H^\pm} \geq m_W$. However, as described in the review on "Status of Higgs Boson Physics" in this Volume these relations are violated by radiative corrections.

The observed signal at about 125 GeV, see section "H", can be interpreted as one of the neutral Higgs bosons of supersymmetric models. Unless otherwise noted, we identify the lighter scalar H_1^0 with the Higgs discovered at 125 GeV at the LHC (AAD 12AI, CHATRCHYAN 12N).

Unless otherwise noted, the experiments in e^+e^- collisions search for the processes $e^+e^- \rightarrow H_1^0 Z^0$ in the channels used for the Standard Model Higgs searches and $e^+e^- \rightarrow H_1^0 A^0$ in the final states $b\overline{b}b\overline{b}$ and $b\overline{b}\tau^+\tau^-$. Unless otherwise stated, the following results assume no invisible H_1^0 or A^0 decays. Unless otherwise noted, the results are given in the m_h^{max} scenario, CARENA 13.

In $p\overline{p}$ and pp collisions the experiments search for a variety of processes, as explicitly specified for each entry. Limits on the A^0 mass arise from these direct searches, as well as from the relations valid in the minimal supersymmetric model between m_{A^0} and $m_{H_1^0}$. As discussed in the review on "Status of Higgs Boson Physics" in this Volume, these relations depend, via potentially large radiative corrections, on the mass of the t quark and on the supersymmetric parameters, in particular those of the stop sector. These indirect limits are weaker for larger t and $\widetilde{t}$ masses.

To include the radiative corrections to the Higgs masses, unless otherwise stated, the listed papers use theoretical predictions incorporating two-loop corrections and beyond (SLAVICH 21), and the results are given for the M_h^{125} benchmark scenario, see BAGNASCHI 19.

Mass Limits for heavy neutral Higgs bosons (H_2^0, A^0) in the MSSM

The limits rely on $pp \to H_2^0/A^0 \to \tau^+\tau^-$ and assume that H_2^0 and A^0 are (sufficiently) mass degenerate. The limits depend on tanβ.

VALUE (GeV)	CL%	DOCUMENT ID		TECN	COMMENT
> 835	95	[1] TUMASYAN	23S	CMS	tanβ = 10 GeV
>1240	95	[1] TUMASYAN	23S	CMS	tanβ = 20 GeV
>1605	95	[1] TUMASYAN	23S	CMS	tanβ = 30 GeV
>1820	95	[1] TUMASYAN	23S	CMS	tanβ = 40 GeV
>1950	95	[1] TUMASYAN	23S	CMS	tanβ = 50 GeV
>2062	95	[1] TUMASYAN	23S	CMS	tanβ = 60 GeV
>1121	95	[2] AAD	20AA	ATLS	tanβ = 10 GeV
>1475	95	[2] AAD	20AA	ATLS	tanβ = 20 GeV
>1677	95	[2] AAD	20AA	ATLS	tanβ = 30 GeV
>1826	95	[2] AAD	20AA	ATLS	tanβ = 40 GeV
>1937	95	[2] AAD	20AA	ATLS	tanβ = 50 GeV
>2033	95	[2] AAD	20AA	ATLS	tanβ = 60 GeV
• • • We do not use the following data for averages, fits, limits, etc. • • •					
		[3] AAD	24AP	ATLS	H_2^0, $A^0 \to t\bar{t}$
		[4] AAD	24CB	ATLS	H properties
		[5] AAD	24H	ATLS	$H_2^0 \to HH$
		[6] AAD	20	ATLS	H properties
		[7] AAD	20C	ATLS	$H_2^0 \to HH$
		[8] AAD	20L	ATLS	$H_2^0 \to b\bar{b}$
		[9] SIRUNYAN	20AC	CMS	$A^0 \to ZH$
		[10] SIRUNYAN	20AF	CMS	$H_2^0/A^0 \to t\bar{t}$
		[11] SIRUNYAN	20Y	CMS	$H_2^0 \to W^+W^-$
		[12] SIRUNYAN	19CR	CMS	$H_2^0/A^0 \to \mu^+\mu^-$
> 377	95	[13] AABOUD	18G	ATLS	tanβ = 10 GeV
> 863	95	[13] AABOUD	18G	ATLS	tanβ = 20 GeV
>1157	95	[13] AABOUD	18G	ATLS	tanβ = 30 GeV
>1328	95	[13] AABOUD	18G	ATLS	tanβ = 40 GeV
>1483	95	[13] AABOUD	18G	ATLS	tanβ = 50 GeV
>1613	95	[13] AABOUD	18G	ATLS	tanβ = 60 GeV
		[14] SIRUNYAN	18A	CMS	$H_2^0 \to H^0 H^0$
		[15] SIRUNYAN	18BP	CMS	$pp \to H_2^0/A^0 + b + X$, $H_2^0/A^0 \to b\bar{b}$
> 389	95	[16] SIRUNYAN	18CX	CMS	tanβ = 10 GeV
> 832	95	[16] SIRUNYAN	18CX	CMS	tanβ = 20 GeV
>1148	95	[16] SIRUNYAN	18CX	CMS	tanβ = 30 GeV
>1341	95	[16] SIRUNYAN	18CX	CMS	tanβ = 40 GeV
>1496	95	[16] SIRUNYAN	18CX	CMS	tanβ = 50 GeV
>1613	95	[16] SIRUNYAN	18CX	CMS	tanβ = 60 GeV
		[17] AABOUD	16AA	ATLS	$A^0 \to \tau^+\tau^-$

		[18] KHACHATRY...	16A	CMS	$H^0_{1,2}/A^0 \to \mu^+\mu^-$
		[19] KHACHATRY...	16P	CMS	$H^0_2 \to H^0 H^0$, $A^0 \to Z H^0$
		[20] KHACHATRY...	15AY	CMS	$pp \to H^0_{1,2}/A^0 + b + X$, $H^0_{1,2}/A^0 \to b\overline{b}$
		[21] AAD	14AW	ATLS	$pp \to H^0_{1,2}/A^0 + X$, $H^0_{1,2}/A^0 \to \tau\tau$
		[22] KHACHATRY...	14M	CMS	$pp \to H^0_{1,2}/A^0 + X$, $H^0_{1,2}/A^0 \to \tau\tau$
		[23] AAD	13O	ATLS	$pp \to H^0_{1,2}/A^0 + X$, $H^0_{1,2}/A^0 \to \tau^+\tau^-$, $\mu^+\mu^-$
		[24] AAIJ	13T	LHCB	$pp \to H^0_{1,2}/A^0 + X$, $H^0_{1,2}/A^0 \to \tau^+\tau^-$
		[25] CHATRCHYAN	13AG	CMS	$pp \to H^0_{1,2}/A^0 + b + X$, $H^0_{1,2}/A^0 \to b\overline{b}$
		[26] AALTONEN	12AQ	TEVA	$p\overline{p} \to H^0_{1,2}/A^0 + b + X$, $H^0_{1,2}/A^0 \to b\overline{b}$
		[27] AALTONEN	12X	CDF	$p\overline{p} \to H^0_{1,2}/A^0 + b + X$, $H^0_{1,2}/A^0 \to b\overline{b}$
		[28] ABAZOV	12G	D0	$p\overline{p} \to H^0_{1,2}/A^0 + X$, $H^0_{1,2}/A^0 \to \tau^+\tau^-$
		[29] CHATRCHYAN	12K	CMS	$pp \to H^0_{1,2}/A^0 + X$, $H^0_{1,2}/A^0 \to \tau^+\tau^-$
		[30] ABAZOV	11K	D0	$p\overline{p} \to H^0_{1,2}/A^0 + b + X$, $H^0_{1,2}/A^0 \to b\overline{b}$
		[31] ABAZOV	11W	D0	$p\overline{p} \to H^0_{1,2}/A^0 + b + X$, $H^0_{1,2}/A^0 \to \tau^+\tau^-$
		[32] AALTONEN	09AR	CDF	$p\overline{p} \to H^0_{1,2}/A^0 + X$, $H^0_{1,2}/A^0 \to \tau^+\tau^-$
$>$ 90.4		[33] ABDALLAH	08B	DLPH	$E_{cm} \le 209$ GeV
$>$ 93.4	95	[34] SCHAEL	06B	LEP	$E_{cm} \le 209$ GeV
		[35] ACOSTA	05Q	CDF	$p\overline{p} \to H^0_{1,2}/A^0 + X$
$>$ 85.0	95	[36,37] ABBIENDI	04M	OPAL	$E_{cm} \le 209$ GeV
		[38] ABBIENDI	03G	OPAL	$H^0_1 \to A^0 A^0$
$>$ 86.5	95	[36,39] ACHARD	02H	L3	$E_{cm} \le 209$ GeV, $\tan\beta > 0.4$
		[40] AKEROYD	02	RVUE	
$>$ 90.1	95	[36,41] HEISTER	02	ALEP	$E_{cm} \le 209$ GeV, $\tan\beta > 0.5$

[1] TUMASYAN 23S search for production of $H^0_2/A^0 \to \tau^+\tau^-$ by gluon fusion and b-associated prodution using 138 fb^{-1} of pp collisions at $E_{cm} = 13$ TeV. See their Fig.

13 for excluded regions in the m_{A^0}–tanβ plane in M_h^{125} and M_{hEFT}^{125} MSSM scenarios. In both scenarios $m_{A^0} < 350$ GeV is excluded at 95% CL.

[2] AAD 20AA search for $H_2^0/A^0 \to \tau^+\tau^-$ produced by gluon fusion or b-associated production using 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 2(c) for excluded region in the M_h^{125} scenario of MSSM. Values of tanβ > 8 (21) are excluded for m_{A^0} = 1.0 (1.5) TeV at 95%CL.

[3] AAD 24AP search for production of a heavy H_2^0 and A^0 decaying to $t\overline{t}$ in 140 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 13(b) for excluded parameter regions in hMSSM.

[4] AAD 24CB use ATLAS measurements from up to 139 fb^{-1} of H production cross sections and decays to constrain the MSSM Higgs sector. See their Figs. 22 and 23 for excluded regions in various MSSM scenarios.

[5] AAD 24H combine searches for a scalar resonance decaying to HH using up to 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV from AAD 22F, AAD 23Z, and AAD 22Y. See their Fig. 3 for the excluded region in the MSSM $M_{h,EFT}^{125}$ and $M_{h,EFT}^{125}(\widetilde{\chi})$ benchmark scenarios (where only the resonant H-exchange contribution has been taken into account in the signal modelling).

[6] AAD 20 combine measurements on H production and decay using data taken in years 2015–2017 (up to 79.8 fb^{-1}) of pp collisions at E_{cm} = 13 TeV. See their Fig. 19 for excluded region in the hMSSM parameter space.

[7] AAD 20C combine searches for a scalar resonance decaying to HH in 36.1 fb^{-1} of pp collisions at E_{cm} = 13 TeV from AABOUD 19A, AABOUD 19O, AABOUD 18CQ, AABOUD 19T, AABOUD 18CW, and AABOUD 18BU. See their Fig. 7(b) for the excluded region in the hMSSM parameter space.

[8] AAD 20L search for b-associated production of H_2^0 decaying to $b\overline{b}$ in 27.8 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 9 for excluded regions in hMSSM, $\mathrm{m}_h^{\mathrm{mod}+}$ and $\mathrm{m}_h^{\mathrm{mod}-}$ scenarios of MSSM.

[9] SIRUNYAN 20AC search for gluon-fusion and b-associated production of A^0 decaying to ZH in 35.9 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 6 for excluded regions in the M_{hEFT}^{125} and hMSSM scenarios of the MSSM.

[10] SIRUNYAN 20AF search for $H_2^0/A^0 \to t\overline{t}$ with one or two charged leptons in the final state using kinematic variables in 35.9 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 8 for excluded region in the hMSSM scenario of MSSM. Values of tanβ below 1.0–1.5 are excluded for m_{A^0} = 0.4–0.75 TeV at 95%CL.

[11] SIRUNYAN 20Y search for gluon-fusion and vector-boson-fusion production of H_2^0 decaying to W^+W^- in the final states $\ell\nu\ell\nu$ and $\ell\nu qq$ in 35.9 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Figs. 8 and 9 for excluded regions in various MSSM scenarios.

[12] SIRUNYAN 19CR search for production of H_2^0/A^0 in gluon fusion and in association with a $b\overline{b}$ pair, decaying to $\mu^+\mu^-$ in 35.9 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 5 for the excluded region in the MSSM parameter space in the $m_h^{\mathrm{mod}+}$ and hMSSM scenarios.

[13] AABOUD 18G search for production of $H_2^0/A^0 \to \tau^+\tau^-$ by gluon fusion and b-associated prodution in 36.1 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 10 for excluded regions in the m_{A^0}– tanβ plane in several MSSM scenarios.

[14] SIRUNYAN 18A search for production of a scalar resonance decaying to $H^0H^0 \to b\overline{b}\tau^+\tau^-$ in 35.9 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 5 (lower) for excluded regions in the m_{A^0} – tanβ plane in the hMSSM scenario.

[15] SIRUNYAN 18BP search for production of $H_2^0/A^0 \to b\overline{b}$ by b-associated prodution in 35.7 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 6 for the limits on cross section times branching ratio for $m_{H_2^0}$, m_{A^0} = 0.3–1.3 TeV, and Fig. 7 for excluded regions in the m_{A^0}– tan(β) plane in several MSSM scenarios.

[16] SIRUNYAN 18CX search for production of $H_{1,2}^0/A^0 \to \tau^+\tau^-$ by gluon fusion and b-associated prodution in 35.9 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 9 for excluded regions in the m_{A^0}– tan(β) plane in several MSSM scenarios.

[17] AABOUD 16AA search for production of a Higgs boson in gluon fusion and in association with a $b\overline{b}$ pair followed by the decay $A^0 \to \tau^+\tau^-$ in 3.2 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 5(a, b) for limits on cross section times branching ratio for m_{A^0} = 200–1200 GeV, and Fig. 5(c, d) for the excluded region in the MSSM parameter space in the $m_h^{\text{mod}+}$ and hMSSM scenarios.

[18] KHACHATRYAN 16A search for production of a Higgs boson in gluon fusion and in association with a $b\overline{b}$ pair followed by the decay $H_{1,2}^0/A^0 \to \mu^+\mu^-$ in 5.1 fb^{-1} of pp collisions at E_{cm} = 7 TeV and 19.3 fb^{-1} at E_{cm} = 8 TeV. See their Fig. 7 for the excluded region in the MSSM parameter space in the $m_h^{\text{mod}+}$ benchmark scenario and Fig. 9 for limits on cross section times branching ratio.

[19] KHACHATRYAN 16P search for gluon fusion production of an H_2^0 decaying to $H^0 H^0 \to b\overline{b}\tau^+\tau^-$ and an A^0 decaying to $Z H^0 \to \ell^+\ell^-\tau^+\tau^-$ in 19.7 fb^{-1} of pp collisions at E_{cm} = 8 TeV. See their Fig. 12 for excluded region in the tanβ − cos($\beta-\alpha$) plane for $m_{H_2^0} = m_{A^0}$ = 300 GeV.

[20] KHACHATRYAN 15AY search for production of a Higgs boson in association with a b quark in the decay $H_{1,2}^0/A^0 \to b\overline{b}$ in 19.7 fb^{-1} of pp collisions at E_{cm} = 8 TeV and combine with CHATRCHYAN 13AG 7 TeV data. See their Fig. 6 for the limits on cross section times branching ratio for m_{A^0} = 100–900 GeV and Figs. 7–9 for the excluded region in the MSSM parameter space in various benchmark scenarios.

[21] AAD 14AW search for production of a Higgs boson followed by the decay $H_{1,2}^0/A^0 \to \tau^+\tau^-$ in 19.5–20.3 fb^{-1} of pp collisions at E_{cm} = 8 TeV. See their Fig. 11 for the limits on cross section times branching ratio and their Figs. 9 and 10 for the excluded region in the MSSM parameter space. For m_{A^0} = 140 GeV, the region tan$\beta > 5.4$ is excluded at 95% CL in the m_h^{max} scenario.

[22] KHACHATRYAN 14M search for production of a Higgs boson in gluon fusion and in association with a b quark followed by the decay $H_{1,2}^0/A^0 \to \tau^+\tau^-$ in 4.9 fb^{-1} of pp collisions at E_{cm} = 7 TeV and 19.7 fb^{-1} at E_{cm} = 8 TeV. See their Figs. 7 and 8 for one- and two-dimensional limits on cross section times branching ratio and their Figs. 5 and 6 for the excluded region in the MSSM parameter space. For m_{A^0} = 140 GeV, the region tan$\beta > 3.8$ is excluded at 95% CL in the m_h^{max} scenario.

[23] AAD 13O search for production of a Higgs boson in the decay $H_{1,2}^0/A^0 \to \tau^+\tau^-$ and $\mu^+\mu^-$ with 4.7–4.8 fb^{-1} of pp collisions at E_{cm} = 7 TeV. See their Fig. 6 for the excluded region in the MSSM parameter space and their Fig. 7 for the limits on cross section times branching ratio. For m_{A^0} = 110–170 GeV, tan$\beta \gtrsim 10$ is excluded, and for tan$\beta = 50$, m_{A^0} below 470 GeV is excluded at 95% CL in the m_h^{max} scenario.

[24] AAIJ 13T search for production of a Higgs boson in the forward region in the decay $H_{1,2}^0/A^0 \to \tau^+\tau^-$ in 1.0 fb^{-1} of pp collisions at E_{cm} = 7 TeV. See their Fig. 2 for the limits on cross section times branching ratio and the excluded region in the MSSM parameter space.

[25] CHATRCHYAN 13AG search for production of a Higgs boson in association with a b quark in the decay $H^0_{1,2}/A^0 \rightarrow b\overline{b}$ in 2.7–4.8 fb^{-1} of pp collisions at $E_{\rm cm}$ = 7 TeV. See their Fig. 6 for the excluded region in the MSSM parameter space and Fig. 5 for the limits on cross section times branching ratio. For m_{A^0} = 90–350 GeV, upper bounds on tanβ of 18–42 at 95% CL are obtained in the $m_h^{\rm max}$ scenario with $\mu = +200$ GeV.

[26] AALTONEN 12AQ combine AALTONEN 12X and ABAZOV 11K. See their Table I and Fig. 1 for the limit on cross section times branching ratio and Fig. 2 for the excluded region in the MSSM parameter space.

[27] AALTONEN 12X search for associated production of a Higgs boson and a b quark in the decay $H^0_{1,2}/A^0 \rightarrow b\overline{b}$, with 2.6 fb^{-1} of $p\overline{p}$ collisions at $E_{\rm cm}$ = 1.96 TeV. See their Table III and Fig. 15 for the limit on cross section times branching ratio and Figs. 17, 18 for the excluded region in the MSSM parameter space.

[28] ABAZOV 12G search for production of a Higgs boson in the decay $H^0_{1,2}/A^0 \rightarrow \tau^+\tau^-$ with 7.3 fb^{-1} of $p\overline{p}$ collisions at $E_{\rm cm}$ = 1.96 TeV and combine with ABAZOV 11W and ABAZOV 11K. See their Figs. 4, 5, and 6 for the excluded region in the MSSM parameter space. For m_{A^0} = 90–180 GeV, tan$\beta \gtrsim$ 30 is excluded at 95% CL. in the $m_h^{\rm max}$ scenario.

[29] CHATRCHYAN 12K search for production of a Higgs boson in the decay $H^0_{1,2}/A^0 \rightarrow \tau^+\tau^-$ with 4.6 fb^{-1} of pp collisions at $E_{\rm cm}$ = 7 TeV. See their Fig. 3 and Table 4 for the excluded region in the MSSM parameter space. For m_{A^0} = 160 GeV, the region tanβ > 7.1 is excluded at 95% CL in the $m_h^{\rm max}$ scenario. Superseded by KHACHATRYAN 14M.

[30] ABAZOV 11K search for associated production of a Higgs boson and a b quark, followed by the decay $H^0_{1,2}/A^0 \rightarrow b\overline{b}$, in 5.2 fb^{-1} of $p\overline{p}$ collisions at $E_{\rm cm}$ = 1.96 TeV. See their Fig. 5/Table 2 for the limit on cross section times branching ratio and Fig. 6 for the excluded region in the MSSM parameter space for $\mu = -200$ GeV.

[31] ABAZOV 11W search for associated production of a Higgs boson and a b quark, followed by the decay $H^0_{1,2}/A^0 \rightarrow \tau\tau$, in 7.3 fb^{-1} of $p\overline{p}$ collisions at $E_{\rm cm}$ = 1.96 TeV. See their Fig. 2 for the limit on cross section times branching ratio and for the excluded region in the MSSM parameter space.

[32] AALTONEN 09AR search for Higgs bosons decaying to $\tau^+\tau^-$ in two doublet models in 1.8 fb^{-1} of $p\overline{p}$ collisions at $E_{\rm cm}$ = 1.96 TeV. See their Fig. 2 for the limit on $\sigma \cdot$ B($H^0_{1,2}/A^0 \rightarrow \tau^+\tau^-$) for different Higgs masses, and see their Fig. 3 for the excluded region in the MSSM parameter space.

[33] ABDALLAH 08B give limits in eight CP-conserving benchmark scenarios and some CP-violating scenarios. See paper for excluded regions for each scenario. Supersedes ABDALLAH 04.

[34] SCHAEL 06B make a combined analysis of the LEP data. The quoted limit is for the $m_h^{\rm max}$ scenario with m_t = 174.3 GeV. In the CP-violating CPX scenario no lower bound on $m_{H^0_1}$ can be set at 95% CL. See paper for excluded regions in various scenarios. See Figs. 2–6 and Tabs. 14–21 for limits on $\sigma(Z H^0)\cdot$ B($H^0 \rightarrow b\overline{b}, \tau^+\tau^-$) and $\sigma(H^0_1 H^0_2)\cdot$ B($H^0_1,H^0_2 \rightarrow b\overline{b},\tau^+\tau^-$).

[35] ACOSTA 05Q search for $H^0_{1,2}/A^0$ production in $p\overline{p}$ collisions at $E_{\rm cm}$ = 1.8 TeV with $H^0_{1,2}/A^0 \rightarrow \tau^+\tau^-$. At m_{A^0} = 100 GeV, the obtained cross section upper limit is above theoretical expectation.

[36] Search for $e^+e^- \rightarrow H^0_1 A^0$ in the final states $b\overline{b}b\overline{b}$ and $b\overline{b}\tau^+\tau^-$, and $e^+e^- \rightarrow H^0_1 Z$. Universal scalar mass of 1 TeV, SU(2) gaugino mass of 200 GeV, and $\mu = -200$

GeV are assumed, and two-loop radiative corrections incorporated. The limits hold for m_t=175 GeV, and for the m_h^{max} scenario.

[37] ABBIENDI 04M exclude $0.7 < \tan\beta < 1.9$, assuming $m_t = 174.3$ GeV. Limits for other MSSM benchmark scenarios, as well as for CP violating cases, are also given.

[38] ABBIENDI 03G search for $e^+e^- \to H_1^0 Z$ followed by $H_1^0 \to A^0 A^0$, $A^0 \to c\overline{c}$, gg, or $\tau^+\tau^-$. In the no-mixing scenario, the region $m_{H_1^0}$ = 45-85 GeV and m_{A^0} = 2-9.5 GeV is excluded at 95% CL.

[39] ACHARD 02H also search for the final state $H_1^0 Z \to 2A^0 q\overline{q}$, $A^0 \to q\overline{q}$. In addition, the MSSM parameter set in the "large-μ" and "no-mixing" scenarios are examined.

[40] AKEROYD 02 examine the possibility of a light A^0 with $\tan\beta < 1$. Electroweak measurements are found to be inconsistent with such a scenario.

[41] HEISTER 02 excludes the range $0.7 < \tan\beta < 2.3$. A wider range is excluded with different stop mixing assumptions. Updates BARATE 01C.

Mass Limits for H_1^0 (Higgs Boson) in Supersymmetric Models

VALUE (GeV)	CL%	DOCUMENT ID		TECN	COMMENT
>89.7		[1] ABDALLAH	08B	DLPH	$E_{cm} \le 209$ GeV
>92.8	95	[2] SCHAEL	06B	LEP	$E_{cm} \le 209$ GeV
>84.5	95	[3,4] ABBIENDI	04M	OPAL	$E_{cm} \le 209$ GeV
>86.0	95	[3,5] ACHARD	02H	L3	$E_{cm} \le 209$ GeV, $\tan\beta > 0.4$
>89.8	95	[3,6] HEISTER	02	ALEP	$E_{cm} \le 209$ GeV, $\tan\beta > 0.5$

• • • We do not use the following data for averages, fits, limits, etc. • • •

		[7] AALTONEN	12AQ	TEVA	$p\overline{p} \to H_{1,2}^0/A^0 + b + X$, $H_{1,2}^0/A^0 \to b\overline{b}$

[1] ABDALLAH 08B give limits in eight CP-conserving benchmark scenarios and some CP-violating scenarios. See paper for excluded regions for each scenario. Supersedes ABDALLAH 04.

[2] SCHAEL 06B make a combined analysis of the LEP data. The quoted limit is for the m_h^{max} scenario with $m_t = 174.3$ GeV. In the CP-violating CPX scenario no lower bound on $m_{H_1^0}$ can be set at 95% CL. See paper for excluded regions in various scenarios. See Figs. 2–6 and Tabs. 14–21 for limits on $\sigma(ZH^0)\cdot$ B($H^0 \to b\overline{b}, \tau^+\tau^-$) and $\sigma(H_1^0 H_2^0)\cdot$ B($H_1^0,H_2^0 \to b\overline{b},\tau^+\tau^-$).

[3] Search for $e^+e^- \to H_1^0 A^0$ in the final states $b\overline{b}b\overline{b}$ and $b\overline{b}\tau^+\tau^-$, and $e^+e^- \to H_1^0 Z$. Universal scalar mass of 1 TeV, SU(2) gaugino mass of 200 GeV, and $\mu = -200$ GeV are assumed, and two-loop radiative corrections incorporated. The limits hold for m_t=175 GeV, and for the m_h^{max} scenario.

[4] ABBIENDI 04M exclude $0.7 < \tan\beta < 1.9$, assuming $m_t = 174.3$ GeV. Limits for other MSSM benchmark scenarios, as well as for CP violating cases, are also given.

[5] ACHARD 02H also search for the final state $H_1^0 Z \to 2A^0 q\overline{q}$, $A^0 \to q\overline{q}$. In addition, the MSSM parameter set in the "large-μ" and "no-mixing" scenarios are examined.

[6] HEISTER 02 excludes the range $0.7 < \tan\beta < 2.3$. A wider range is excluded with different stop mixing assumptions. Updates BARATE 01C.

[7] AALTONEN 12AQ combine AALTONEN 12X and ABAZOV 11K. See their Table I and Fig. 1 for the limit on cross section times branching ratio and Fig. 2 for the excluded region in the MSSM parameter space.

MASS LIMITS FOR NEUTRAL HIGGS BOSONS IN EXTENDED HIGGS MODELS

This Section covers models which do not fit into either the Standard Model or its simplest minimal Supersymmetric extension (MSSM), leading to anomalous production rates, or nonstandard final states and branching ratios. In particular, this Section covers limits which may apply to generic two-Higgs-doublet models (2HDM), or to special regions of the MSSM parameter space where decays to invisible particles or to photon pairs are dominant (see the review on "Status of Higgs Boson Physics"). Concerning the mass limits for H^0 and A^0 listed below, see the footnotes or the comment lines for details on the nature of the models to which the limits apply.

The observed signal at about 125 GeV, see section "H", can be interpreted as one of the neutral Higgs bosons of an extended Higgs sector.

Mass Limits in General two-Higgs-doublet Models

VALUE (GeV)	CL%	DOCUMENT ID		TECN	COMMENT
• • • We do not use the following data for averages, fits, limits, etc. • • •					
		[1] AAD	24AB	ATLS	$A^0 \rightarrow Z H_2^0$
		[2] AAD	24AP	ATLS	H_2^0, $A^0 \rightarrow t\overline{t}$
		[3] AAD	24CB	ATLS	H properties
		[4] AAD	24H	ATLS	$H_2^0 \rightarrow HH$
		[5] HAYRAPETY...	24H	CMS	flavor changing H_2^0, A^0
		[6] AAD	23AD	ATLS	$A^0 \rightarrow Z H_2^0$, $H_2^0 \rightarrow HH$
		[7] AAD	23BG	ATLS	$t\overline{t} H_2^0/A^0$
		[8] AAD	23O	ATLS	$A^0 \rightarrow ZH$
		[9] AAD	21AF	ATLS	$H_2^0 \rightarrow ZZ$
		[10] AAD	21AI	ATLS	$A^0 \rightarrow Z H_2^0$
		[11] AAD	20	ATLS	H^0 properties
		[12] AAD	20L	ATLS	$H_2^0 \rightarrow b\overline{b}$
		[13] SIRUNYAN	20AA	CMS	$H_2^0 \rightarrow Z A^0$ or $A^0 \rightarrow Z H_2^0$
		[14] SIRUNYAN	20Y	CMS	$H_2^0 \rightarrow W^+ W^-$
		[15] SIRUNYAN	19AE	CMS	$A^0 \rightarrow \tau^+ \tau^-$
		[16] SIRUNYAN	19AV	CMS	$A^0 \rightarrow Z H^0$
		[17] AABOUD	18AH	ATLS	$A^0 \rightarrow Z H_2^0$
		[18] AABOUD	18AI	ATLS	$A^0 \rightarrow Z H^0$
		[19] AABOUD	18BF	ATLS	$H_2^0 \rightarrow ZZ$
		[20] AABOUD	18CE	ATLS	$pp \rightarrow H_2^0/A^0 t\overline{t}$, $H_2^0/A^0 \rightarrow t\overline{t}$
		[21] HALLER	18	RVUE	global fits
		[22] SIRUNYAN	18BP	CMS	$pp \rightarrow H_2^0/A^0 + b + X$, $H_2^0/A^0 \rightarrow b\overline{b}$
		[23] SIRUNYAN	18ED	CMS	$A^0 \rightarrow Z H^0$
		[24] AABOUD	17AN	ATLS	H_2^0, $A^0 \rightarrow t\overline{t}$
		[25] SIRUNYAN	17AX	CMS	$A^0 b\overline{b}$, $A^0 \rightarrow \mu^+ \mu^-$
		[26] AAD	16AX	ATLS	$H_2^0 \rightarrow ZZ$

		[27] KHACHATRY...	16P	CMS	$H_2^0 \rightarrow H^0 H^0$, $A^0 \rightarrow Z H^0$
		[28] KHACHATRY...	16W	CMS	$A^0 b\overline{b}$, $A^0 \rightarrow \tau^+\tau^-$
		[29] KHACHATRY...	16Z	CMS	$H_2^0 \rightarrow Z A^0$ or $A^0 \rightarrow Z H_2^0$
		[30] AAD	15BK	ATLS	$H_2^0 \rightarrow H^0 H^0$
		[31] AAD	15S	ATLS	$A^0 \rightarrow Z H^0$
		[32] KHACHATRY...	15BB	CMS	H_2^0, $A^0 \rightarrow \gamma\gamma$
		[33] KHACHATRY...	15N	CMS	$A^0 \rightarrow Z H^0$
		[34] AAD	14M	ATLS	$H_2^0 \rightarrow H^\pm W^\mp \rightarrow H^0 W^\pm W^\mp$, $H^0 \rightarrow b\overline{b}$
		[35] KHACHATRY...	14Q	CMS	$H_2^0 \rightarrow H^0 H^0$, $A^0 \rightarrow Z H^0$
		[36] AALTONEN	09AR	CDF	$p\overline{p} \rightarrow H_{1,2}^0/A^0 + X$, $H_{1,2}^0/A^0 \rightarrow \tau^+\tau^-$
none 1–55	95	[37] ABBIENDI	05A	OPAL	H_1^0, Type II model
>110.6	95	[38] ABDALLAH	05D	DLPH	$H^0 \rightarrow$ 2 jets
		[39] ABDALLAH	04O	DLPH	$Z \rightarrow f\overline{f}H$
		[40] ABDALLAH	04O	DLPH	$e^+e^- \rightarrow H^0 Z$, $H^0 A^0$
		[41] ABBIENDI	02D	OPAL	$e^+e^- \rightarrow b\overline{b}H$
none 1–44	95	[42] ABBIENDI	01E	OPAL	H_1^0, Type-II model
> 68.0	95	[43] ABBIENDI	99E	OPAL	$\tan\beta > 1$
		[44] ABREU	95H	DLPH	$Z \rightarrow H^0 Z^*$, $H^0 A^0$
		[45] PICH	92	RVUE	Very light Higgs

[1] AAD 24AB search for gluon-fusion and $b\overline{b}$ associated production of A^0 decaying to $A^0 \rightarrow Z H_2^0$, in the final states $\ell^+\ell^- t\overline{t}$ and $\nu\overline{\nu} b\overline{b}$ using 140 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Figs. 5 and 6 for limits on cross section times branching ratios in terms of two-Higgs-doublet models, and Fig. 7 for excluded model parameter space.

[2] AAD 24AP search for production of a heavy H_2^0 and A^0 decaying to $t\overline{t}$ in 140 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 13(a) for excluded parameter regions in a Type II two-Higgs-doublet model.

[3] AAD 24CB use ATLAS measurements from up to 139 fb^{-1} of H production cross sections and decays to constrain two Higgs doublet models. See their Figs. 19–21 for excluded regions of 2HDM parameter space.

[4] AAD 24H combine searches for a scalar resonance decaying to HH using up to 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV from AAD 22F, AAD 23Z, and AAD 22Y. See their Fig. 2 for the excluded region in the Type-I two-Higgs-doublet model parameter space (where only the resonant H-exchange contribution has been taken into account in the signal modelling).

[5] HAYRAPETYAN 24H search for H_2^0 and A^0 having flavor-violating couplings to tc or tu, produced in association with top quark(s), using 138 fb^{-1} of pp collisions at E_{cm} = 13 TeV. For the interpretation of the data in terms of flavor-violating two-Higs-doublet models, given in the range m_{A^0} = 0.2–1.0 TeV, see their Figs. 4, 5, and 6.

[6] AAD 23AD search for associated production of $W/Z H_2^0$ and gluon fusion production of A^0 decaying to $Z H_2^0$, with the decay chain $H_2^0 \rightarrow HH \rightarrow b\overline{b}b\overline{b}$, using 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Figs. 12 and 13 for excluded regions in Type-I and lepton-specific 2HDMs.

[7] AAD 23BG search for production of H_2^0/A^0 in association with a $t\overline{t}$ pair, decaying to $t\overline{t}$, using 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 8 for excluded regions in the parameter space of the type II 2HDM.

[8] AAD 23O search for production of an A^0 in gluon-gluon fusion and in association with a $b\overline{b}$, decaying to ZH in the final states $\nu\overline{\nu} b\overline{b}$ and $\ell^+\ell^- b\overline{b}$ using 139 fb^{-1} of pp

collisions at $E_{\rm cm}$ = 13 TeV. See their Figs. 12 and 13 for excluded regions in the parameter space in various 2HDMs.

[9] AAD 21AF search for production of a heavy H_2^0 state decaying to ZZ in the final states $\ell^+\ell^-\ell'^+\ell'^-$ and $\ell^+\ell^-\nu\overline{\nu}$ in 139 fb^{-1} of pp collisions at $E_{\rm cm}$ = 13 TeV. See their Figs. 6 and 7 for excluded parameter regions of the 2HDM Type I and II.

[10] AAD 21AI search for production of an A^0 in gluon-gluon fusion and in association with a $b\overline{b}$, decaying to $ZH_2^0 \rightarrow \ell^+\ell^- b\overline{b}$ or $\ell^+\ell^- W^+W^-$ in 139 fb^{-1} of pp collisions at $E_{\rm cm}$ = 13 TeV. See their Figs. 10 and 14 for excluded regions in the parameter space of various 2HDMs.

[11] AAD 20 combine measurements on H^0 production and decay using data taken in years 2015–2017 (up to 79.8 fb^{-1}) of pp collisions at $E_{\rm cm}$ = 13 TeV. See their Fig. 18 for excluded regions in various 2HDMs.

[12] AAD 20L search for b-associated production of H_2^0 decaying to $b\overline{b}$ in 27.8 fb^{-1} of pp collisions at $E_{\rm cm}$ = 13 TeV. See their Figs. 10 and 11 for excluded regions in the flipped two Higgs doublet model.

[13] SIRUNYAN 20AA search for $H_2^0 \rightarrow ZA^0$, $A^0 \rightarrow b\overline{b}$ or $A^0 \rightarrow ZH_2^0$, $H_2^0 \rightarrow b\overline{b}$ in 35.9 fb^{-1} of pp collisions at $E_{\rm cm}$ = 13 TeV. See their Figs. 8 and 9 for excluded regions in the parameter space of Type-II two Higgs doublet model.

[14] SIRUNYAN 20Y search for gluon-fusion and vector-boson-fusion production of H_2^0 decaying to W^+W^- in the final states $\ell\nu\ell\nu$ and $\ell\nu qq$ in 35.9 fb^{-1} of pp collisions at $E_{\rm cm}$ = 13 TeV. See their Fig. 7 for excluded regions in Type I and II two Higgs doublet models.

[15] SIRUNYAN 19AE search for a pseudoscalar resonance produced in association with a $b\overline{b}$ pair, decaying to $\tau^+\tau^-$ in 35.9 fb^{-1} of pp collisions at $E_{\rm cm}$ = 13 TeV. See their Fig. 4 for cross section limits for m_{A^0} = 25–70 GeV and comparison with some representative 2HDMs.

[16] SIRUNYAN 19AV search for a scalar resonance produced by gluon fusion or b associated production, decaying to $ZH^0 \rightarrow \ell^+\ell^- b\overline{b}$ ($\ell = e$, μ) or $\nu\overline{\nu}b\overline{b}$ in 35.9 fb^{-1} of pp collisions at $E_{\rm cm}$ = 13 TeV. See their Figs. 6 and 7 for excluded regions in the parameter space of various 2HDMs.

[17] AABOUD 18AH search for production of an A^0 in gluon-gluon fusion and in association with a $b\overline{b}$, decaying to $ZH_2^0 \rightarrow \ell^+\ell^- b\overline{b}$ in 36.1 fb^{-1} of pp collisions at $E_{\rm cm}$ = 13 TeV. See their Fig. 6 for excluded regions in the parameter space of various 2HDMs.

[18] AABOUD 18AI search for production of an A^0 in gluon-gluon fusion and in association with a $b\overline{b}$, decaying to ZH^0 in the final states $\nu\overline{\nu}b\overline{b}$ and $\ell^+\ell^- b\overline{b}$ in 36.1 fb^{-1} of pp collisions at $E_{\rm cm}$ = 13 TeV. See their Figs. 7 and 8 for excluded regions in the parameter space in various 2HDMs.

[19] AABOUD 18BF search for production of a heavy H_2^0 state decaying to ZZ in the final states $\ell^+\ell^-\ell^+\ell^-$ and $\ell^+\ell^-\nu\overline{\nu}$ in 36.1 fb^{-1} of pp collisions at $E_{\rm cm}$ = 13 TeV. See their Figs. 8 and 9 for excluded parameter regions in 2HDM Type I and II.

[20] AABOUD 18CE search for the process $pp \rightarrow H_2^0/A^0\, t\overline{t}$ followed by the decay $H_2^0/A^0 \rightarrow t\overline{t}$ in 36.1 fb^{-1} of pp collisions at $E_{\rm cm}$ = 13 TeV. See their Fig. 12 for limits on cross section times branching ratio, and for lower limits on tanβ for $m_{H_2^0}$, m_{A^0} = 0.4–1.0 TeV in the 2HDM type II.

[21] HALLER 18 perform global fits in the framework of two-Higgs-doublet models (type I, II, lepton specific, flipped). See their Fig. 8 for allowed parameter regions from fits to LHC H^0 measurements, Fig. 9 bottom and charm decays, Fig. 10 muon anomalous magnetic moment, Fig. 11 electroweak precision data, and Fig. 12 by combination of all data.

[22] SIRUNYAN 18BP search for production of $H_2^0/A^0 \rightarrow b\overline{b}$ by b-associated prodution in 35.7 fb^{-1} of pp collisions at $E_{\rm cm}$ = 13 TeV. See their Fig. 6 for the limits on cross

section times branching ratio for $m_{H_2^0}$, m_{A^0} = 0.3–1.3 TeV, and Figs. 8 and 9 for excluded regions in the parameter space of type-II and flipped 2HDMs.

[23] SIRUNYAN 18ED search for production of an A^0 in gluon-gluon fusion and in association with a $b\overline{b}$, decaying to $Z H^0$ in the final states $\nu\overline{\nu} b\overline{b}$ or $\ell^+\ell^- b\overline{b}$ in 35.9 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 9 for excluded regions in the parameter space in Type I and II 2HDMs.

[24] AABOUD 17AN search for production of a heavy H_2^0 and/or A^0 decaying to $t\overline{t}$ in 20.3 fb^{-1} of pp collisions at E_{cm} = 8 TeV. See their Fig. 3 and Table III for excluded parameter regions in Type II Two-Higgs-Doublet-Models.

[25] SIRUNYAN 17AX search for $A^0 b\overline{b}$ production followed by the decay $A^0 \rightarrow \mu^+\mu^-$ in 19.7 fb^{-1} of pp collisions at E_{cm} = 8 TeV. Limits are set in the range m_{A^0} = 25–60 GeV. See their Fig. 5 for upper limits on $\sigma(A^0 b\overline{b})\cdot$B$(A^0 \rightarrow \mu^+\mu^-)$.

[26] AAD 16AX search for production of a heavy H^0 state decaying to ZZ in the final states $\ell^+\ell^-\ell^+\ell^-$, $\ell^+\ell^-\nu\overline{\nu}$, $\ell^+\ell^- q\overline{q}$, and $\nu\overline{\nu} q\overline{q}$ in 20.3 fb^{-1} of pp collisions at E_{cm} = 8 TeV. See their Figs. 13 and 14 for excluded parameter regions in Type I and II models.

[27] KHACHATRYAN 16P search for gluon fusion production of an H_2^0 decaying to $H^0 H^0 \rightarrow b\overline{b}\tau^+\tau^-$ and an A^0 decaying to $Z H^0 \rightarrow \ell^+\ell^-\tau^+\tau^-$ in 19.7 fb^{-1} of pp collisions at E_{cm} = 8 TeV. See their Fig. 11 for limits on tanβ for m_{A^0} = 230–350 GeV.

[28] KHACHATRYAN 16W search for $A^0 b\overline{b}$ production followed by the decay $A^0 \rightarrow \tau^+\tau^-$ in 19.7 fb^{-1} of pp collisions at E_{cm} = 8 TeV. See their Fig. 3 for upper limits on $\sigma(A^0 b\overline{b})\cdot$B$(A^0 \rightarrow \tau^+\tau^-)$.

[29] KHACHATRYAN 16Z search for $H_2^0 \rightarrow Z A^0$ followed by $A^0 \rightarrow b\overline{b}$ or $\tau^+\tau^-$, and $A^0 \rightarrow Z H_2^0$ followed by $H_2^0 \rightarrow b\overline{b}$ or $\tau^+\tau^-$, in 19.8 fb^{-1} of pp collisions at E_{cm} = 8 TeV. See their Fig. 4 for cross section limits and Fig. 5 for excluded region in the parameter space.

[30] AAD 15BK search for production of a heavy H_2^0 decaying to $H^0 H^0$ in the final state $b\overline{b}b\overline{b}$ in 19.5 fb^{-1} of pp collisions at E_{cm} = 8 TeV. See their Figs. 15–18 for excluded regions in the parameter space.

[31] AAD 15S search for production of A^0 decaying to $Z H^0 \rightarrow \ell^+\ell^- b\overline{b}$, $\nu\overline{\nu} b\overline{b}$ and $\ell^+\ell^-\tau^+\tau^-$ in 20.3 fb^{-1} of pp collisions at E_{cm} = 8 TeV. See their Figs. 4 and 5 for excluded regions in the parameter space.

[32] KHACHATRYAN 15BB search for H_2^0 , $A^0 \rightarrow \gamma\gamma$ in 19.7 fb^{-1} of pp collisions at E_{cm} = 8 TeV. See their Fig. 10 for excluded regions in the two-Higgs-doublet model parameter space.

[33] KHACHATRYAN 15N search for production of A^0 decaying to $Z H^0 \rightarrow \ell^+\ell^- b\overline{b}$ in 19.7 fb^{-1} of pp collisions at E_{cm} = 8 TeV. See their Fig. 5 for excluded regions in the tanβ − cos$(\beta-\alpha)$ plane for m_{A^0} = 300 GeV.

[34] AAD 14M search for the decay cascade $H_2^0 \rightarrow H^\pm W^\mp \rightarrow H^0 W^\pm W^\mp$, H^0 decaying to $b\overline{b}$ in 20.3 fb^{-1} of pp collisions at E_{cm} = 8 TeV. See their Table IV for limits in a two-Higgs-doublet model for $m_{H_2^0}$= 325–1025 GeV and m_{H^+}= 225–825 GeV.

[35] KHACHATRYAN 14Q search for $H_2^0 \rightarrow H^0 H^0$ and $A^0 \rightarrow Z H^0$ in 19.5 fb^{-1} of pp collisions at E_{cm} = 8 TeV. See their Figs. 4 and 5 for limits on cross section times branching ratio for m_{H_2,A^0}= 260–360 GeV and their Figs. 7–9 for limits in two-Higgs-doublet models.

[36] AALTONEN 09AR search for Higgs bosons decaying to $\tau^+\tau^-$ in two doublet models in 1.8 fb^{-1} of $p\overline{p}$ collisions at E_{cm} = 1.96 TeV. See their Fig. 2 for the limit on $\sigma \cdot$ B$(H_{1,2}^0/A^0 \rightarrow \tau^+\tau^-)$ for different Higgs masses, and see their Fig. 3 for the excluded region in the MSSM parameter space.

[37] ABBIENDI 05A search for $e^+e^- \to H_1^0 A^0$ in general Type-II two-doublet models, with decays H_1^0, $A^0 \to q\overline{q}$, gg, $\tau^+\tau^-$, and $H_1^0 \to A^0 A^0$.

[38] ABDALLAH 05D search for $e^+e^- \to H^0 Z$ and $H^0 A^0$ with H^0, A^0 decaying to two jets of any flavor including gg. The limit is for SM $H^0 Z$ production cross section with $B(H^0 \to jj) = 1$.

[39] ABDALLAH 04O search for $Z \to b\overline{b}H^0$, $b\overline{b}A^0$, $\tau^+\tau^- H^0$ and $\tau^+\tau^- A^0$ in the final states $4b$, $b\overline{b}\tau^+\tau^-$, and 4τ. See paper for limits on Yukawa couplings.

[40] ABDALLAH 04O search for $e^+e^- \to H^0 Z$ and $H^0 A^0$, with H^0, A^0 decaying to $b\overline{b}$, $\tau^+\tau^-$, or $H^0 \to A^0 A^0$ at E_{cm} = 189–208 GeV. See paper for limits on couplings.

[41] ABBIENDI 02D search for $Z \to b\overline{b}H_1^0$ and $b\overline{b}A^0$ with $H_1^0/A^0 \to \tau^+\tau^-$, in the range $4 < m_H < 12$ GeV. See their Fig. 8 for limits on the Yukawa coupling.

[42] ABBIENDI 01E search for neutral Higgs bosons in general Type-II two-doublet models, at $E_{cm} \le 189$ GeV. In addition to usual final states, the decays H_1^0, $A^0 \to q\overline{q}$, gg are searched for. See their Figs. 15,16 for excluded regions.

[43] ABBIENDI 99E search for $e^+e^- \to H^0 A^0$ and $H^0 Z$ at E_{cm} = 183 GeV. The limit is with $m_H = m_A$ in general two Higgs-doublet models. See their Fig. 18 for the exclusion limit in the m_H–m_A plane. Updates the results of ACKERSTAFF 98S.

[44] See Fig. 4 of ABREU 95H for the excluded region in the $m_{H^0} - m_{A^0}$ plane for general two-doublet models. For $\tan\beta > 1$, the region $m_{H^0} + m_{A^0} \lesssim 87$ GeV, $m_{H^0} < 47$ GeV is excluded at 95% CL.

[45] PICH 92 analyse H^0 with $m_{H^0} < 2m_\mu$ in general two-doublet models. Excluded regions in the space of mass-mixing angles from LEP, beam dump, and $\pi^\pm$, η rare decays are shown in Figs. 3,4. The considered mass region is not totally excluded.

Mass Limits for H^0 with Vanishing Yukawa Couplings

These limits assume that H^0 couples to gauge bosons with the same strength as the Standard Model Higgs boson, but has no coupling to quarks and leptons (this is often referred to as "fermiophobic").

VALUE (GeV)	*CL%*	*DOCUMENT ID*		*TECN*	*COMMENT*
• • • We do not use the following data for averages, fits, limits, etc. • • •					
		[1] AALTONEN	13K	CDF	$H^0 \to WW^{(*)}$
none 100–113	95	[2] AALTONEN	13L	CDF	$H^0 \to \gamma\gamma$, WW^*, ZZ^*
none 100–116	95	[3] AALTONEN	13M	TEVA	$H^0 \to \gamma\gamma$, WW^*, ZZ^*
		[4] ABAZOV	13G	D0	$H^0 \to WW^{(*)}$
none 100–113	95	[5] ABAZOV	13H	D0	$H^0 \to \gamma\gamma$
		[6] ABAZOV	13I	D0	$H^0 \to WW^{(*)}$
		[7] ABAZOV	13J	D0	$H^0 \to WW^{(*)}$, $ZZ^{(*)}$
none 100–114	95	[8] ABAZOV	13L	D0	$H^0 \to \gamma\gamma$, WW^*, ZZ^*
none 110–147	95	[9] CHATRCHYAN	13AL	CMS	$H^0 \to \gamma\gamma$
none 110–118, 119.5–121	95	[10] AAD	12N	ATLS	$H^0 \to \gamma\gamma$
none 100–114	95	[11] AALTONEN	12AN	CDF	$H^0 \to \gamma\gamma$
none 110–194	95	[12] CHATRCHYAN	12AO	CMS	$H^0 \to \gamma\gamma$, $WW^{(*)}$, $ZZ^{(*)}$
none 70–106	95	[13] AALTONEN	09AB	CDF	$H^0 \to \gamma\gamma$
none 70–100	95	[14] ABAZOV	08U	D0	$H^0 \to \gamma\gamma$
>105.8	95	[15] SCHAEL	07	ALEP	$e^+e^- \to H^0 Z$, $H^0 \to WW^*$
>104.1	95	[16,17] ABDALLAH	04L	DLPH	$e^+e^- \to H^0 Z$, $H^0 \to \gamma\gamma$
>107	95	[18] ACHARD	03C	L3	$H^0 \to WW^*$,ZZ^*, $\gamma\gamma$
>105.5	95	[16,19] ABBIENDI	02F	OPAL	$H^0 \to \gamma\gamma$

>105.4	95	[20] ACHARD	02C	L3	$H^0 \to \gamma\gamma$
none 60–82	95	[21] AFFOLDER	01H	CDF	$p\overline{p} \to H^0 W/Z$, $H^0 \to \gamma\gamma$
> 94.9	95	[22] ACCIARRI	00S	L3	$e^+e^- \to H^0 Z$, $H^0 \to \gamma\gamma$
>100.7	95	[23] BARATE	00L	ALEP	$e^+e^- \to H^0 Z$, $H^0 \to \gamma\gamma$
> 96.2	95	[24] ABBIENDI	99O	OPAL	$e^+e^- \to H^0 Z$, $H^0 \to \gamma\gamma$
> 78.5	95	[25] ABBOTT	99B	D0	$p\overline{p} \to H^0 W/Z$, $H^0 \to \gamma\gamma$
		[26] ABREU	99P	DLPH	$e^+e^- \to H^0\gamma$ and/or $H^0 \to \gamma\gamma$

[1] AALTONEN 13K search for $H^0 \to WW^{(*)}$ in 9.7 fb^{-1} of $p\overline{p}$ collisions at E_{cm} = 1.96 TeV. A limit on cross section times branching ratio which corresponds to (1.3–6.6) times the expected cross section is given in the range m_{H^0} = 110–200 GeV at 95% CL.

[2] AALTONEN 13L combine all CDF searches with 9.45–10.0 fb^{-1} of $p\overline{p}$ collisions at E_{cm} = 1.96 TeV.

[3] AALTONEN 13M combine all Tevatron data from the CDF and D0 Collaborations of $p\overline{p}$ collisions at E_{cm} = 1.96 TeV.

[4] ABAZOV 13G search for $H^0 \to WW^{(*)}$ in 9.7 fb^{-1} of $p\overline{p}$ collisions at E_{cm} = 1.96 TeV. A limit on cross section times branching ratio which corresponds to (2–9) times the expected cross section is given for m_{H^0} = 100–200 GeV at 95% CL.

[5] ABAZOV 13H search for $H^0 \to \gamma\gamma$ in 9.6 fb^{-1} of $p\overline{p}$ collisions at E_{cm} = 1.96 TeV.

[6] ABAZOV 13I search for H^0 production in the final state with one lepton and two or more jets plus missing E_T in 9.7 fb^{-1} of $p\overline{p}$ collisions at E_{cm} = 1.96 TeV. The search is sensitive to WH^0, ZH^0 and vector-boson fusion Higgs production with $H^0 \to WW^{(*)}$. A limit on cross section times branching ratio which corresponds to (8–30) times the expected cross section is given in the range m_{H^0} = 100–200 GeV at 95% CL.

[7] ABAZOV 13J search for H^0 production in the final states $ee\mu$, $e\mu\mu$, $\mu\tau\tau$, and $e^\pm\mu^\pm$ in 8.6–9.7 fb^{-1} of $p\overline{p}$ collisions at E_{cm} = 1.96 TeV. The search is sensitive to WH^0, ZH^0 production with $H^0 \to WW^{(*)}$, $ZZ^{(*)}$, decaying to leptonic final states. A limit on cross section times branching ratio which corresponds to (2.4–13.0) times the expected cross section is given in the range m_{H^0} = 100–200 GeV at 95% CL.

[8] ABAZOV 13L combine all D0 results with up to 9.7 fb^{-1} of $p\overline{p}$ collisions at E_{cm} = 1.96 TeV.

[9] CHATRCHYAN 13AL search for $H^0 \to \gamma\gamma$ in 5.1 fb^{-1} and 5.3 fb^{-1} of pp collisions at E_{cm} = 7 and 8 TeV.

[10] AAD 12N search for $H^0 \to \gamma\gamma$ with 4.9 fb^{-1} of pp collisions at E_{cm} = 7 TeV in the mass range m_{H^0} = 110–150 GeV.

[11] AALTONEN 12AN search for $H^0 \to \gamma\gamma$ with 10 fb^{-1} of $p\overline{p}$ collisions at E_{cm} = 1.96 TeV in the mass range m_{H^0} = 100–150 GeV.

[12] CHATRCHYAN 12AO use data from CHATRCHYAN 12G, CHATRCHYAN 12E, CHATRCHYAN 12H, CHATRCHYAN 12I, CHATRCHYAN 12D, and CHATRCHYAN 12C.

[13] AALTONEN 09AB search for $H^0 \to \gamma\gamma$ in 3.0 fb^{-1} of $p\overline{p}$ collisions at E_{cm} = 1.96 TeV in the mass range m_{H^0} = 70–150 GeV. Associated $H^0 W$, $H^0 Z$ production and WW, ZZ fusion are considered.

[14] ABAZOV 08U search for $H^0 \to \gamma\gamma$ in $p\overline{p}$ collisions at E_{cm} = 1.96 TeV in the mass range m_{H^0} = 70–150 GeV. Associated $H^0 W$, $H^0 Z$ production and WW, ZZ fusion are considered. See their Tab. 1 for the limit on $\sigma \cdot$ B($H^0 \to \gamma\gamma$), and see their Fig. 3 for the excluded region in the m_{H^0} — B($H^0 \to \gamma\gamma$) plane.

[15] SCHAEL 07 search for Higgs bosons in association with a fermion pair and decaying to WW^*. The limit is from this search and HEISTER 02L for a H^0 with SM production cross section.

[16] Search for associated production of a $\gamma\gamma$ resonance with a Z boson, followed by $Z \to q\overline{q}$, $\ell^+\ell^-$, or $\nu\overline{\nu}$, at $E_{cm} \leq 209$ GeV. The limit is for a H^0 with SM production cross section.

[17] Updates ABREU 01F.

[18] ACHARD 03C search for $e^+e^- \to ZH^0$ followed by $H^0 \to WW^*$ or ZZ^* at E_{cm}= 200-209 GeV and combine with the ACHARD 02C result. The limit is for a H^0 with SM production cross section. For B($H^0 \to WW^*$) + B($H^0 \to ZZ^*$) = 1, $m_{H^0} > 108.1$ GeV is obtained. See fig. 6 for the limits under different BR assumptions.

[19] For B($H^0 \to \gamma\gamma$)=1, m_{H^0} >117 GeV is obtained.

[20] ACHARD 02C search for associated production of a $\gamma\gamma$ resonance with a Z boson, followed by $Z \to q\overline{q}$, $\ell^+\ell^-$, or $\nu\overline{\nu}$, at $E_{cm} \leq 209$ GeV. The limit is for a H^0 with SM production cross section. For B($H^0 \to \gamma\gamma$)=1, m_{H^0} >114 GeV is obtained.

[21] AFFOLDER 01H search for associated production of a $\gamma\gamma$ resonance and a W or Z (tagged by two jets, an isolated lepton, or missing E_T). The limit assumes Standard Model values for the production cross section and for the couplings of the H^0 to W and Z bosons. See their Fig. 11 for limits with B($H^0 \to \gamma\gamma$)< 1.

[22] ACCIARRI 00S search for associated production of a $\gamma\gamma$ resonance with a $q\overline{q}$, $\nu\overline{\nu}$, or $\ell^+\ell^-$ pair in e^+e^- collisions at E_{cm}= 189 GeV. The limit is for a H^0 with SM production cross section. For B($H^0 \to \gamma\gamma$)=1, $m_{H^0} > 98$ GeV is obtained. See their Fig. 5 for limits on B($H \to \gamma\gamma$)$\cdot\sigma(e^+e^- \to Hf\overline{f})/\sigma(e^+e^- \to Hf\overline{f})$ (SM).

[23] BARATE 00L search for associated production of a $\gamma\gamma$ resonance with a $q\overline{q}$, $\nu\overline{\nu}$, or $\ell^+\ell^-$ pair in e^+e^- collisions at E_{cm}= 88–202 GeV. The limit is for a H^0 with SM production cross section. For B($H^0 \to \gamma\gamma$)=1, $m_{H^0} > 109$ GeV is obtained. See their Fig. 3 for limits on B($H \to \gamma\gamma$)$\cdot\sigma(e^+e^- \to Hf\overline{f})/\sigma(e^+e^- \to Hf\overline{f})$ (SM).

[24] ABBIENDI 99O search for associated production of a $\gamma\gamma$ resonance with a $q\overline{q}$, $\nu\overline{\nu}$, or $\ell^+\ell^-$ pair in e^+e^- collisions at 189 GeV. The limit is for a H^0 with SM production cross section. See their Fig. 4 for limits on $\sigma(e^+e^- \to H^0Z^0)\times$B($H^0 \to \gamma\gamma$)$\times$B($X^0 \to f\overline{f}$) for various masses. Updates the results of ACKERSTAFF 98Y.

[25] ABBOTT 99B search for associated production of a $\gamma\gamma$ resonance and a dijet pair. The limit assumes Standard Model values for the production cross section and for the couplings of the H^0 to W and Z bosons. Limits in the range of $\sigma(H^0+Z/W)\cdot$B($H^0 \to \gamma\gamma$)= 0.80–0.34 pb are obtained in the mass range m_{H^0}= 65–150 GeV.

[26] ABREU 99P search for $e^+e^- \to H^0\gamma$ with $H^0 \to b\overline{b}$ or $\gamma\gamma$, and $e^+e^- \to H^0q\overline{q}$ with $H^0 \to \gamma\gamma$. See their Fig. 4 for limits on $\sigma\times$B. Explicit limits within an effective interaction framework are also given.

Mass Limits for H^0 Decaying to Invisible Final States

These limits are for a neutral scalar H^0 which predominantly decays to invisible final states. Standard Model values are assumed for the couplings of H^0 to ordinary particles unless otherwise stated.

VALUE (GeV)	CL%	DOCUMENT ID		TECN	COMMENT
• • • We do not use the following data for averages, fits, limits, etc. • • •					
		[1] AABOUD	19AI	ATLS	WW/ZZ fusion
		[2] AAD	15BD	ATLS	$pp \to H^0WX$, H^0ZX
		[3] AAD	15BH	ATLS	jet + missing E_T
		[4] AAD	14BA	ATLS	secondary vertex
		[5] AAD	14O	ATLS	$pp \to H^0ZX$
		[6] CHATRCHYAN	14B	CMS	$pp \to H^0ZX$, qqH^0X
		[7] AAD	13AG	ATLS	secondary vertex
		[8] AAD	13AT	ATLS	electron jets

		[9] CHATRCHYAN	13BJ	CMS	
		[10] AAD	12AQ	ATLS	secondary vertex
		[11] AALTONEN	12AB	CDF	secondary vertex
		[12] AALTONEN	12U	CDF	secondary vertex
>108.2	95	[13] ABBIENDI	10	OPAL	
		[14] ABBIENDI	07	OPAL	large width
>112.3	95	[15] ACHARD	05	L3	
>112.1	95	[15] ABDALLAH	04B	DLPH	
>114.1	95	[15] HEISTER	02	ALEP	$E_{cm} \leq 209$ GeV
>106.4	95	[15] BARATE	01C	ALEP	$E_{cm} \leq 202$ GeV
> 89.2	95	[16] ACCIARRI	00M	L3	

[1] AABOUD 19AI search for $H^0_{1,2}$ production by vector boson fusion and decay to invisible final states in 36.1 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 6(b) for limits on cross section times branching ratios for $m_{H^0_{1,2}}$ = 0.1–3 TeV.

[2] AAD 15BD search for $pp \rightarrow H^0 W X$ and $pp \rightarrow H^0 Z X$ with W or Z decaying hadronically and H^0 decaying to invisible final states in 20.3 fb^{-1} at E_{cm} = 8TeV. See their Fig. 6 for a limit on the cross section times branching ratio for m_{H^0} = 115–300 GeV.

[3] AAD 15BH search for events with a jet and missing E_T in 20.3 fb^{-1} of pp collisions at E_{cm} = 8 TeV. Limits on $\sigma(H'^0)$ B($H'^0 \rightarrow$ invisible) < (44–10) pb (95%CL) is given for $m_{H'^0}$ = 115–300 GeV.

[4] AAD 14BA search for H^0 production in the decay mode $H^0 \rightarrow X^0 X^0$, where X^0 is a long-lived particle which decays to collimated pairs of $e^+ e^-$, $\mu^+ \mu^-$, or $\pi^+ \pi^-$ plus invisible particles, in 20.3 fb^{-1} of pp collisions at E_{cm} = 8 TeV. See their Figs. 15 and 16 for limits on cross section times branching ratio.

[5] AAD 14O search for $pp \rightarrow H^0 Z X$, $Z \rightarrow \ell\ell$, with H^0 decaying to invisible final states in 4.5 fb^{-1} at E_{cm} = 7 TeV and 20.3 fb^{-1} at E_{cm} = 8 TeV. See their Fig. 3 for a limit on the cross section times branching ratio for m_{H^0} = 110–400 GeV.

[6] CHATRCHYAN 14B search for $pp \rightarrow H^0 Z X$, $Z \rightarrow \ell\ell$ and $Z \rightarrow b\overline{b}$, and also $pp \rightarrow q q H^0 X$ with H^0 decaying to invisible final states using data at E_{cm} = 7 and 8 TeV. See their Figs. 10, 11 for limits on the cross section times branching ratio for m_{H^0} = 100–400 GeV.

[7] AAD 13AG search for H^0 production in the decay mode $H^0 \rightarrow X^0 X^0$, where X^0 is a long-lived particle which decays to $\mu^+ \mu^- X'^0$, in 1.9 fb^{-1} of pp collisions at E_{cm} = 7 TeV. See their Fig. 7 for limits on cross section times branching ratio.

[8] AAD 13AT search for H^0 production in the decay $H^0 \rightarrow X^0 X^0$, where X^0 eventually decays to clusters of collimated $e^+ e^-$ pairs, in 2.04 fb^{-1} of pp collisions at E_{cm} = 7 TeV. See their Fig. 3 for limits on cross section times branching ratio.

[9] CHATRCHYAN 13BJ search for H^0 production in the decay chain $H^0 \rightarrow X^0 X^0$, $X^0 \rightarrow \mu^+ \mu^- X'^0$ in 5.3 fb^{-1} of pp collisions at E_{cm} = 7 TeV. See their Fig. 2 for limits on cross section times branching ratio.

[10] AAD 12AQ search for H^0 production in the decay mode $H^0 \rightarrow X^0 X^0$, where X^0 is a long-lived particle which decays mainly to $b\overline{b}$ in the muon detector, in 1.94 fb^{-1} of pp collisions at E_{cm} = 7 TeV. See their Fig. 3 for limits on cross section times branching ratio for m_{H^0} = 120, 140 GeV, m_{X^0} = 20, 40 GeV in the $c\tau$ range of 0.5–35 m.

[11] AALTONEN 12AB search for H^0 production in the decay $H^0 \rightarrow X^0 X^0$, where X^0 eventually decays to clusters of collimated $\ell^+ \ell^-$ pairs, in 5.1 fb^{-1} of $p\overline{p}$ collisions at E_{cm} = 1.96 TeV. Cross section limits are provided for a benchmark MSSM model incorporating the parameters given in Table VI.

[12] AALTONEN 12U search for H^0 production in the decay mode $H^0 \rightarrow X^0 X^0$, where X^0 is a long-lived particle with $c\tau \approx 1$ cm which decays mainly to $b\overline{b}$, in 3.2 fb^{-1} of $p\overline{p}$ collisions at E_{cm} = 1.96 TeV. See their Figs. 9 and 10 for limits on cross section times branching ratio for m_{H^0} = (130–170) GeV, m_{X^0} = 20, 40 GeV.

[13] ABBIENDI 10 search for $e^+ e^- \rightarrow H^0 Z$ with H^0 decaying invisibly. The limit assumes SM production cross section and B($H^0 \rightarrow$ invisible) = 1.

[14] ABBIENDI 07 search for $e^+ e^- \rightarrow H^0 Z$ with $Z \rightarrow q\overline{q}$ and H^0 decaying to invisible final states. The H^0 width is varied between 1 GeV and 3 TeV. A limit $\sigma \cdot$ B($H^0 \rightarrow$ invisible) < (0.07–0.57) pb (95%CL) is obtained at E_{cm} = 206 GeV for m_{H^0} = 60–114 GeV.

[15] Search for $e^+ e^- \rightarrow H^0 Z$ with H^0 decaying invisibly. The limit assumes SM production cross section and B($H^0 \rightarrow$ invisible) = 1.

[16] ACCIARRI 00M search for $e^+ e^- \rightarrow Z H^0$ with H^0 decaying invisibly at E_{cm}=183–189 GeV. The limit assumes SM production cross section and B($H^0 \rightarrow$ invisible)=1. See their Fig. 6 for limits for smaller branching ratios.

Mass Limits for Light A^0

These limits are for a pseudoscalar A^0 in the mass range below $\mathcal{O}(10)$ GeV.

VALUE (GeV)	DOCUMENT ID		TECN	COMMENT
• • • We do not use the following data for averages, fits, limits, etc. • • •				
	[1] ADACHI	23A	BEL2	$\tau \rightarrow e A^0$, $\tau \rightarrow \mu A^0$
	[2] TUMASYAN	23AR	CMS	$H \rightarrow A^0 A^0 \rightarrow 4\gamma$
	[3] ABLIKIM	22H	BES3	$J/\psi \rightarrow A^0 \gamma$
	[4] JIA	22	BELL	$\Upsilon(1S) \rightarrow A^0 \gamma$
	[5] AAD	20AE	ATLS	$H^0 \rightarrow Z A^0$
	[6] AABOUD	18AP	ATLS	$H^0 \rightarrow A^0 A^0$
	[7] KHACHATRY...	17AZ	CMS	$H^0 \rightarrow A^0 A^0$
	[8] ABLIKIM	16E	BES3	$J/\psi \rightarrow A^0 \gamma$
	[9] KHACHATRY...	16F	CMS	$H^0 \rightarrow A^0 A^0$
	[10] LEES	15H	BABR	$\Upsilon(1S) \rightarrow A^0 \gamma$
	[11] LEES	13C	BABR	$\Upsilon(1S) \rightarrow A^0 \gamma$
	[12] LEES	13L	BABR	$\Upsilon(1S) \rightarrow A^0 \gamma$
	[13] LEES	13R	BABR	$\Upsilon(1S) \rightarrow A^0 \gamma$
	[14] ABLIKIM	12	BES3	$J/\psi \rightarrow A^0 \gamma$
	[15] CHATRCHYAN	12V	CMS	$A^0 \rightarrow \mu^+ \mu^-$
	[16] AALTONEN	11P	CDF	$t \rightarrow b H^+$, $H^+ \rightarrow W^+ A^0$
	[17,18] ABOUZAID	11A	KTEV	$K_L \rightarrow \pi^0 \pi^0 A^0$, $A^0 \rightarrow \mu^+ \mu^-$
	[19] DEL-AMO-SA...	11J	BABR	$\Upsilon(1S) \rightarrow A^0 \gamma$
	[20] LEES	11H	BABR	$\Upsilon(2S, 3S) \rightarrow A^0 \gamma$
	[21] ANDREAS	10	RVUE	
	[18,22] HYUN	10	BELL	$B^0 \rightarrow K^{*0} A^0$, $A^0 \rightarrow \mu^+ \mu^-$
	[18,23] HYUN	10	BELL	$B^0 \rightarrow \rho^0 A^0$, $A^0 \rightarrow \mu^+ \mu^-$
	[24] AUBERT	09P	BABR	$\Upsilon(3S) \rightarrow A^0 \gamma$
	[25] AUBERT	09Z	BABR	$\Upsilon(2S) \rightarrow A^0 \gamma$
	[26] AUBERT	09Z	BABR	$\Upsilon(3S) \rightarrow A^0 \gamma$
	[18,27] TUNG	09	K391	$K_L \rightarrow \pi^0 \pi^0 A^0$, $A^0 \rightarrow \gamma\gamma$
	[28] LOVE	08	CLEO	$\Upsilon(1S) \rightarrow A^0 \gamma$
	[29] BESSON	07	CLEO	$\Upsilon(1S) \rightarrow \eta_b \gamma$
	[30] PARK	05	HYCP	$\Sigma^+ \rightarrow p A^0$, $A^0 \rightarrow \mu^+ \mu^-$
	[31] BALEST	95	CLE2	$\Upsilon(1S) \rightarrow A^0 \gamma$
	[32] ANTREASYAN	90C	CBAL	$\Upsilon(1S) \rightarrow A^0 \gamma$

[1] ADACHI 23A search for flavor-changing τ decays $\tau \to e A^0$ and $\tau \to \mu A^0$, with A^0 invisible, using 62.8 fb^{-1} of $e^+ e^-$ collisions at E_{cm} = 10.58 GeV. Limits on B($\tau \to e A^0$)/B($\tau \to e \nu \nu$) in the range 1.1×10^{-3}–9.7×10^{-3} (95% CL) and B($\tau \to \mu A^0$)/B($\tau \to \mu \nu \nu$) in the range 0.7×10^{-3}–12.2×10^{-3} (95% CL) are given for m_{A^0} = 0–1.6 GeV. See their Fig. 2.

[2] TUMASYAN 23AR search for the decay $H \to A^0 A^0$ with $A^0 \to \gamma\gamma$ (detected as a merged photonlike object) using 136 fb^{-1} of pp collisions at E_{cm} = 13 TeV. Limits on B($H \to A^0 A^0$)$\cdot$B^2($A^0 \to \gamma\gamma$) in the range 0.9×10^{-3}–3.3×10^{-3} (95% CL) are given for m_{A^0} = 0.1–1.2 GeV. See their Fig. 2.

[3] ABLIKIM 22H search for the process $J/\psi \to A^0 \gamma$ with A^0 decaying to $\mu^+ \mu^-$ in 9×10^9 J/ψ events and give limits on B($J/\psi \to A^0 \gamma$)$\cdot$B($A^0 \to \mu^+ \mu^-$) in the range 1.2×10^{-9}–7.78×10^{-7} (90% CL) for 0.212 GeV $\leq m_{A^0} \leq$ 3.0 GeV. See their Fig. 4.

[4] JIA 22 search for the process $\Upsilon(2S) \to \Upsilon(1S)\pi^+\pi^- \to A^0\gamma\pi^+\pi^-$ with A^0 decaying to $\tau^+\tau^-$ or $\mu^+\mu^-$ in 158×10^6 $\Upsilon(2S)$ events and give limits on B($\Upsilon(1S) \to A^0\gamma$)$\cdot$B($A^0 \to \tau^+\tau^-$) in the range 3.8×10^{-6}–1.5×10^{-4} (90% CL) for m_{A^0} = 3.6–9.2 GeV, and B($\Upsilon(1S) \to A^0\gamma$)$\cdot$B($A^0 \to \mu^+\mu^-$) in the range 3.1×10^{-7}–1.6×10^{-5} (90% CL) for m_{A^0} = 0.21–9.2 GeV. See their Fig. 4.

[5] AAD 20AE search for the decay $H^0 \to Z A^0$, $Z \to \ell^+\ell^-$, A^0 decaying hadronically ($A^0 \to gg$ or $s\overline{s}$), in 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV. Limit on the product of production cross section and the $H^0 \to Z A^0$ branching ratio in the range 17–340 pb (95% CL) is given for m_{A^0} = 0.5–4.0 GeV, see their Table I.

[6] AABOUD 18AP search for the decay $H^0 \to A^0 A^0 \to \mu^+\mu^-\mu^+\mu^-$ in 36.1 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 10(b) for limits on B($H^0 \to A^0 A^0$) in the range m_{A^0} = 1–2.5, 4.5–8 GeV, assuming a type-II two-doublet plus singlet model with tan(β) = 5.

[7] KHACHATRYAN 17AZ search for the decay $H^0 \to A^0 A^0 \to \tau^+\tau^-\tau^+\tau^-$, $\mu^+\mu^- b\overline{b}$, and $\mu^+\mu^-\tau^+\tau^-$ in 19.7 fb^{-1} of pp collisions at E_{cm} = 8 TeV. See their Figs. 4, 5, and 6 for cross section limits in the range m_{A^0} = 5–62.5 GeV. See also their Figs. 7, 8, and 9 for interpretation of the data in terms of models with two Higgs doublets and a singlet.

[8] ABLIKIM 16E search for the process $J/\psi \to A^0\gamma$ with A^0 decaying to $\mu^+\mu^-$ and give limits on B($J/\psi \to A^0\gamma$)$\cdot$B($A^0 \to \mu^+\mu^-$) in the range 2.8×10^{-8}–5.0×10^{-6} (90% CL) for 0.212 $\leq m_{A^0} \leq$ 3.0 GeV. See their Fig. 5.

[9] KHACHATRYAN 16F search for the decay $H^0 \to A^0 A^0 \to \tau^+\tau^-\tau^+\tau^-$ in 19.7 fb^{-1} of pp collisions at E_{cm} = 8 TeV. See their Fig. 8 for cross section limits for m_{A^0} = 4–8 GeV.

[10] LEES 15H search for the process $\Upsilon(2S) \to \Upsilon(1S)\pi^+\pi^- \to A^0\gamma\pi^+\pi^-$ with A^0 decaying to $c\overline{c}$ and give limits on B($\Upsilon(1S) \to A^0\gamma$)$\cdot$B($A^0 \to c\overline{c}$) in the range 7.4×10^{-5}–2.4×10^{-3} (90% CL) for 4.00 $\leq m_{A^0} \leq$ 8.95 and 9.10 $\leq m_{A^0} \leq$ 9.25 GeV. See their Fig. 6.

[11] LEES 13C search for the process Υ(2S, 3S)$\to \Upsilon(1S)\pi^+\pi^- \to A^0\gamma\pi^+\pi^-$ with A^0 decaying to $\mu^+\mu^-$ and give limits on B($\Upsilon(1S) \to A^0\gamma$)$\cdot$B($A^0 \to \mu^+\mu^-$) in the range (0.3–9.7) $\times 10^{-6}$ (90% CL) for 0.212 $\leq m_{A^0} \leq$ 9.20 GeV. See their Fig. 5(e) for limits on the $b-A^0$ Yukawa coupling derived by combining this result with AUBERT 09Z.

[12] LEES 13L search for the process $\Upsilon(2S) \to \Upsilon(1S)\pi^+\pi^- \to A^0\gamma\pi^+\pi^-$ with A^0 decaying to gg or $s\overline{s}$ and give limits on B($\Upsilon(1S) \to A^0\gamma$)$\cdot$B($A^0 \to gg$) between 1×10^{-6} and 2×10^{-2} (90% CL) for 0.5 $\leq m_{A^0} \leq$ 9.0 GeV, and B($\Upsilon(1S) \to$

$A^0\gamma$)·B($A^0 \to s\overline{s}$) between 4×10^{-6} and 1×10^{-3} (90%CL) for 1.5 $\leq m_{A^0} \leq$ 9.0 GeV. See their Fig. 4.

[13] LEES 13R search for the process $\Upsilon(2S) \to \Upsilon(1S)\pi^+\pi^- \to A^0\gamma\pi^+\pi^-$ with A^0 decaying to $\tau^+\tau^-$ and give limits on B($\Upsilon(1S) \to A^0\gamma$)·B($A^0 \to \tau^+\tau^-$) in the range 0.9–13 $\times 10^{-5}$ (90% CL) for 3.6 $\leq m_{A^0} \leq$ 9.2 GeV. See their Fig. 4 for limits on the $b-A^0$ Yukawa coupling derived by combining this result with AUBERT 09P.

[14] ABLIKIM 12 searches for the process $\psi(3686) \to \pi\pi J/\psi$, $J/\psi \to A^0\gamma$ with A^0 decaying to $\mu^+\mu^-$. It gives mass dependent limits on B($J/\psi \to A^0\gamma$)·B($A^0 \to \mu^+\mu^-$) in the range 4×10^{-7}–2.1×10^{-5} (90% C.L.) for 0.212 $\leq m_{A^0} \leq$ 3.0 GeV. See their Fig. 2.

[15] CHATRCHYAN 12V search for A^0 production in the decay $A^0 \to \mu^+\mu^-$ with 1.3 fb^{-1} of pp collisions at E_{cm} = 7 TeV. A limit on $\sigma(A^0)$·B($A^0 \to \mu^+\mu^-$) in the range (1.5–7.5) pb is given for m_{A^0} = (5.5–8.7) and (11.5–14) GeV at 95% CL.

[16] AALTONEN 11P search in 2.7 fb^{-1} of $p\overline{p}$ collisions at E_{cm} = 1.96 TeV for the decay chain $t \to bH^+$, $H^+ \to W^+A^0$, $A^0 \to \tau^+\tau^-$ with m_{A^0} between 4 and 9 GeV. See their Fig. 4 for limits on B($t \to bH^+$) for 90 $< m_{H^+} <$ 160 GeV.

[17] ABOUZAID 11A search for the decay chain $K_L \to \pi^0\pi^0A^0$, $A^0 \to \mu^+\mu^-$ and give a limit B($K_L \to \pi^0\pi^0A^0$) · B($A^0 \to \mu^+\mu^-$) $<$ 1.0 $\times 10^{-10}$ at 90% CL for m_{A^0} = 214.3 MeV.

[18] The search was motivated by PARK 05.

[19] DEL-AMO-SANCHEZ 11J search for the process $\Upsilon(2S) \to \Upsilon(1S)\pi^+\pi^- \to A^0\gamma\pi^+\pi^-$ with A^0 decaying to invisible final states. They give limits on B($\Upsilon(1S) \to A^0\gamma$)·B($A^0 \to$ invisible) in the range (1.9–4.5) $\times 10^{-6}$ (90% CL) for 0 $\leq m_{A^0} \leq$ 8.0 GeV, and (2.7–37) $\times 10^{-6}$ for 8.0 $\leq m_{A^0} \leq$ 9.2 GeV.

[20] LEES 11H search for the process Υ(2S, 3S) $\to A^0\gamma$ with A^0 decaying hadronically and give limits on B(Υ(2S, 3S) $\to A^0\gamma$)·B($A^0 \to$ hadrons) in the range 1×10^{-6}–8×10^{-5} (90% CL) for 0.3 $< m_{A^0} <$ 7 GeV. The decay rates for $\Upsilon(2S)$ and $\Upsilon(3S)$ are assumed to be equal up to the phase space factor. See their Fig. 5.

[21] ANDREAS 10 analyze constraints from rare decays and other processes on a light A^0 with $m_{A^0} < 2m_\mu$ and give limits on its coupling to fermions at the level of 10^{-4} times the Standard Model value.

[22] HYUN 10 search for the decay chain $B^0 \to K^{*0}A^0$, $A^0 \to \mu^+\mu^-$ and give a limit on B($B^0 \to K^{*0}A^0$) · B($A^0 \to \mu^+\mu^-$) in the range (2.26–5.53) $\times 10^{-8}$ at 90%CL for m_{A^0} = 212–300 MeV. The limit for m_{A^0} = 214.3 MeV is 2.26 $\times 10^{-8}$.

[23] HYUN 10 search for the decay chain $B^0 \to \rho^0A^0$, $A^0 \to \mu^+\mu^-$ and give a limit on B($B^0 \to \rho^0A^0$) · B($A^0 \to \mu^+\mu^-$) in the range (1.73–4.51) $\times 10^{-8}$ at 90%CL for m_{A^0} = 212–300 MeV. The limit for m_{A^0} = 214.3 MeV is 1.73 $\times 10^{-8}$.

[24] AUBERT 09P search for the process $\Upsilon(3S) \to A^0\gamma$ with $A^0 \to \tau^+\tau^-$ for 4.03 $< m_{A^0} <$ 9.52 and 9.61 $< m_{A^0} <$ 10.10 GeV, and give limits on B($\Upsilon(3S) \to A^0\gamma$)·B($A^0 \to \tau^+\tau^-$) in the range (1.5–16) $\times 10^{-5}$ (90% CL).

[25] AUBERT 09Z search for the process $\Upsilon(2S) \to A^0\gamma$ with $A^0 \to \mu^+\mu^-$ for 0.212 $< m_{A^0} <$ 9.3 GeV and give limits on B($\Upsilon(2S) \to A^0\gamma$)·B($A^0 \to \mu^+\mu^-$) in the range (0.3–8) $\times 10^{-6}$ (90% CL).

[26] AUBERT 09Z search for the process $\Upsilon(3S) \to A^0\gamma$ with $A^0 \to \mu^+\mu^-$ for 0.212 $< m_{A^0} <$ 9.3 GeV and give limits on B($\Upsilon(3S) \to A^0\gamma$)·B($A^0 \to \mu^+\mu^-$) in the range (0.3–5) $\times 10^{-6}$ (90% CL).

[27] TUNG 09 search for the decay chain $K_L \to \pi^0 \pi^0 A^0$, $A^0 \to \gamma\gamma$ and give a limit on $B(K_L \to \pi^0 \pi^0 A^0) \cdot B(A^0 \to \gamma\gamma)$ in the range $(2.4–10.7) \times 10^{-7}$ at 90%CL for m_{A^0} = 194.3–219.3 MeV. The limit for m_{A^0} = 214.3 MeV is 2.4×10^{-7}.

[28] LOVE 08 search for the process $\Upsilon(1S) \to A^0 \gamma$ with $A^0 \to \mu^+\mu^-$ (for $m_{A^0} < 2m_\tau$) and $A^0 \to \tau^+\tau^-$. Limits on $B(\Upsilon(1S) \to A^0 \gamma) \cdot B(A^0 \to \ell^+\ell^-)$ in the range 10^{-6}–10^{-4} (90% CL) are given.

[29] BESSON 07 give a limit $B(\Upsilon(1S) \to \eta_b \gamma) \cdot B(\eta_b \to \tau^+\tau^-) < 0.27\%$ (95% CL), which constrains a possible A^0 exchange contribution to the η_b decay.

[30] PARK 05 found three candidate events for $\Sigma^+ \to p \mu^+\mu^-$ in the HyperCP experiment. Due to a narrow spread in dimuon mass, they hypothesize the events as a possible signal of a new boson. It can be interpreted as a neutral particle with $m_{A^0} = 214.3 \pm 0.5$ MeV and the branching fraction $B(\Sigma^+ \to p A^0) \cdot B(A^0 \to \mu^+\mu^-) = (3.1^{+2.4}_{-1.9} \pm 1.5) \times 10^{-8}$.

[31] BALEST 95 give limits $B(\Upsilon(1S) \to A^0 \gamma)$ ¡ 1.5×10^{-5} at 90% CL for $m_{A^0} < 5$ GeV. The limit becomes $< 10^{-4}$ for $m_{A^0} < 7.7$ GeV.

[32] ANTREASYAN 90C give limits $B(\Upsilon(1S) \to A^0 \gamma)$ ¡ 5.6×10^{-5} at 90% CL for $m_{A^0} <$ 7.2 GeV. A^0 is assumed not to decay in the detector.

Other Mass Limits

We use a symbol H_1^0 if mass < 125 GeV or H_2^0 if mass > 125 GeV. The notation H is reserved for the 125 GeV particle.

VALUE (GeV)	CL%	DOCUMENT ID		TECN	COMMENT
• • • We do not use the following data for averages, fits, limits, etc. • • •					
		[1] HAYRAPETY...	25A	CMS	$H_1^0 \to \gamma\gamma$
		[2] AAD	24A	ATLS	$H_2^0 \to Z\gamma$
		[3] AAD	24AP	ATLS	H_2^0, $A^0 \to t\bar{t}$
		[4] AAD	24BQ	ATLS	$H \to A^0 A^0$
		[5] AAD	24BV	ATLS	$H_2^0 \to HH$
		[6] AAD	24BW	ATLS	$H_3^0 \to H_2^0 H$
		[7] AAD	24BX	ATLS	$H_4^0 \to H_2^0 H_3^0$
		[8] AAD	24BX	ATLS	$A^0 \to H_2^0 Z$
		[9] AAD	24CA	ATLS	$H_3^0 \to H_{1,2}^0 H$
		[10] AAD	24CD	ATLS	$H_2^0 \to ZA$
		[11] AAD	24CF	ATLS	$A^0 \to \tau^+\tau^-$
		[12] AAD	24CI	ATLS	H_2^0, $A^0 \to$ invisible
		[13] AAD	24CK	ATLS	two doublet + pseudoscalar + Dirac DM
		[14] AAD	24H	ATLS	$H_2^0 \to HH$
		[15] AAD	24K	ATLS	$H \to Z A^0$
		[16] ADACHI	24F	BEL2	$H_1^0 \to \mu^+\mu^-$
		[17] HAYRAPETY...	24AA	CMS	$H \to A^0 A^0$
		[18] HAYRAPETY...	24AE	CMS	$H_2^0 \to HH$
		[19] HAYRAPETY...	24AF	CMS	$H \to A^0 A^0$
		[20] HAYRAPETY...	24AI	CMS	$H \to H_1^0 H_1^0$, long-lived
		[21] HAYRAPETY...	24AJ	CMS	$H_2^0 \to \gamma\gamma$
		[22] HAYRAPETY...	24AZ	CMS	$A^0 A^0$, $A^0 \to \mu^+\mu^-$

[23] HAYRAPETY...	24AZ	CMS	H, $H^0_{1,2} \to A^0 A^0$, $A^0 \to \mu^+ \mu^-$
[24] HAYRAPETY...	24I	CMS	$H \to Z A^0$
[25] TUMASYAN	24A	CMS	$H^0_{1,2} \to e^+ e^-$, $\mu^+ \mu^-$, $\tau^+ \tau^-$
[26] TUMASYAN	24B	CMS	$H^0_2 \to H H$
[27] TUMASYAN	24B	CMS	$H^0_3 \to H^0_{1,2} H$
[28] AAD	23AD	ATLS	$H^0_2 \to H H$
[29] AAD	23AD	ATLS	$A^0 \to Z H^0_2 \to Z H H$
[30] AAD	23AJ	ATLS	$H^\pm \to W^\pm A^0$
[31] AAD	23BD	ATLS	$t \to q H^0_{1,2}$
[32] AAD	23BE	ATLS	$H^0_2 \to W^+ W^-$
[33] AAD	23BG	ATLS	$t \overline{t} H^0_2 / A^0$
[34] AAD	23BW	ATLS	$A^0 t \overline{t}$, $A^0 \to \mu^+ \mu^-$
[35] AAD	23BX	ATLS	H + invisible A^0
[36] AAD	23CA	ATLS	$H^0_3 \to H^0_2 H$
[37] AAD	23CR	ATLS	flavor changing H^0_2
[38] AAD	23O	ATLS	$A^0 \to Z H$
[39] AAD	23R	ATLS	$A^0 \to \gamma \gamma$
[40] AAD	23U	ATLS	$H^0_2 \to Z \gamma$
[41] AAD	23Z	ATLS	$H^0_2 \to H H$
[42] HAYRAPETY...	23C	CMS	$H^0_{1,2} \to e \mu$
[43] HAYRAPETY...	23G	CMS	$A^0 \to \mu^+ \mu^-$
[44] TUMASYAN	23	CMS	$H^0_3 \to H^0_{1,2} H$
[45] TUMASYAN	23M	CMS	$H \to A^0 A^0$
[46] TUMASYAN	23O	CMS	$H^0_2 \to H H$
[47] TUMASYAN	23S	CMS	$H^0_{1,2} \to \tau^+ \tau^-$
[48] AAD	22A	ATLS	$H \to A^0 A^0$
[49] AAD	22D	ATLS	$Z A^0$, $A^0 \to$ invisible
[50] AAD	22F	ATLS	$H^0_2 \to H H$
[51] AAD	22I	ATLS	$H \to \widetilde{\chi}^0_2 \widetilde{\chi}^0_1$, $\widetilde{\chi}^0_2 \to A^0 \widetilde{\chi}^0_1$, $A^0 \to b \overline{b}$
[52] AAD	22J	ATLS	$H \to Z A^0$
[53] AAD	22J	ATLS	$H \to A^0 A^0$, $H^0_1 H^0_1$
[54] AAD	22P	ATLS	H^0_1, $H^0_2 \to$ invisible
[55] AAD	22Y	ATLS	$H^0_2 \to H H$
[56] ABRATENKO	22A	MCBN	$K^+ \to H^0_1 \pi^+$
[57] TUMASYAN	22AK	CMS	$H^0_3 \to H^0_1 H^0_1$
[58] TUMASYAN	22D	CMS	$H^0_2 \to W^+ W^-$
[59] AAD	21AF	ATLS	$H^0_2 \to Z Z$
[60] AAD	21AI	ATLS	$A^0 \to Z H^0_2$
[61] AAD	21AY	ATLS	$H^0_2 \to \gamma \gamma$
[62] AAD	21AZ	ATLS	$A^0_2 \to H A^0_1$

[63] AAD	21BB	ATLS	$A_2^0 \to H A_1^0$
[64] AAD	21BE	ATLS	$A_1^0 \to$ invisible
[65] ABRATENKO	21	MCBN	$K^+ \to H_1^0 \pi^+$
[66] SIRUNYAN	21A	CMS	$H_2^0 \to Z A^0$, $A^0 \to$ invisible
[67] TUMASYAN	21F	CMS	$H_3^0 \to H H_{1,2}^0$
[68] AAD	20AA	ATLS	$H_2^0/A^0 \to \tau^+ \tau^-$
[69] AAD	20AI	ATLS	$H \to A^0 A^0$
[70] AAD	20AO	ATLS	$H_2^0 \to H H$
[71] AAD	20C	ATLS	$H_2^0 \to H H$
[72] AAD	20L	ATLS	$H_2^0 \to b \overline{b}$
[73] AAD	20X	ATLS	$H_2^0 \to H H$
[74] AAIJ	20AL	LHCB	$A^0 \to \mu^+ \mu^-$
[75] SIRUNYAN	20	CMS	$H \to A^0 A^0$
[76] SIRUNYAN	20AA	CMS	$H_2^0 \to Z A^0$ or $A^0 \to Z H_2^0$
[77] SIRUNYAN	20AC	CMS	$A^0 \to Z H$
[78] SIRUNYAN	20AD	CMS	$H_2^0 \to \mu \tau,\ e \tau$
[79] SIRUNYAN	20AF	CMS	$H_2^0/A^0 \to t \overline{t}$
[80] SIRUNYAN	20AP	CMS	H, $H_2^0 \to A^0 A^0$
[81] SIRUNYAN	20Y	CMS	$H_2^0 \to W^+ W^-$
[82] SIRUNYAN	20Z	CMS	$t \overline{t} H_{1,2}^0$ or $t \overline{t} A^0$, $H_{1,2}^0/A^0 \to e^+ e^-, \mu^+ \mu^-$
[83] AABOUD	19A	ATLS	$H_2^0 \to H H$
[84] AABOUD	19AG	ATLS	$H \to A^0 A^0$
[85] AABOUD	19O	ATLS	$H_2^0 \to H H$
[86] AABOUD	19T	ATLS	$H_2^0 \to H H$
[87] AABOUD	19V	ATLS	two doublet + pseudoscalar model
[88] AABOUD	19Y	ATLS	$H_2^0 \to \mu^+ \mu^-$
[89] AALTONEN	19	CDF	$H_{1,2}^0 \to b \overline{b}$
[90] SIRUNYAN	19	CMS	$H_2^0 \to H H$
[91] SIRUNYAN	19AE	CMS	$A^0 \to \tau^+ \tau^-$
[92] SIRUNYAN	19AN	CMS	$A_2^0 \to H A_1^0$
[93] SIRUNYAN	19AV	CMS	$A^0 \to Z H$
[94] SIRUNYAN	19B	CMS	$H_{1,2}^0/A^0 \to b \overline{b}$
[95] SIRUNYAN	19BB	CMS	$H_1^0 \to \gamma \gamma$
[96] SIRUNYAN	19BD	CMS	$H \to A^0 A^0$
[97] SIRUNYAN	19BE	CMS	$H_2^0 \to H H$
[98] SIRUNYAN	19BQ	CMS	$H_{1,2}^0 \to A^0 A^0$
[99] SIRUNYAN	19CR	CMS	$H_2^0/A^0 \to \mu^+ \mu^-$
[100] SIRUNYAN	19H	CMS	$H_2^0 \to H H$
[101] AABOUD	18AA	ATLS	$H_2^0 \to Z \gamma$
[102] AABOUD	18AG	ATLS	$H \to A^0 A^0$
[103] AABOUD	18AH	ATLS	$A^0 \to Z H_2^0$
[104] AABOUD	18AI	ATLS	$A^0 \to Z H$

[105] AABOUD	18BF	ATLS	$H_2^0 \to ZZ$
[106] AABOUD	18BU	ATLS	$H_2^0 \to HH$
[107] AABOUD	18BX	ATLS	$H \to A^0 A^0$
[108] AABOUD	18CQ	ATLS	$H_2^0 \to HH$
[109] AABOUD	18F	ATLS	$H_2^0 \to W^+ W^-$, ZZ
[110] AAIJ	18AM	LHCB	$H_{1,2}^0 \to \mu\tau$
[111] AAIJ	18AQ	LHCB	$A^0 \to \mu^+\mu^-$
[112] AAIJ	18AQ	LHCB	$H \to A^0 A^0$, $A^0 \to \mu^+\mu^-$
[113] SIRUNYAN	18AF	CMS	$H_2^0 \to HH$
[114] SIRUNYAN	18BA	CMS	$H_2^0 \to ZZ$
[115] SIRUNYAN	18CW	CMS	$H_2^0 \to HH$
[116] SIRUNYAN	18DK	CMS	$H_2^0 \to Z\gamma$
[117] SIRUNYAN	18DT	CMS	$H \to A^0 A^0$
[118] SIRUNYAN	18DU	CMS	$H_2^0 \to \gamma\gamma$
[119] SIRUNYAN	18ED	CMS	$A^0 \to ZH$
[120] SIRUNYAN	18EE	CMS	$H \to A^0 A^0$
[121] SIRUNYAN	18F	CMS	pp, 13 TeV, $H_2^0 \to HH$
[122] AABOUD	17	ATLS	$H_2^0 \to Z\gamma$
[123] AABOUD	17AP	ATLS	$H_2^0 \to \gamma\gamma$
[124] AABOUD	17AW	ATLS	$H_2^0 \to Z\gamma$
[125] KHACHATRY...	17AZ	CMS	$H \to A^0 A^0$
[126] KHACHATRY...	17D	CMS	pp, 8, 13 TeV, $H_2^0 \to Z\gamma$
[127] KHACHATRY...	17R	CMS	$H_2^0 \to \gamma\gamma$
[128] SIRUNYAN	17CN	CMS	pp, 8 TeV, $H_2^0 \to HH$
[129] SIRUNYAN	17Y	CMS	pp, 8, 13 TeV, $H_2^0 \to Z\gamma$
[130] AABOUD	16AB	ATLS	$H \to A^0 A^0$
[131] AABOUD	16AE	ATLS	$H_2^0 \to W^+ W^-$, ZZ
[132] AABOUD	16H	ATLS	$H_2^0 \to \gamma\gamma$
[133] AABOUD	16I	ATLS	$H_2^0 \to HH$
[134] AAD	16AX	ATLS	$H \to ZZ$
[135] AAD	16C	ATLS	$H \to W^+ W^-$
[136] AAD	16L	ATLS	$H \to A^0 A^0$
[137] AAD	16L	ATLS	$H_2^0 \to A^0 A^0$
[138] AALTONEN	16C	CDF	$H_1^0 H^\pm \to H_1^0 H_1^0 W^*$, $H_1^0 \to \gamma\gamma$
[139] KHACHATRY...	16BG	CMS	$H_2^0 \to HH$
[140] KHACHATRY...	16BQ	CMS	pp, 8 TeV, $H_2^0 \to HH$
[141] KHACHATRY...	16F	CMS	$H \to H_1 H_1$
[142] KHACHATRY...	16M	CMS	$H_2^0 \to \gamma\gamma$
[143] KHACHATRY...	16P	CMS	$H_2^0 \to HH$
[144] KHACHATRY...	16P	CMS	$A^0 \to ZH$
[145] AAD	15BK	ATLS	$H_2^0 \to HH$
[146] AAD	15BZ	ATLS	$H \to A^0 A^0$
[147] AAD	15BZ	ATLS	$H_2^0 \to A^0 A^0$
[148] AAD	15CE	ATLS	$H_2^0 \to HH$

		[149] AAD	15H	ATLS	$H_2^0 \to HH$
		[150] AAD	15S	ATLS	$A^0 \to ZH$
		[151] KHACHATRY...	15AW	CMS	$H_2^0 \to W^+W^-$, ZZ
		[152] KHACHATRY...	15BB	CMS	$H \to \gamma\gamma$
		[153] KHACHATRY...	15N	CMS	$A^0 \to ZH$
		[154] KHACHATRY...	15O	CMS	$A^0 \to ZH$
		[155] KHACHATRY...	15R	CMS	$H_2^0 \to HH$
		[156] AAD	14AP	ATLS	$H \to \gamma\gamma$
		[157] AAD	14M	ATLS	$H_2^0 \to H^\pm W^\mp \to HW^\pm W^\mp$, $H \to b\overline{b}$
		[158] CHATRCHYAN	14G	CMS	$H \to WW^{(*)}$
		[159] KHACHATRY...	14P	CMS	$H \to \gamma\gamma$
		[160] AALTONEN	13P	CDF	$H'^0 \to H^\pm W^\mp \to HW^+W^-$
		[161] CHATRCHYAN	13BJ	CMS	$H \to A^0 A^0$
		[162] AALTONEN	11P	CDF	$t \to bH^+$, $H^+ \to W^+ A^0$
		[163] ABBIENDI	10	OPAL	$H \to \widetilde{\chi}_1^0 \widetilde{\chi}_2^0$
		[164] SCHAEL	10	ALEP	$H \to A^0 A^0$
		[165] ABAZOV	09V	D0	$H \to A^0 A^0$
none 3–63	95	[166] ABBIENDI	05A	OPAL	A^0, Type II model
$>$104	95	[167] ABBIENDI	04K	OPAL	$H \to$ 2 jets
		[168] ABDALLAH	04	DLPH	HVV couplings
$>$110.3	95	[169] ACHARD	04B	L3	$H \to$ 2 jets
		[170] ACHARD	04F	L3	Anomalous coupling
		[171] ABBIENDI	03F	OPAL	$e^+e^- \to HZ$, $H \to$ any
		[172] ABBIENDI	03G	OPAL	$H_1^0 \to A^0 A^0$
$>$105.4	95	[173,174] HEISTER	02L	ALEP	$H_1^0 \to \gamma\gamma$
$>$109.1	95	[175] HEISTER	02M	ALEP	$H \to$ 2 jets or $\tau^+\tau^-$
none 12–56	95	[176] ABBIENDI	01E	OPAL	A^0, Type-II model
		[177] ACCIARRI	00R	L3	$e^+e^- \to H\gamma$ and/or $H \to \gamma\gamma$
		[178] ACCIARRI	00R	L3	$e^+e^- \to e^+e^- H$
		[179] GONZALEZ...	98B	RVUE	Anomalous coupling
		[180] KRAWCZYK	97	RVUE	$(g-2)_\mu$
		[181] ALEXANDER	96H	OPAL	$Z \to H\gamma$

[1] HAYRAPETYAN 25A search for the decay $H_1^0 \to \gamma\gamma$ in 132 fb^{-1} of pp collisions at $E_{\rm cm}$ = 13 TeV. See their Fig. 5 for limits on cross section times branching ratio for $m_{H_1^0}$ = 70–110 GeV.

[2] AAD 24A search for the decay $H_2^0 \to Z\gamma$ with Z decaying to e^+e^- or $\mu^+\mu^-$ using 140 fb^{-1} of pp collisions at $E_{\rm cm}$ = 13 TeV. See their Fig. 4 for limits on production cross section times branching ratios for $m_{H_2^0}$ = 0.22–3.4 TeV.

[3] AAD 24AP search for production of a heavy H_2^0 or A^0 decaying to $t\overline{t}$ in 140 fb^{-1} of pp collisions at $E_{\rm cm}$ = 13 TeV. See their Figs. 15 and 16 for limits on the Yukawa coupling for m_{A^0}, $m_{H_2^0}$ = 0.4–1.4 TeV for different assumptions on width over mass.

[4] AAD 24BQ search for the decay chain $H \to A^0 A^0 \to b\overline{b}\tau^+\tau^-$, produced by gluon fusion, vector boson fusion, or in association with a weak vector boson, in 140 fb^{-1} of pp collisions at $E_{\rm cm}$ = 13 TeV. See their Fig. 10 for limits on the product of branching ratios in the range m_{A^0} = 12–60 GeV.

[5] AAD 24BV search for H^0_2 produced by vector-boson fusion, decaying to $HH \to b\overline{b}b\overline{b}$, using 140 fb^{-1} of pp collisions at $E_{\rm cm}$ = 13 TeV. See their Fig. 8 for limits on the product of production cross section times branching ratios for $m_{H^0_2}$ = 1–5 TeV.

[6] AAD 24BW search for production of H^0_3 decaying to $H^0_2 H$, $H^0_2 \to W^+W^-$ or ZZ, and $H \to \gamma\gamma$ using 140 fb^{-1} of pp collisions at $E_{\rm cm}$ = 13 TeV. See their Figs. 6–8 for limits on production cross section times branching ratios in the ranges $m_{H^0_3}$ = 0.3–1.0 TeV and $m_{H^0_2}$ = 0.17–0.5 TeV.

[7] AAD 24BX search for production of H^0_4 decaying to $H^0_2 H^0_3$, $H^0_2 \to$ invisible, and $H^0_3 \to ZZ \to 4\ell$, using 139 fb^{-1} of pp collisions at $E_{\rm cm}$ = 13 TeV. See their Fig. 8 (a), (b) for limits on production cross section times branching ratios in the ranges $m_{H^0_4}$ = 0.39–1.3 TeV, $m_{H^0_3}$ = 0.22–1.0 TeV, and $m_{H^0_2}$ = 160 GeV.

[8] AAD 24BX search for production of A^0 decaying to $H^0_2 Z$, $H^0_2 \to ZZ$, resulting in 4ℓ final states, using 139 fb^{-1} of pp collisions at $E_{\rm cm}$ = 13 TeV. See their Fig. 8 (c), (d) for limits on production cross section times branching ratios in the ranges m_{A^0} = 0.32–1.3 TeV and $m_{H^0_2}$ = 0.22–1.0 TeV.

[9] AAD 24CA search for production of H^0_3 decaying to $H^0_{1,2} H$, $H^0_{1,2} \to b\overline{b}$, and $H \to \gamma\gamma$ using 140 fb^{-1} of pp collisions at $E_{\rm cm}$ = 13 TeV. See their Fig. 4 for limits on production cross section times branching ratios in the ranges $m_{H^0_3}$ = 0.17–1.0 TeV and $m_{H^0_{1,2}}$ = 0.015–0.5 TeV.

[10] AAD 24CD search for production of H^0_2 decaying to ZA, $A \to$ invisible, using 140 fb^{-1} of pp collisions at $E_{\rm cm}$ = 13 TeV. See their Figs. 8–12 for excluded parameter regions in two-Higgs-doublet plus singlet pseudoscalar model.

[11] AAD 24CF search for gluon-fusion production of A^0 decaying to $\tau^+\tau^-$ in 140 fb^{-1} of pp collisions at $E_{\rm cm}$ = 13 TeV. See their Fig. 7 for limits on cross section times branching ratio in the range m_{A^0} = 20–90 GeV, and Fig. 8 for A^0 coupling with the top quark.

[12] AAD 24CI search for production of H^0_2, A^0 decaying to invisible final states, in various channels including $t\overline{t}$ + missing E_T and $t\overline{t}t\overline{t}$, using up to 140 fb^{-1} of pp collisions at $E_{\rm cm}$ = 13 TeV. See their Fig. 3 for cross section limits in a simplified model with fixed Dark Matter mass and couplings to quarks and Dark Matter.

[13] AAD 24CK combines published ATLAS data up to 139 fb^{-1} to constrain two-Higgs-doublet plus singlet pseudoscalar model with A^0_1 decaying to invisibly into Dirac Dark Matter final states. See their Figs. 4–9 for excluded parameter regions.

[14] AAD 24H combine searches for a scalar resonance decaying to HH using up to 139 fb^{-1} of pp collisions at $E_{\rm cm}$ = 13 TeV from AAD 22F, AAD 23Z, and AAD 22Y. See their Fig. 1 for limits on cross section times branching ratio for $m_{H^0_2}$ = 0.251–5.0 TeV.

[15] AAD 24K search for the decay $H \to ZA^0$ with $A^0 \to \gamma\gamma$ and $Z \to \ell^+\ell^-$ in 139 fb^{-1} of pp collisions at $E_{\rm cm}$ = 13 TeV. Merged $\gamma\gamma$ events as well as resolved ones are looked for. See their Fig. 3 for limits on the product of the branching ratios for m_{A^0} = 0.1–33 GeV.

[16] ADACHI 24F search for H^0_1, which couples preferentially to muon pairs, in the process $e^+ e^- \to \mu^+ \mu^- H^0_1 \to \mu^+ \mu^- \mu^+ \mu^-$ using 178 fb^{-1} of $e^+ e^-$ collisions at E_{cm} = 10.58 GeV. See their Fig. 13 (lower) for limits on cross section times branching ratio for $m_{H^0_1}$ = 0.212–9.0 GeV. Limits on the model parameter are given in Fig. 14 (lower).

[17] HAYRAPETYAN 24AA search for the decay chain $H \to A^0 A^0 \to b\overline{b}b\overline{b}$, produced in association with $W^\pm$ or Z, in 138 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 3 for limits on the product of branching ratios in the range m_{A^0} = 15–60 GeV.

[18] HAYRAPETYAN 24AE search for $H^0_2 \to HH$ in the final state $b\overline{b}W^+ W^-$ using 138 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 19 (upper) for limit on the product of production cross section times branching ratios for $m_{H^0_2}$ = 0.25–0.9 TeV.

[19] HAYRAPETYAN 24AF search for $H \to A^0 A^0$ in the final state $b\overline{b}\mu^+ \mu^-$ and $b\overline{b}\tau^+ \tau^-$ in 138 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Figs. 10, 11, and 12 for limits on the product of branching ratios in the range m_{A^0} = 12–60 GeV. Interpretation of the limits in terms of two Higgs doublet plus singlet models is given in their Figs. 13 and 14.

[20] HAYRAPETYAN 24AI search for $H \to H^0_1 H^0_1$, with long-lived H^0_1 decaying in the muon detectors to $\gamma\gamma$, $e^+ e^-$, $\pi^0 \pi^0$, $\pi^+ \pi^-$, $K^+ K^-$, $K^0 K^0$, $d\overline{d}$, $\tau^+ \tau^-$, or $b\overline{b}$, using 138 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Figs. 13–17 for limits on B($H \to H^0_1 H^0_1$) in the mass range between 0.4 and 55 GeV for each of the decay modes.

[21] HAYRAPETYAN 24AJ search for production of a scalar resonance decaying to $\gamma\gamma$ in 138 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 3 (right) for limits on production cross section times branching ratio for $m_{H^0_2}$ = 0.6–5.0 TeV, and Fig. 4 (lower) for their dependence on the width.

[22] HAYRAPETYAN 24AZ search for pair production of A^0 decaying to $\mu^+ \mu^-$ using 101–137 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 5 for limits on cross section times branching ratio for m_{A^0} = 0.21–60 GeV.

[23] HAYRAPETYAN 24AZ search for H or $H^0_{1,2}$ decaying to $A^0 A^0 \to \mu^+ \mu^- \mu^+ \mu^-$ using 101–137 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 8 for limits on cross section times branching ratio for m_{A^0} = 0.5–2.7 GeV and several choices of $m_{H^0_{1,2}}$ in a NMSSM scenario.

[24] HAYRAPETYAN 24I search for the decay $H \to Z A^0$ with $A^0 \to \gamma\gamma$ and $Z \to \ell^+ \ell^-$ in 138 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 7 for limits on the product of cross section times the branching ratios for m_{A^0} = 1–30 GeV.

[25] TUMASYAN 24A search for production of $H^0_{1,2}$ in association with W, Z, or $t\overline{t}$, decaying to $e^+ e^-$, $\mu^+ \mu^-$, or $\tau^+ \tau^-$, using 138 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 13 for limits on cross section times branching ratio for $m_{H^0_{1,2}}$ = 15–350 GeV. See also Figs. 16 and 17 for limits on the mixing with the Standard Model Higgs boson.

[26] TUMASYAN 24B search for $H^0_2 \to HH$ in the final state $\gamma\gamma b\overline{b}$ using 138 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 6 (upper) for limit on the product of production cross section times branching ratios for $m_{H^0_2}$ = 0.26–1.0 TeV.

[27] TUMASYAN 24B search for production of H^0_3 decaying to $H^0_{1,2} H \to b\overline{b}\gamma\gamma$ using 138 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 7 for limits on production cross section times branching ratios for $m_{H^0_3}$ = 0.3–1.0 TeV and $m_{H^0_{1,2}}$ = 0.09–0.8 TeV.

[28] AAD 23AD search for associated production of $W/Z H_2^0$ with the decay chain $H_2^0 \to HH \to b\overline{b}b\overline{b}$ using 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 9 for limits on cross section times branching ratios for $m_{H_2^0}$ = 260–1000 GeV.

[29] AAD 23AD search for gluon fusion production of A^0 with the decay chain $A^0 \to Z H_2^0$, $H_2^0 \to HH \to b\overline{b}b\overline{b}$ using 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 10 for limits on cross section times branching ratios for m_{A^0} = 350–800 GeV and $m_{H_2^0}$ = 260–400 GeV.

[30] AAD 23AJ search for production of $H^\pm$ in association with a top quark, followed by $H^\pm \to W^\pm A^0$, $A^0 \to$ invisible, using 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 10 for excluded parameter regions of 2HDM + CP-odd singlet model.

[31] AAD 23BD search for a top quark decaying to $qH_{1,2}^0$ ($q = u$, c), $H_{1,2}^0 \to b\overline{b}$, using 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 9 for limits on production cross section times branching ratios for $m_{H_{1,2}^0}$ = 20–160 GeV.

[32] AAD 23BE search for associated production of $H_2^0 W$ and decay $H_2^0 \to W^+ W^-$ assuming the presence of higher dimensional $H_2^0 W^+ W^-$ interactions, using 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 6 for excluded parameter region of higher dimensional operators, and Fig. 7 for limits on cross section times branching ratio for $m_{H_2^0}$ = 0.3–1.5 TeV.

[33] AAD 23BG search for production of H_2^0/A^0 in association with a $t\overline{t}$ pair, decaying to $t\overline{t}$, using 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 7 for limits on cross section times branching ratios for $m_{H_2^0} = m_{A^0}$ = 0.4–1.0 TeV.

[34] AAD 23BW search for A^0 production in association with a $t\overline{t}$ pair, decaying to $\mu^+\mu^-$, using 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 5(a) for limits on production cross section times branching ratio for m_{A^0} = 15–72 GeV.

[35] AAD 23BX search for production of $H \to \tau^+\tau^-$ with missing transverse momentum using 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 8 for interpretation of the data in terms of 2HDM + a model.

[36] AAD 23CA search for production of H_3^0 decaying to $H_2^0 H$, $H_2^0 \to W^+ W^-$ or ZZ, and $H \to \tau^+\tau^-$ using 140 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Figs. 4, 5 for limits on production cross section times branching ratios in the ranges $m_{H_3^0}$ = 0.5–1.5 TeV and $m_{H_2^0}$ = 0.2–0.5 TeV.

[37] AAD 23CR search for H_2^0 having flavor-violating couplings to tc or tu, produced in association with top quark(s), using 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 14 for limits on production cross section times branching ratios for $m_{H_2^0}$ = 0.2–1.5 TeV with various assumptions on the flavor-changing couplings.

[38] AAD 23O search for production of an A^0 in gluon-gluon fusion and in association with a $b\overline{b}$, decaying to ZH in the final states $\nu\overline{\nu}b\overline{b}$ and $\ell^+\ell^- b\overline{b}$ using 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 9 for limits on cross section times branching ratio for m_{A^0} = 0.22–2.0 TeV, and Fig. 11 for limits with both production components.

[39] AAD 23R search for the decay $A^0 \to \gamma\gamma$ in 138 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 7 for limits on cross section times branching ratio for m_{A^0} = 10–70 GeV.

[40] AAD 23U search for the decay $H_2^0 \to Z\gamma$ with Z decaying hadronically in 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 8(a) for limits on production cross section times branching ratios for $m_{H_2^0}$ = 1.0–6.8 TeV.

[41] AAD 23Z search for the decay chain $H_2^0 \to HH \to b\overline{b}\tau^+\tau^-$ using 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 10 for limits on the product of production cross section times branching ratios for $m_{H_2^0}$ = 0.251–1.6 TeV.

[42] HAYRAPETYAN 23C search for $H_{1,2}^0 \to e\mu$ using 138 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 7 for limits on production cross section times branching ratio for $m_{H_{1,2}^0}$ = 110–160 GeV.

[43] HAYRAPETYAN 23G search for dimuon resonance in the mass range 1.1–2.6 or 4.2–7.9 GeV in 96.6 fb^{-1} of pp collisions at E_{cm} = 13 TeV, in inclusive and high p_T selections. See their Fig. 5 for cross section times branching ratio limits and Fig. 7 for mixing angle limits in two Higgs doublet plus singlet model (at 90% CL).

[44] TUMASYAN 23 search for production of H_3^0 decaying to $H_{1,2}^0 H \to b\overline{b}b\overline{b}$ using 138 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 4 for limits on production cross section times branching ratios for $m_{H_3^0}$ = 0.9–4.0 TeV and $m_{H_{1,2}^0}$ = 60–600 GeV, and their interpretation in the NMSSM and the Two Real Singlet Model (TRSM).

[45] TUMASYAN 23M search for the decay chain $H \to A^0 A^0 \to \gamma\gamma\gamma\gamma$ in 132 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 6 for limits on cross section times branching ratio in the range m_{A^0} = 15–62 GeV.

[46] TUMASYAN 23O search for $H_2^0 \to HH$, each H decaying to either WW^* or $\tau^+\tau^-$ using 138 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 14 (upper) for limit on the product of production cross section times branching ratios for $m_{H_2^0}$ = 0.25–1.0 TeV.

[47] TUMASYAN 23S search for gluon fusion and b-associated production of $H_{1,2}^0$ decaying to $\tau^+\tau^-$ using 138 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 10 for limits on production cross section times branching ratios for $m_{H_{1,2}^0}$ = 0.06–3.5 TeV.

[48] AAD 22A search for the decay chain $H \to A^0 A^0 \to \mu^+\mu^- b\overline{b}$ in 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 9 for limits on the overall branching fraction in the range m_{A^0} = 16–62 GeV. See also Fig. 11 for limits without assuming A^0 is pseudoscalar.

[49] AAD 22D search for ZA^0 associate production with $Z \to \ell^+\ell^-$, A^0 decaying invisibly, in 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 5 for excluded regions in the mass parameter space of two Higgs doublet plus singlet (2HDM+A^0) model with a certain choice of the model parameters.

[50] AAD 22F search for gluon fusion production of H_2^0 decaying to $HH \to b\overline{b}b\overline{b}$ using 126–139 fb^{-1} of pp collisions at E_{cm} = 13 TeV. B($H \to b\overline{b}$) = 0.582 is assumed. See their Fig. 14 for limit on the product of production cross section times branching ratios for $m_{H_2^0}$ = 0.251–5.0 TeV.

[51] AAD 22I search for ZH associate production with the decay chain $H \to \widetilde{\chi}_2^0 \widetilde{\chi}_1^0$, $\widetilde{\chi}_2^0 \to A^0 \widetilde{\chi}_1^0$, $A^0 \to b\overline{b}$, and $Z \to \ell^+\ell^-$, in 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Figs. 3 and 4 for limits on the product of cross section times the branching ratios for m_{A^0} = 20–65 GeV with various choices of NMSSM model parameters.

[52] AAD 22J search for the decay $H \to ZA^0$ with $A^0 \to \mu^+\mu^-$ and $Z \to e^+e^-$, $\mu^+\mu^-$ in 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV assuming SM gluon-gluon fusion production of the H. See their Fig. 17(b) for limits on the product of cross section times the branching ratios for m_{A^0} = 15–30 GeV.

[53] AAD 22J search for the decay $H \to A^0 A^0$ with $A^0 \to \mu^+\mu^-$ in 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV assuming SM gluon-gluon fusion production of the H in the range of m_{A^0} = 1–60 GeV . See their Fig. 14(b) for limits on the product of cross

section times the branching ratios for m_{A^0} = 1.5–60 GeV (excluding ψ and Υ regions). The limit also applies to the decay $H \to H_1^0 H_1^0$.

[54] AAD 22P search for invisibly decaying H_1^0, H_2^0 produced by vector boson fusion in 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV. Limit on the product of cross section times branching ratio in the range 0.1–1 pb (95% CL) is given for the mass range 0.05–2 TeV. See their Fig. 14.

[55] AAD 22Y search for gluon fusion production of H_2^0 decaying to $HH \to b\overline{b}\gamma\gamma$ in 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 15 for limit on the product of production cross section times branching ratios to HH for $m_{H_2^0}$ = 0.251–1.0 TeV.

[56] ABRATENKO 22A search for a singlet scalar boson H_1^0 having a small mixing with the SM Higgs boson in the decay chain $K^+ \to H_1^0 \pi^+$, $H_1^0 \to \mu^+\mu^-$ from data corresponding to 7.01 × 10^{20} protons on NuMI target. See their Fig. 13 (right) and Table V for limits on the SM Higgs component of H_1^0 for $m_{H_1^0}$ = 212–279 MeV.

[57] TUMASYAN 22AK search for gluon-fusion production of H_3^0 decaying to $H_1^0 H_1^0 \to b\overline{b}b\overline{b}$ in 138 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 5 for limits on cross section times branching ratio for $m_{H_3^0}$ = 1–3 TeV, $m_{H_1^0}$ = 25–100 GeV.

[58] TUMASYAN 22D search for production of an H_2^0 (denoted radion in the paper) in gluon fusion and vector boson fusion, decaying to W^+W^- in the final states $\ell\nu$ + hadrons, using 137 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 7 for limits on cross section times branching ratio for $m_{H_2^0}$ = 1.0–4.5 TeV.

[59] AAD 21AF search for production of a heavy H_2^0 state decaying to ZZ in the final states $\ell^+\ell^-\ell'^+\ell'^-$ and $\ell^+\ell^-\nu\overline{\nu}$ in 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 4 for upper limits on cross section times branching ratio for $m_{H_2^0}$ = 0.2–2.0 TeV assuming ggF or VBF with narrow width approximation, and Fig. 5 for upper limits on cross section times branching ratio for $m_{H_2^0}$ = 0.4–2.0 TeV assuming ggF, and with several assumptions on its width.

[60] AAD 21AI search for production of an A^0 in gluon-gluon fusion and in association with a $b\overline{b}$, decaying to $ZH_2^0 \to \ell^+\ell^- b\overline{b}$ or $\ell^+\ell^- W^+W^-$ in 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Figs. 9 and 13 for cross section limits for m_{A^0} = 230–800 GeV and $m_{H_2^0}$ = 130–700 GeV.

[61] AAD 21AY search for production of a scalar resonance decaying to $\gamma\gamma$ in 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 5(a) for limits on fiducial cross section times branching ratio for $m_{H_2^0}$ = 0.16–3 TeV with narrow width approximation, and Table 2 with several assumptions on the width.

[62] AAD 21AZ search for production of A_2^0 decaying to HA_1^0 followed by $H \to \gamma\gamma$, $A_1^0 \to$ invisible in 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Figs. 10–12 for limits in terms of two-Higgs-doublet model plus singlet pseudoscalar and a fermionic Dark Matter particle.

[63] AAD 21BB search for production of A_2^0 by gluon fusion or associated $A_2^0 b\overline{b}$ production, decaying to HA_1^0 followed by $H \to b\overline{b}$, $A_1^0 \to$ invisible in 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 8 for limits in terms of two-Higgs-doublet plus singlet pseudoscalar model.

[64] AAD 21BE search for production of A_1^0 associated with a single top quark and either a light quark or a W boson, decaying to invisible final states, in 139 fb^{-1} of pp collisions

at E_{cm} = 13 TeV. See their Figs. 13–15 for limits in terms of two-Higgs-doublet model plus singlet pseudoscalar, which is assumed to decay to a pair of Dark Matter particles.

[65] ABRATENKO 21 search for a singlet scalar boson H^0_1 having a small mixing with the SM Higgs boson in the decay chain $K^+ \rightarrow H^0_1 \pi^+$, $H^0_1 \rightarrow e^+ e^-$ from data corresponding to 1.93×10^{20} protons on NuMI target. See their Fig. 2 for limits on the SM Higgs component of H^0_1 for $m_{H^0_1}$ = 3–210 MeV.

[66] SIRUNYAN 21A search for $H^0_2 \rightarrow Z A^0$ with $Z \rightarrow \ell^+ \ell^-$, A^0 decaying invisibly, in 137 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 8 for excluded regions in the mass parameter space of two Higgs doublet plus singlet model with a certain choice of the model parameters.

[67] TUMASYAN 21F search for gluon fusion production of H^0_3 decaying to $H H^0_{1,2} \rightarrow \tau^+ \tau^- b\overline{b}$ in 137 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Figs. 5 and 6 for limits on cross section times branching ratios for $m_{H^0_{1,2}}$ = 0.06–2.8 TeV and $m_{H^0_3}$ = 0.24–3.0 TeV.

[68] AAD 20AA search for $H^0_2/A^0 \rightarrow \tau^+ \tau^-$ produced by gluon fusion or b-associated production using 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 2(a), 2(b) for limits on the product of cross section and branching ratio for $m_{H^0_2}$, m_{A^0} = 0.2–2.5 TeV.

[69] AAD 20AI search for ZH production followed by the decay $H \rightarrow A^0 A^0 \rightarrow b\overline{b} b\overline{b}$ in 36 fb^{-1} of pp collisions at E_{cm} = 13 TeV. The search looks for collimated $A^0 \rightarrow b\overline{b}$ decays and is complementary to AABOUD 18BX. See their Fig. 10 for limits on the product of production cross section and branching ratios in the range m_{A^0} = 15–30 GeV.

[70] AAD 20AO search for gluon fusion production of H^0_2 decaying to $HH \rightarrow \tau^+ \tau^- b\overline{b}$ (with hadronically decaying $\tau^+ \tau^-$) using 139 fb^{-1} of pp collisions at E_{cm} = 13 TeV. Limit on the product of production cross section times branching ratios in the range 28–817 fb (95% CL) is given for m_{A^0} = 1.0–3.0 TeV, see their Fig. 13.

[71] AAD 20C combine searches for a scalar resonance decaying to HH in 36.1 fb^{-1} of pp collisions at E_{cm} = 13 TeV from AABOUD 19A, AABOUD 19O, AABOUD 18CQ, AABOUD 19T, AABOUD 18CW, and AABOUD 18BU. See their Fig. 5(a) for limits on cross section times branching ratio for $m_{H^0_2}$ = 0.26–3 TeV.

[72] AAD 20L search for b-associated production of H^0_2 decaying to $b\overline{b}$ in 27.8 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 8 for limits on the product of cross section and branching ratio for $m_{H^0_2}$ = 0.45–1.4 TeV.

[73] AAD 20X search for vector-boson-fusion production of H^0_2 decaying to HH using 126 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 5 for limits on the product of cross section and branching ratio for the assumptions of a narrow- and broad-width resonance.

[74] AAIJ 20AL search for dimuon resonance in the mass range 0.2–60 GeV in 5.1 fb^{-1} of pp collisions at E_{cm} = 13 TeV, in inclusive and b quark associated production. Displaced decays are searched for for masses below 3 GeV. See their Figs. 7–9 for cross section limits and Fig. 10 for limits for mixing angle in two Higgs doublet plus singlet model (at 90% CL).

[75] SIRUNYAN 20 search for the decay $H \rightarrow A^0 A^0 \rightarrow \tau^+ \tau^- \tau^+ \tau^-$ or $\tau^+ \tau^- \mu^+ \mu^-$ in 35.9 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 10 for limits on the product of production cross section (normalized to the SM) and branching ratios in the range m_{A^0} = 4–15 GeV.

[76] SIRUNYAN 20AA search for $H_2^0 \rightarrow Z A^0$, $A^0 \rightarrow b\overline{b}$ or $A^0 \rightarrow Z H_2^0$, $H_2^0 \rightarrow b\overline{b}$ in 35.9 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 7 for limits on the product of cross section and branching ratio for $m_{H_2^0}$ = 0.12–1 TeV and m_{A^0} = 0.03–1 TeV.

[77] SIRUNYAN 20AC search for gluon-fusion production of A^0 decaying to ZH in 35.9 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 5 for limits on the product of cross section and branching ratios for m_{A^0} = 220–400 GeV.

[78] SIRUNYAN 20AD search for lepton-flavor violating decays $H_2^0 \rightarrow \mu\tau$, $e\tau$ of gluon-fusion-produced H_2^0 in 35.9 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 5 (9) and Table 5 (6) for limits on production cross section times branching ratio for $m_{H_2^0}$ = 0.2–0.9 TeV for the $\mu\tau$ ($e\tau$) final state.

[79] SIRUNYAN 20AF search for $H_2^0/A^0 \rightarrow t\overline{t}$ with one or two charged leptons in the final state using kinematic variables in 35.9 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Figs. 5 and 6 for limits on top Yukawa coupling of H_2^0 and A^0 for $m_{H_2^0}$, m_{A^0} = 0.4–0.75 TeV for various width assumptions.

[80] SIRUNYAN 20AP search for the decay H or $H_2^0 \rightarrow A^0 A^0 \rightarrow \mu^+\mu^-\tau^+\tau^-$ (for $m_{H_2^0}$ = 300 GeV) with boosted final-state topology in 35.9 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 7 for limits on the product of production cross section (normalized to the SM) and branching ratios in the range m_{A^0} = 3.6–21 GeV, and Figs. 8 and 9 for its interpretation in terms of models with two Higgs doublets plus a singlet.

[81] SIRUNYAN 20Y search for gluon-fusion and vector-boson-fusion production of H_2^0 decaying to W^+W^- in the final states $\ell\nu\ell\nu$ and $\ell\nu qq$ in 35.9 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 6 for limits on the product of cross section and branching ratio for $m_{H_2^0}$ = 0.2–3 TeV.

[82] SIRUNYAN 20Z search for $H_{1,2}^0$ or A^0 production in association with a $t\overline{t}$ pair, decaying to e^+e^- or $\mu^+\mu^-$, in 137 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 12 for limits on production cross section times branching ratio for $m_{H_{1,2}^0}$, m_{A^0} = 15–75 GeV and 108–340 GeV.

[83] AABOUD 19A search for a narrow scalar resonance decaying to $HH \rightarrow b\overline{b}b\overline{b}$ in 27.5–36.1 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 9(a) for limits on cross section times branching ratios for $m_{H_2^0}$ = 0.26–3 TeV.

[84] AABOUD 19AG search for the decay $H \rightarrow A^0 A^0 \rightarrow \mu^+\mu^- b\overline{b}$ in 36.7 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 6 (a) for limits on the product of production cross section (normalized to the SM) and branching ratios in the range m_{A^0} = 20–60 GeV.

[85] AABOUD 19O search for a scalar resonance decaying to $HH \rightarrow b\overline{b}WW^*$ in 36.1 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 12 (left) for limits on cross section times branching ratio for $m_{H_2^0}$ = 0.5–3 TeV.

[86] AABOUD 19T search for a scalar resonance decaying to $HH \rightarrow WW^*WW^*$ in 36.1 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 3 for limits on cross section times branching ratio for $m_{H_2^0}$ = 260–500 GeV, assuming SM decay rates for the H.

[87] AABOUD 19V combine published ATLAS data to constrain two-Higgs-doublet plus singlet pseudoscalar model with A_1^0 decaying to invisible final states. See their Fig. 19 for excluded parameter regions.

[88] AABOUD 19Y search for a narrow scalar resonance produced by gluon fusion or b associated production, decaying to $\mu^+\mu^-$ in 36.1 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Figs. 4 and 5(a) for cross section limits for $m_{H_2^0}$ = 0.2–1.0 TeV.

[89] AALTONEN 19 search for b associated production of a scalar particle decaying to $b\overline{b}$ in 5.4 fb^{-1} of $p\overline{p}$ collisions at E_{cm} = 1.96 TeV. See their Fig. 3 for limits on cross section times branching ratio for $m_{H^0_{1,2}}$ = 100–300 GeV.

[90] SIRUNYAN 19 search for a narrow scalar resonance decaying to $HH \to \gamma\gamma b\overline{b}$ in 35.9 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 9 (left) for limits on cross section times branching ratios for $m_{H^0_2}$ = 260–900 GeV.

[91] SIRUNYAN 19AE search for a scalar resonance produced in association with a $b\overline{b}$ pair, decaying to $\tau^+\tau^-$ in 35.9 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 4 for cross section limits for m_{A^0} = 25–70 GeV.

[92] SIRUNYAN 19AN search for production of A^0_2 decaying to HA^0_1 followed by $H \to b\overline{b}$, $A^0_1 \to$ invisible in 35.9 fb^{-1} of pp collisions at E_{cm} = 13 TeV, in the mass range $m_{A^0_2}$ = 0.2–1.6 TeV, $m_{A^0_1}$ = 0.15–0.5 TeV. See their Fig. 6 for limits in terms of two-Higgs-doublet plus singlet pseudoscalar model.

[93] SIRUNYAN 19AV search for a scalar resonance produced by gluon fusion or b-associated production, decaying to $ZH \to \ell^+\ell^- b\overline{b}$ ($\ell = e, \mu$) or $\nu\overline{\nu}b\overline{b}$ in 35.9 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 5 for cross section limits for m_{A^0} = 0.22–1.0 TeV.

[94] SIRUNYAN 19B search for gluon fusion production of narrow scalar resonance with large transverse momentum, decaying to $b\overline{b}$, in 35.9 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Figs. 7 and 8 for limits on cross section times branching ratio for the resonance mass of 50–350 GeV.

[95] SIRUNYAN 19BB search for the decay $H^0_1 \to \gamma\gamma$ in 19.7 fb^{-1} of pp collisions at E_{cm} = 8 TeV and 35.9 fb^{-1} at E_{cm} = 13 TeV. See their Figs. 4–6 for limits on cross section times branching ratio for $m_{H^0_1}$ = 80–110 GeV (some results in Fig. 5 for $m_{H^0_1}$ = 70–110 GeV).

[96] SIRUNYAN 19BD search for the decay $H \to A^0A^0 \to \mu^+\mu^- b\overline{b}$ in 35.9 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 5 for limits on the product of cross section times branching ratios in the range m_{A^0} = 20–62.5 GeV. See also their Figs. 6 and 7 for interpretation of the data in terms of models with two Higgs doublets and a singlet.

[97] SIRUNYAN 19BE combine searches for $H^0_2 \to HH$ in 35.9 fb^{-1} of pp collisions at E_{cm} = 13 TeV in various H decay modes, from SIRUNYAN 18A, SIRUNYAN 18AF, SIRUNYAN 18CW, SIRUNYAN 19, and SIRUNYAN 19H. See their Fig. 3 for limits on cross section times branching ratios for $m_{H^0_2}$ = 0.25–3 TeV.

[98] SIRUNYAN 19BQ search for production of $H^0_{1,2}$ decaying to $A^0A^0 \to \mu^+\mu^-\mu^+\mu^-$ in 35.9 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 2 for limits on cross section times branching ratio for $m_{H^0_{1,2}}$ = 90–150 GeV, m_{A^0} = 0.25–3.55 GeV.

[99] SIRUNYAN 19CR search for production of H^0_2/A^0 in gluon fusion and in association with a $b\overline{b}$ pair, decaying to $\mu^+\mu^-$ in 35.9 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 6 for limits on cross section times branching ratio.

[100] SIRUNYAN 19H search for a narrow scalar resonance decaying to $HH \to b\overline{b}b\overline{b}$ in 35.9 fb^{-1} of pp collisions at E_{cm} = 13 TeV, where one $b\overline{b}$ pair is resolved and the other not. Limits on cross section times branching ratios for $m_{H^0_2}$ = 0.75–1.6 TeV are obtained and combined with data from SIRUNYAN 18AF. See their Fig. 5 (right).

[101] AABOUD 18AA search for production of a scalar resonance decaying to $Z\gamma$, with Z decaying hadronically, in 36.1 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 8(a) for limits on cross section times branching ratio for $m_{H^0_2}$ = 1.0–6.8 TeV.

[102] AABOUD 18AG search for the decay $H \to A^0 A^0 \to \gamma\gamma g g$ in 36.7 fb^{-1} of $p p$ collisions at E_{cm} = 13 TeV. See their Fig. 2 and Table 6 for cross section limits in the range m_{A^0} = 20–60 GeV.

[103] AABOUD 18AH search for production of an A^0 in gluon-gluon fusion and in association with a $b\overline{b}$, decaying to $Z H^0_2 \to \ell^+ \ell^- b\overline{b}$ in 36.1 fb^{-1} of $p p$ collisions at E_{cm} = 13 TeV. See their Fig. 5 for cross section limits for m_{A^0} = 230–800 GeV and $m_{H^0_2}$ = 130–700 GeV.

[104] AABOUD 18AI search for production of an A^0 in gluon-gluon fusion and in association with a $b\overline{b}$, decaying to $Z H$ in the final states $\nu\overline{\nu} b\overline{b}$ and $\ell^+ \ell^- b\overline{b}$ in 36.1 fb^{-1} of $p p$ collisions at E_{cm} = 13 TeV. See their Fig. 6 for cross section limits for m_{A^0} = 0.2–2 TeV. See also AABOUD 18CC.

[105] AABOUD 18BF search for production of a heavy H^0_2 state decaying to $Z Z$ in the final states $\ell^+ \ell^- \ell^+ \ell^-$ and $\ell^+ \ell^- \nu\overline{\nu}$ in 36.1 fb^{-1} of $p p$ collisions at E_{cm} = 13 TeV. See their Fig. 6 for upper limits on cross section times branching ratio for $m_{H^0_2}$ = 0.2–1.2 TeV assuming ggF or VBF with the NWA. See their Fig. 7 for upper limits on cross section times branching ratio for $m_{H^0_2}$ = 0.4–1.0 TeV assuming ggF, and with several assumptions on its width.

[106] AABOUD 18BU search for a narrow scalar resonance decaying to $H H \to \gamma\gamma W W^*$ in 36.1 fb^{-1} of $p p$ collisions at E_{cm} = 13 TeV. See their Fig. 4 for limits on cross section times branching ratios for $m_{H^0_2}$ = 260–500 GeV.

[107] AABOUD 18BX search for associated production of $W H$ or $Z H$ followed by the decay $H \to A^0 A^0 \to b\overline{b} b\overline{b}$ in 36.1 fb^{-1} of $p p$ collisions at E_{cm} = 13 TeV. See their Fig. 9 for limits on cross section times branching ratios for m_{A^0} = 20–60 GeV. See also their Fig. 10 for the dependence of the limit on A^0 lifetime.

[108] AABOUD 18CQ search for a narrow scalar resonance decaying to $H H \to b\overline{b}\tau^+\tau^-$ in 36.1 fb^{-1} of $p p$ collisions at E_{cm} = 13 TeV. See their Fig. 2 (above) for limits on cross section times branching ratios for $m_{H^0_2}$ = 260–1000 GeV.

[109] AABOUD 18F search for production of a narrow scalar resonance decaying to $W^+ W^-$ and $Z Z$, followed by hadronic decays of W and Z, in 36.7 fb^{-1} of $p p$ collisions at E_{cm} = 13 TeV. See their Fig. 5(c) for limits on cross section times branching ratio for $m_{H^0_2}$ = 1.2–3.0 TeV.

[110] AAIJ 18AM search for gluon-fusion production of $H^0_{1,2}$ decaying to $\mu\tau$ in 2 fb^{-1} of $p p$ collisions at E_{cm} = 8 TeV. See their Fig. 2 for limits on cross section times branching ratio for $m_{H^0_{1,2}}$ = 45–195 GeV.

[111] AAIJ 18AQ search for gluon-fusion production of a scalar particle A^0 decaying to $\mu^+ \mu^-$ in 1.99 fb^{-1} of $p p$ collisions at E_{cm} = 8 TeV and 0.98 fb^{-1} at E_{cm} = 7 TeV. See their Fig. 4 for limits on cross section times branching ratio for m_{A^0} = 5.5–15 GeV (using the E_{cm} = 8 TeV data set).

[112] AAIJ 18AQ search for the decay $H \to A^0 A^0$, with one of the A^0 decaying to $\mu^+ \mu^-$, in 1.99 fb^{-1} of $p p$ collisions at E_{cm} = 8 TeV and 0.98 fb^{-1} at E_{cm} = 7 TeV. See their Fig. 5 (right) for limits on the product of branching ratios for m_{A^0} = 5.5–15 GeV (using the E_{cm} = 8 TeV data set).

[113] SIRUNYAN 18AF search for a narrow scalar resonance decaying to $H H \to b\overline{b} b\overline{b}$ in 35.9 fb^{-1} of $p p$ collisions at E_{cm} = 13 TeV, where both $b\overline{b}$ pairs are not resolved. See their Fig. 9 for limits on cross section times branching ratios for $m_{H^0_2}$ = 0.75–3 TeV.

[114] SIRUNYAN 18BA search for production of a heavy H^0_2 state decaying to $Z Z$ in the final states $\ell^+ \ell^- \ell^+ \ell^-$, $\ell^+ \ell^- q\overline{q}$, and $\ell^+ \ell^- \nu\overline{\nu}$ in 35.9 fb^{-1} of $p p$ collisions at E_{cm} =

13 TeV. See their Figs. 10 and 11 for upper limits on cross section times branching ratio for $m_{H_2^0}$ = 0.13–3 TeV with several assumptions on its width and on the fraction of Vector-Boson-Fusion of the total production cross section.

115 SIRUNYAN 18CW search for a narrow scalar resonance decaying to $HH \to b\overline{b}b\overline{b}$ in 35.9 fb^{-1} of pp collisions at E_{cm} = 13 TeV, where both $b\overline{b}$ pairs are resolved. See their Fig. 9 for limits on cross section times branching ratios for $m_{H_2^0}$ = 260–1200 GeV.

116 SIRUNYAN 18DK search for production of a scalar resonance decaying to $Z\gamma$, with Z decaying to $\ell^+\ell^-$ or hadronically, in 35.9 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 7 for limits on cross section times branching ratio for $m_{H_2^0}$ = 0.35–4 TeV for different assumptions on the width of the resonance.

117 SIRUNYAN 18DT search for the decay $H \to A^0 A^0 \to \tau^+\tau^- b\overline{b}$ in 35.9 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 7 for limits on the product of branching ratios in the range m_{A^0} = 15–60 GeV. See also their Fig. 8 for interpretation of the data in terms of models with two Higgs doublets and a singlet.

118 SIRUNYAN 18DU search for production of a narrow scalar resonance decaying to $\gamma\gamma$ in 35.9 fb^{-1} (taken in 2016) of pp collisions at E_{cm} = 13 TeV. See their Fig. 3 (right) for limits on cross section times branching ratio for $m_{H_2^0}$ = 0.5–5 TeV for several values of its width-to-mass ratio.

119 SIRUNYAN 18ED search for production of an A^0 in gluon-gluon fusion and in association with a $b\overline{b}$, decaying to ZH in the final states $\nu\overline{\nu}b\overline{b}$ or $\ell^+\ell^- b\overline{b}$ in 35.9 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 8 for cross section limits for m_{A^0} = 0.8–2 TeV.

120 SIRUNYAN 18EE search for the decay $H \to A^0 A^0 \to \mu^+\mu^-\tau^+\tau^-$ in 35.9 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 4 for limits on the product of branching ratios in the range m_{A^0} = 15–62.5 GeV, normalized to the SM production cross section. See also their Fig. 5 for interpretation of the data in terms of models with two Higgs doublets and a singlet.

121 SIRUNYAN 18F search for a narrow scalar resonance decaying to $HH \to WWb\overline{b}$ or $ZZb\overline{b}$ in the final state $\ell\ell\nu\nu b\overline{b}$ in 35.9 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 7 for limits on cross section times branching ratios for $m_{H_2^0}$ = 250–900 GeV.

122 AABOUD 17 search for production of a scalar resonance decaying to $Z\gamma$ in 3.2 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 4 for the limits on cross section times branching ratio for $m_{H_2^0}$ = 0.25–3.0 TeV.

123 AABOUD 17AP search for production of a scalar resonance decaying to $\gamma\gamma$ in 36.7 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 4(a) for limits on fiducial cross section times branching ratio for $m_{H_2^0}$ = 0.2–2.7 TeV with narrow width approximation.

124 AABOUD 17AW search for production of a scalar resonance decaying to $Z\gamma$ in 36.1 fb^{-1} of pp collisions at E_{cm} = 13 TeV. See their Fig. 7 for limits on cross section times branching ratio for $m_{H_2^0}$ = 0.25–2.4 TeV.

125 KHACHATRYAN 17AZ search for the decay $H \to A^0 A^0 \to \tau^+\tau^-\tau^+\tau^-$, $\mu^+\mu^- b\overline{b}$, and $\mu^+\mu^-\tau^+\tau^-$ in 19.7 fb^{-1} of pp collisions at E_{cm} = 8 TeV. See their Figs. 4, 5, and 6 for cross section limits in the range m_{A^0} = 5–62.5 GeV. See also their Figs. 7, 8, and 9 for interpretation of the data in terms of models with two Higgs doublets and a singlet.

126 KHACHATRYAN 17D search for production of a scalar resonance decaying to $Z\gamma$ in 19.7 fb^{-1} of pp collisions at E_{cm} = 8 TeV and 2.7 fb^{-1} at E_{cm} = 13 TeV. See their Figs. 3 and 4 for the limits on cross section times branching ratio for $m_{H_2^0}$ = 0.2–2.0 TeV.

127 KHACHATRYAN 17R search for production of a narrow scalar resonance decaying to $\gamma\gamma$ in 12.9 fb^{-1} (taken in 2016) of pp collisions at E_{cm} = 13 TeV. See their Fig. 2 for

limits on cross section times branching ratio for $m_{H_2^0}$ = 0.5–4.5 TeV for several values of its width-to-mass ratio. Limits from combination with KHACHATRYAN 16M are shown in their Figs. 4 and 6.

[128] SIRUNYAN 17CN search for a narrow scalar resonance decaying to $H H \rightarrow b\overline{b}\tau^+\tau^-$ in 18.3 fb^{-1} of $p p$ collisions at E_{cm} = 8 TeV. See their Fig. 5 (above) and Table II for limits on the cross section times branching ratios for $m_{H_2^0}$ = 0.3–1 TeV, and Fig. 6 (above) and Table III for the corresponding limits by combining with data from KHACHATRYAN 16BQ and KHACHATRYAN 15R.

[129] SIRUNYAN 17Y search for production of a scalar resonance decaying to $Z\gamma$ in 19.7 fb^{-1} of $p p$ collisions at E_{cm} = 8 TeV and 2.7 fb^{-1} at E_{cm} = 13 TeV. See their Figs. 3, 4 and Table 3 for limits on cross section times branching ratio for $m_{H_2^0}$ = 0.7–3.0 TeV, and Fig. 5 for the corresponding limits for $m_{H_2^0}$ = 0.2–3.0 TeV from combination with KHACHATRYAN 17D data.

[130] AABOUD 16AB search for associated production of $W H$ with the decay $H \rightarrow A^0 A^0 \rightarrow b\overline{b}b\overline{b}$ in 3.2 fb^{-1} of $p p$ collisions at E_{cm} = 13 TeV. See their Fig. 8 for limits on cross section times branching ratios for m_{A^0} = 20–60 GeV.

[131] AABOUD 16AE search for production of a narrow scalar resonance decaying to $W^+ W^-$ and $Z Z$ in 3.2 fb^{-1} of $p p$ collisions at E_{cm} = 13 TeV. See their Fig. 4 for limits on cross section times branching ratio for $m_{H_2^0}$ = 0.5–3 TeV.

[132] AABOUD 16H search for production of a scalar resonance decaying to $\gamma\gamma$ in 3.2 fb^{-1} of $p p$ collisions at E_{cm} = 13 TeV. See their Fig. 12 for limits on cross section times branching ratio for $m_{H_2^0}$ = 0.2–2 TeV with different assumptions on the width.

[133] AABOUD 16I search for a narrow scalar resonance decaying to $H H \rightarrow b\overline{b}b\overline{b}$ in 3.2 fb^{-1} of $p p$ collisions at E_{cm} = 13 TeV. See their Fig. 10(c) for limits on cross section times branching ratios for $m_{H_2^0}$ = 0.5–3 TeV.

[134] AAD 16AX search for production of a heavy H state decaying to $Z Z$ in the final states $\ell^+\ell^-\ell^+\ell^-$, $\ell^+\ell^-\nu\overline{\nu}$, $\ell^+\ell^- q\overline{q}$, and$\nu\overline{\nu}q\overline{q}$ in 20.3 fb^{-1} of $p p$ collisions at E_{cm} = 8 TeV. See their Fig.12 for upper limits on $\sigma(H)$ B($H \rightarrow Z Z$) for m_H ranging from 140 GeV to 1000 GeV.

[135] AAD 16C search for production of a heavy H state decaying to $W^+ W^-$ in the final states $\ell\nu\ell\nu$ and $\ell\nu q q$ in 20.3 fb^{-1} of $p p$ collisions at E_{cm} = 8 TeV. See their Figs. 12, 13, and 16 for upper limits on $\sigma(H)$ B($H \rightarrow W^+ W^-$) for m_H ranging from 300 GeV to 1000 or 1500 GeV with various assumptions on the total width of H.

[136] AAD 16L search for the decay $H \rightarrow A^0 A^0 \rightarrow \gamma\gamma\gamma\gamma$ in 20.3 fb^{-1} of $p p$ collisions at E_{cm} = 8 TeV. See their Fig. 4 (upper right) for limits on cross section times branching ratios (normalized to the SM H cross section) for m_{A^0} = 10–60 GeV.

[137] AAD 16L search for the decay $H_2^0 \rightarrow A^0 A^0 \rightarrow \gamma\gamma\gamma\gamma$ in 20.3 fb^{-1} of $p p$ collisions at E_{cm} = 8 TeV. See their Fig. 4 (lower right) for limits on cross section times branching ratios for $m_{H_2^0}$ = 600 GeV and m_{A^0} = 10–245 GeV, and Table 5 for limits for $m_{H_2^0}$= 300 and 900 GeV.

[138] AALTONEN 16C search for electroweak associated production of $H_1^0 H^\pm$ followed by the decays $H^\pm \rightarrow H_1^0 W^*$, $H_1^0 \rightarrow \gamma\gamma$ for $m_{H_1^0}$ = 10–105 GeV and $m_{H^\pm}$ = 30–300 GeV. See their Fig. 3 for excluded parameter region in a two-doublet model in which H_1^0 has no direct decay to fermions.

[139] KHACHATRYAN 16BG search for a narrow scalar resonance decaying to $H H \rightarrow b\overline{b}b\overline{b}$ in 19.7 fb^{-1} of $p p$ collisions at E_{cm} = 8 TeV. See their Fig. 6 for limits on the cross section times branching ratios for $m_{H_2^0}$ = 1.15–3 TeV.

[140] KHACHATRYAN 16BQ search for a resonance decaying to $HH \rightarrow \gamma\gamma b\overline{b}$ in 19.7 fb^{-1} of pp collisions at E_{cm} = 8 TeV. See their Fig. 9 for limits on the cross section times branching ratios for $m_{H_2^0}$ = 0.26–1.1 TeV.

[141] KHACHATRYAN 16F search for the decay $H \rightarrow H_1^0 H_1^0 \rightarrow \tau^+\tau^-\tau^+\tau^-$ in 19.7 fb^{-1} of pp collisions at E_{cm} = 8 TeV. See their Fig. 8 for cross section limits for $m_{H_1^0}$ = 4–8 GeV.

[142] KHACHATRYAN 16M search for production of a narrow resonance decaying to $\gamma\gamma$ in 19.7 fb^{-1} of pp collisions at E_{cm} = 8 TeV and 3.3 fb^{-1} at E_{cm} = 13 TeV. See their Fig. 3 (top) for limits on cross section times branching ratio for $m_{H_2^0}$ = 0.5–4 TeV.

[143] KHACHATRYAN 16P search for gluon fusion production of an H_2^0 decaying to $HH \rightarrow b\overline{b}\tau^+\tau^-$ in 19.7 fb^{-1} of pp collisions at E_{cm} = 8 TeV. See their Fig. 8 (lower right) for cross section limits for $m_{H_2^0}$ = 260–350 GeV.

[144] KHACHATRYAN 16P search for gluon fusion production of an A^0 decaying to $ZH \rightarrow \ell^+\ell^-\tau^+\tau^-$ in 19.7 fb^{-1} of pp collisions at E_{cm} = 8 TeV. See their Fig. 10 for cross section limits for $m_{H_2^0}$ = 220–350 GeV.

[145] AAD 15BK search for production of a heavy H_2^0 decaying to HH in the final state $b\overline{b}b\overline{b}$ in 19.5 fb^{-1} of pp collisions at E_{cm} = 8 TeV. See their Fig. 14(c) for $\sigma(H_2^0)$ B($H_2^0 \rightarrow HH$) for $m_{H_2^0}$ = 500–1500 GeV with $\Gamma_{H_2^0}$ = 1 GeV.

[146] AAD 15BZ search for the decay $H \rightarrow A^0 A^0 \rightarrow \mu^+\mu^-\tau^+\tau^-$ (m_H = 125 GeV) in 20.3 fb^{-1} of pp collisions at E_{cm} = 8 TeV. See their Fig. 6 for limits on cross section times branching ratio for m_{A^0} = 3.7–50 GeV.

[147] AAD 15BZ search for a state H_2^0 via the decay $H_2^0 \rightarrow A^0 A^0 \rightarrow \mu^+\mu^-\tau^+\tau^-$ in 20.3 fb^{-1} of pp collisions at E_{cm} = 8 TeV. See their Fig. 6 for limits on cross section times branching ratio for $m_{H_2^0}$ = 100–500 GeV and m_{A^0} = 5 GeV.

[148] AAD 15CE search for production of a heavy H_2^0 decaying to HH in the final states $b\overline{b}\tau^+\tau^-$ and $\gamma\gamma W W^*$ in 20.3 fb^{-1} of pp collisions at E_{cm} = 8 TeV and combine with data from AAD 15H and AAD 15BK. A limit $\sigma(H_2^0)$ B($H_2^0 \rightarrow HH$) $<$ 2.1–0.011 pb (95% CL) is given for $m_{H_2^0}$ = 260–1000 GeV. See their Fig. 6.

[149] AAD 15H search for production of a heavy H_2^0 decaying to HH in the finalstate $\gamma\gamma b\overline{b}$ in 20.3 fb^{-1} of pp collisions at E_{cm} = 8 TeV.A limit of $\sigma(H_2^0)$ B($H_2^0 \rightarrow HH$) $<$ 3.5–0.7 pb is given for $m_{H_2^0}$ = 260–500 GeV at 95% CL. See their Fig. 3.

[150] AAD 15S search for production of A^0 decaying to $ZH \rightarrow \ell^+\ell^- b\overline{b}$, $\nu\overline{\nu} b\overline{b}$ and $\ell^+\ell^-\tau^+\tau^-$ in 20.3 fb^{-1} of pp collisions at E_{cm} = 8 TeV. See their Fig. 3 for cross section limits for m_{A^0} = 200–1000 GeV.

[151] KHACHATRYAN 15AW search for production of a heavy state H_2^0 of an electroweak singlet extension of the Standard Model via the decays of H_2^0 to W^+W^- and ZZ in up to 5.1 fb^{-1} of pp collisions at E_{cm} = 7 TeV and up to 19.7 fb^{-1} at E_{cm} = 8 TeV in the range $m_{H_2^0}$ = 145–1000 GeV. See their Figs. 8 and 9 for limits in the parameter space of the model.

[152] KHACHATRYAN 15BB search for production of a resonance H decaying to $\gamma\gamma$ in 19.7 fb^{-1} of pp collisions at E_{cm} = 8 TeV. See their Fig. 7 for limits on cross section times branching ratio for m_H = 150–850 GeV.

[153] KHACHATRYAN 15N search for production of A^0 decaying to $ZH \rightarrow \ell^+\ell^- b\overline{b}$ in 19.7 fb^{-1} of pp collisions at $E_{cm} = 8$ TeV. See their Fig. 3 for limits on cross section times branching ratios for $m_{A^0} = 225$–600 GeV.

[154] KHACHATRYAN 15O search for production of a high-mass narrow resonance A^0 decaying to $ZH \rightarrow q\overline{q}\tau^+\tau^-$ in 19.7 fb^{-1} of pp collisions at $E_{cm} = 8$ TeV. See their Fig. 6 for limits on cross section times branching ratios for $m_{A^0} = 800$–2500 GeV.

[155] KHACHATRYAN 15R search for a narrow scalar resonance decaying to $HH \rightarrow b\overline{b}b\overline{b}$ in 17.9 fb^{-1} of pp collisions at $E_{cm} = 8$ TeV. See their Fig. 5 (top) for limits on cross section times branching ratios for $m_{H_2^0} = 0.27$–1.1 TeV.

[156] AAD 14AP search for a second H state decaying to $\gamma\gamma$ in addition to the state at about 125 GeV in 20.3 fb^{-1} of pp collisions at $E_{cm} = 8$ TeV. See their Fig. 4 for limits on cross section times branching ratio for $m_H = 65$–600 GeV.

[157] AAD 14M search for the decay cascade $H_2^0 \rightarrow H^\pm W^\mp \rightarrow HW^\pm W^\mp$, H decaying to $b\overline{b}$ in 20.3 fb^{-1} of pp collisions at $E_{cm} = 8$ TeV. See their Table III for limits on cross section times branching ratio for $m_{H_2^0} = 325$–1025 GeV and $m_{H^+} = 225$–925 GeV.

[158] CHATRCHYAN 14G search for a second H state decaying to $WW^{(*)}$ in addition to the observed signal at about 125 GeV using 4.9 fb^{-1} of pp collisions at $E_{cm} = 7$ TeV and 19.4 fb^{-1} at $E_{cm} = 8$ TeV. See their Fig. 21 (right) for cross section limits in the mass range 110–600 GeV.

[159] KHACHATRYAN 14P search for a second H state decaying to $\gamma\gamma$ in addition to the observed signal at about 125 GeV using 5.1 fb^{-1} of pp collisions at $E_{cm} = 7$ TeV and 19.7 fb^{-1} at $E_{cm} = 8$ TeV. See their Figs. 27 and 28 for cross section limits in the mass range 110–150 GeV.

[160] AALTONEN 13P search for production of a heavy Higgs boson H'^0 that decays into a charged Higgs boson $H^\pm$ and a lighter Higgs boson H via the decay chain $H'^0 \rightarrow H^\pm W^\mp$, $H^\pm \rightarrow W^\pm H$, $H \rightarrow b\overline{b}$ in the final state $\ell\nu$ plus 4 jets in 8.7 fb^{-1} of $p\overline{p}$ collisions at $E_{cm} = 1.96$ TeV. See their Fig. 4 for limits on cross section times branching ratio in the $m_{H^\pm}-m_{H'^0}$ plane for $m_H = 126$ GeV.

[161] CHATRCHYAN 13BJ search for H production in the decay chain $H \rightarrow A^0 A^0$, $A^0 \rightarrow \mu^+\mu^-$ in 5.3 fb^{-1} of pp collisions at $E_{cm} = 7$ TeV. See their Fig. 2 for limits on cross section times branching ratio.

[162] AALTONEN 11P search in 2.7 fb^{-1} of $p\overline{p}$ collisions at $E_{cm} = 1.96$ TeV for the decay chain $t \rightarrow bH^+$, $H^+ \rightarrow W^+ A^0$, $A^0 \rightarrow \tau^+\tau^-$ with m_{A^0} between 4 and 9 GeV. See their Fig. 4 for limits on B($t \rightarrow bH^+$) for $90 < m_{H^+} < 160$ GeV.

[163] ABBIENDI 10 search for $e^+e^- \rightarrow ZH$ with the decay chain $H \rightarrow \widetilde{\chi}_1^0\widetilde{\chi}_2^0$, $\widetilde{\chi}_2^0 \rightarrow \widetilde{\chi}_1^0$ + (γ or Z^*), when $\widetilde{\chi}_1^0$ and $\widetilde{\chi}_2^0$ are nearly degenerate. For a mass difference of 2 (4) GeV, a lower limit on m_H of 108.4 (107.0) GeV (95% CL) is obtained for SM ZH cross section and B($H \rightarrow \widetilde{\chi}_1^0\widetilde{\chi}_2^0$) = 1.

[164] SCHAEL 10 search for the process $e^+e^- \rightarrow HZ$ followed by the decay chain $H \rightarrow A^0 A^0 \rightarrow \tau^+\tau^-\tau^+\tau^-$ with $Z \rightarrow \ell^+\ell^-$, $\nu\overline{\nu}$ at $E_{cm} = 183$–209 GeV. For a HZZ coupling equal to the SM value, B($H \rightarrow A^0 A^0$) = B($A^0 \rightarrow \tau^+\tau^-$) = 1, and $m_{A^0} =$ 4–10 GeV, m_H up to 107 GeV is excluded at 95% CL.

[165] ABAZOV 09V search for H production followed by the decay chain $H \rightarrow A^0 A^0 \rightarrow \mu^+\mu^-\mu^+\mu^-$ or $\mu^+\mu^-\tau^+\tau^-$ in 4.2 fb^{-1} of $p\overline{p}$ collisions at $E_{cm} = 1.96$ TeV. See their Fig. 3 for limits on $\sigma(H)$·B($H \rightarrow A^0 A^0$) for $m_{A^0} = 3.6$–19 GeV.

[166] ABBIENDI 05A search for $e^+e^- \rightarrow H_1^0 A^0$ in general Type-II two-doublet models, with decays H_1^0, $A^0 \rightarrow q\overline{q}$, gg, $\tau^+\tau^-$, and $H_1^0 \rightarrow A^0 A^0$.

[167] ABBIENDI 04K search for $e^+e^- \to HZ$ with H decaying to two jets of any flavor including gg. The limit is for SM production cross section with B($H \to jj$) = 1.
[168] ABDALLAH 04 consider the full combined LEP and LEP2 datasets to set limits on the Higgs coupling to W or Z bosons, assuming SM decays of the Higgs. Results in Fig. 26.
[169] ACHARD 04B search for $e^+e^- \to HZ$ with H decaying to $b\overline{b}$, $c\overline{c}$, or gg. The limit is for SM production cross section with B($H \to jj$) = 1.
[170] ACHARD 04F search for H with anomalous coupling to gauge boson pairs in the processes $e^+e^- \to H\gamma$, e^+e^-H, HZ with decays $H \to f\overline{f}$, $\gamma\gamma$, $Z\gamma$, and W^*W at E_{cm} = 189–209 GeV. See paper for limits.
[171] ABBIENDI 03F search for $H \to$ anything in $e^+e^- \to HZ$, using the recoil mass spectrum of $Z \to e^+e^-$ or $\mu^+\mu^-$. In addition, it searched for $Z \to \nu\overline{\nu}$ and $H \to e^+e^-$ or photons. Scenarios with large width or continuum H mass distribution are considered. See their Figs. 11–14 for the results.
[172] ABBIENDI 03G search for $e^+e^- \to H^0_1 Z$ followed by $H^0_1 \to A^0 A^0$, $A^0 \to c\overline{c}$, gg, or $\tau^+\tau^-$ in the region $m_{H^0_1}$ = 45-86 GeV and m_{A^0} = 2-11 GeV. See their Fig. 7 for the limits.
[173] Search for associated production of a $\gamma\gamma$ resonance with a Z boson, followed by $Z \to q\overline{q}$, $\ell^+\ell^-$, or $\nu\overline{\nu}$, at $E_{cm} \leq$ 209 GeV. The limit is for a H with SM production cross section and B($H \to f\overline{f}$)=0 for all fermions f.
[174] For B($H \to \gamma\gamma$)=1, $m_H >$ 113.1 GeV is obtained.
[175] HEISTER 02M search for $e^+e^- \to HZ$, assuming that H decays to $q\overline{q}$, gg, or $\tau^+\tau^-$ only. The limit assumes SM production cross section.
[176] ABBIENDI 01E search for neutral Higgs bosons in general Type-II two-doublet models, at $E_{cm} \leq$ 189 GeV. In addition to usual final states, the decays H^0_1, $A^0 \to q\overline{q}$, gg are searched for. See their Figs. 15,16 for excluded regions.
[177] ACCIARRI 00R search for $e^+e^- \to H\gamma$ with $H \to b\overline{b}$, $Z\gamma$, or $\gamma\gamma$. See their Fig. 3 for limits on $\sigma \cdot$ B. Explicit limits within an effective interaction framework are also given, for which the Standard Model Higgs search results are used in addition.
[178] ACCIARRI 00R search for the two-photon type processes $e^+e^- \to e^+e^-H$ with $H \to b\overline{b}$ or $\gamma\gamma$. See their Fig. 4 for limits on $\Gamma(H \to \gamma\gamma)\cdot$B($H \to \gamma\gamma$ or $b\overline{b}$) for m_H=70–170 GeV.
[179] GONZALEZ-GARCIA 98B use DØ limit for $\gamma\gamma$ events with missing E_T in $p\overline{p}$ collisions (ABBOTT 98) to constrain possible ZH or WH production followed by unconventional $H \to \gamma\gamma$ decay which is induced by higher-dimensional operators. See their Figs. 1 and 2 for limits on the anomalous couplings.
[180] KRAWCZYK 97 analyse the muon anomalous magnetic moment in a two-doublet Higgs model (with type II Yukawa couplings) assuming no $H^0_1 ZZ$ coupling and obtain $m_{H^0_1} \gtrsim$ 5 GeV or $m_{A^0} \gtrsim$ 5 GeV for $\tan\beta >$ 50. Other Higgs bosons are assumed to be much heavier.
[181] ALEXANDER 96H give B($Z \to H\gamma$)×B($H \to q\overline{q}$) $<$ 1–4 $\times 10^{-5}$ (95%CL) and B($Z \to H\gamma$)×B($H \to b\overline{b}$) $<$ 0.7–2 $\times 10^{-5}$ (95%CL) in the range 20 $< m_H <$ 80 GeV.

Electroweak Constraints on the Standard Model Higgs Boson Mass

Here we list constraints on the mass of the Higgs boson derived from fits to precision electroweak observables, assuming the minimal Standard Model with a doublet Higgs field and three generations of fermions.

VALUE (GeV)	DOCUMENT ID		TECN
$\mathbf{90^{+21}_{-18}}$	[1] HALLER	18	RVUE

• • • We do not use the following data for averages, fits, limits, etc. • • •

91^{+30}_{-23}	[2] BAAK	12	RVUE
94^{+25}_{-22}	[3] BAAK	12A	RVUE
91^{+31}_{-24}	[4] ERLER	10A	RVUE
129^{+74}_{-49}	[5] LEP-SLC	06	RVUE

[1] HALLER 18 make Standard Model fits to Z and neutral current parameters, m_t, m_W, and Γ_W measurements available in 2018. The direct mass measurement at the LHC is not used in the fit.

[2] BAAK 12 make Standard Model fits to Z and neutral current parameters, m_t, m_W, and Γ_W measurements available in 2010 (using also preliminary data). The quoted result is obtained from a fit that does not include the limit from the direct Higgs searches. The result including direct search data from LEP2, the Tevatron and the LHC is 120^{+12}_{-5} GeV.

[3] BAAK 12A make Standard Model fits to Z and neutral current parameters, m_t, m_W, and Γ_W measurements available in 2012 (using also preliminary data). The quoted result is obtained from a fit that does not include the measured mass value of the signal observed at the LHC and also no limits from direct Higgs searches.

[4] ERLER 10A makes Standard Model fits to Z and neutral current parameters, m_t, m_W measurements available in 2009 (using also preliminary data). The quoted result is obtained from a fit that does not include the limits from the direct Higgs searches. With direct search data from LEP2 and Tevatron added to the fit, the 90% CL (99% CL) interval is 115–148 (114–197) GeV.

[5] LEP-SLC 06 make Standard Model fits to Z parameters from LEP/SLC and m_t, m_W, and Γ_W measurements available in 2005 with $\Delta\alpha^{(5)}_{\text{had}}(m_Z) = 0.02758 \pm 0.00035$. The 95% CL limit is 285 GeV.

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AAD	24AB	JHEP 2402 197	G. Aad *et al.*	(ATLAS Collab.)
AAD	24AP	JHEP 2408 013	G. Aad *et al.*	(ATLAS Collab.)
AAD	24BQ	PR D110 052013	G. Aad *et al.*	(ATLAS Collab.)
AAD	24BV	PL B858 139007	G. Aad *et al.*	(ATLAS Collab.)
AAD	24BW	JHEP 2410 104	G. Aad *et al.*	(ATLAS Collab.)
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AAD	24CA	JHEP 2411 047	G. Aad *et al.*	(ATLAS Collab.)
AAD	24CB	JHEP 2411 097	G. Aad *et al.*	(ATLAS Collab.)
AAD	24CD	JHEP 2411 126	G. Aad *et al.*	(ATLAS Collab.)
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ADACHI	24F	PR D109 112015	I. Adachi *et al.*	(BELLE II Collab.)
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CHATRCHYAN	14G	JHEP 1401 096	S. Chatrchyan *et al.*	(CMS Collab.)
KHACHATRY...	14M	JHEP 1410 160	V. Khachatryan *et al.*	(CMS Collab.)
KHACHATRY...	14P	EPJ C74 3076	V. Khachatryan *et al.*	(CMS Collab.)
KHACHATRY...	14Q	PR D90 112013	V. Khachatryan *et al.*	(CMS Collab.)
AAD	13AG	PL B721 32	G. Aad *et al.*	(ATLAS Collab.)
AAD	13AT	NJP 15 043009	G. Aad *et al.*	(ATLAS Collab.)
AAD	13O	JHEP 1302 095	G. Aad *et al.*	(ATLAS Collab.)
AAIJ	13T	JHEP 1305 132	R. Aaij *et al.*	(LHCb Collab.)
AALTONEN	13K	PR D88 052012	T. Aaltonen *et al.*	(CDF Collab.)
AALTONEN	13L	PR D88 052013	T. Aaltonen *et al.*	(CDF Collab.)
AALTONEN	13M	PR D88 052014	T. Aaltonen *et al.*	(CDF and D0 Collabs.)
AALTONEN	13P	PRL 110 121801	T. Aaltonen *et al.*	(CDF Collab.)
ABAZOV	13G	PR D88 052006	V.M. Abazov *et al.*	(D0 Collab.)
ABAZOV	13H	PR D88 052007	V.M. Abazov *et al.*	(D0 Collab.)
ABAZOV	13I	PR D88 052008	V.M. Abazov *et al.*	(D0 Collab.)
ABAZOV	13J	PR D88 052009	V.M. Abazov *et al.*	(D0 Collab.)
ABAZOV	13L	PR D88 052011	V.M. Abazov *et al.*	(D0 Collab.)
CARENA	13	EPJ C73 2552	M. Carena *et al.*	
CHATRCHYAN	13AG	PL B722 207	S. Chatrchyan *et al.*	(CMS Collab.)
CHATRCHYAN	13AL	PL B725 36	S. Chatrchyan *et al.*	(CMS Collab.)
CHATRCHYAN	13BJ	PL B726 564	S. Chatrchyan *et al.*	(CMS Collab.)
LEES	13C	PR D87 031102	J.P. Lees *et al.*	(BABAR Collab.)
LEES	13L	PR D88 031701	J.P. Lees *et al.*	(BABAR Collab.)
LEES	13R	PR D88 071102	J.P. Lees *et al.*	(BABAR Collab.)
AAD	12AI	PL B716 1	G. Aad *et al.*	(ATLAS Collab.)
AAD	12AQ	PRL 108 251801	G. Aad *et al.*	(ATLAS Collab.)
AAD	12N	EPJ C72 2157	G. Aad *et al.*	(ATLAS Collab.)
AALTONEN	12AB	PR D85 092001	T. Aaltonen *et al.*	(CDF Collab.)
AALTONEN	12AN	PL B717 173	T. Aaltonen *et al.*	(CDF Collab.)
AALTONEN	12AQ	PR D86 091101	T. Aaltonen *et al.*	(CDF and D0 Collabs.)
AALTONEN	12U	PR D85 012007	T. Aaltonen *et al.*	(CDF Collab.)
AALTONEN	12X	PR D85 032005	T. Aaltonen *et al.*	(CDF Collab.)
ABAZOV	12G	PL B710 569	V.M. Abazov *et al.*	(D0 Collab.)
ABLIKIM	12	PR D85 092012	M. Ablikim *et al.*	(BESIII Collab.)
BAAK	12	EPJ C72 2003	M. Baak *et al.*	(Gfitter Group)
BAAK	12A	EPJ C72 2205	M. Baak *et al.*	(Gfitter Group)
CHATRCHYAN	12AO	JHEP 1209 111	S. Chatrchyan *et al.*	(CMS Collab.)
CHATRCHYAN	12C	JHEP 1203 081	S. Chatrchyan *et al.*	(CMS Collab.)
CHATRCHYAN	12D	JHEP 1204 036	S. Chatrchyan *et al.*	(CMS Collab.)
CHATRCHYAN	12E	PL B710 91	S. Chatrchyan *et al.*	(CMS Collab.)
CHATRCHYAN	12G	PL B710 403	S. Chatrchyan *et al.*	(CMS Collab.)
CHATRCHYAN	12H	PRL 108 111804	S. Chatrchyan *et al.*	(CMS Collab.)
CHATRCHYAN	12I	JHEP 1203 040	S. Chatrchyan *et al.*	(CMS Collab.)
CHATRCHYAN	12K	PL B713 68	S. Chatrchyan *et al.*	(CMS Collab.)
CHATRCHYAN	12N	PL B716 30	S. Chatrchyan *et al.*	(CMS Collab.)
CHATRCHYAN	12V	PRL 109 121801	S. Chatrchyan *et al.*	(CMS Collab.)
AALTONEN	11P	PRL 107 031801	T. Aaltonen *et al.*	(CDF Collab.)
ABAZOV	11K	PL B698 97	V.M. Abazov *et al.*	(D0 Collab.)
ABAZOV	11W	PRL 107 121801	V.M. Abazov *et al.*	(D0 Collab.)
ABOUZAID	11A	PRL 107 201803	E. Abouzaid *et al.*	(KTeV Collab.)
DEL-AMO-SA...	11J	PRL 107 021804	P. del Amo Sanchez *et al.*	(BABAR Collab.)

LEES	11H	PRL 107 221803	J.P. Lees *et al.*	(BABAR Collab.)
ABBIENDI	10	PL B682 381	G. Abbiendi *et al.*	(OPAL Collab.)
ANDREAS	10	JHEP 1008 003	S. Andreas *et al.*	(DESY)
ERLER	10A	PR D81 051301	J. Erler	(UNAM)
HYUN	10	PRL 105 091801	H.J. Hyun *et al.*	(BELLE Collab.)
SCHAEL	10	JHEP 1005 049	S. Schael *et al.*	(ALEPH Collab.)
AALTONEN	09AB	PRL 103 061803	T. Aaltonen *et al.*	(CDF Collab.)
AALTONEN	09AR	PRL 103 201801	T. Aaltonen *et al.*	(CDF Collab.)
ABAZOV	09V	PRL 103 061801	V.M. Abazov *et al.*	(D0 Collab.)
AUBERT	09P	PRL 103 181801	B. Aubert *et al.*	(BABAR Collab.)
AUBERT	09Z	PRL 103 081803	B. Aubert *et al.*	(BABAR Collab.)
TUNG	09	PRL 102 051802	Y.C. Tung *et al.*	(KEK E391a Collab.)
ABAZOV	08U	PRL 101 051801	V.M. Abazov *et al.*	(D0 Collab.)
ABDALLAH	08B	EPJ C54 1	J. Abdallah *et al.*	(DELPHI Collab.)
Also		EPJ C56 165 (errat.)	J. Abdallah *et al.*	(DELPHI Collab.)
LOVE	08	PRL 101 151802	W. Love *et al.*	(CLEO Collab.)
ABBIENDI	07	EPJ C49 457	G. Abbiendi *et al.*	(OPAL Collab.)
BESSON	07	PRL 98 052002	D. Besson *et al.*	(CLEO Collab.)
SCHAEL	07	EPJ C49 439	S. Schael *et al.*	(ALEPH Collab.)
LEP-SLC	06	PRPL 427 257	ALEPH, DELPHI, L3, OPAL, SLD and working groups	
SCHAEL	06B	EPJ C47 547	S. Schael *et al.*	(LEP Collabs.)
ABBIENDI	05A	EPJ C40 317	G. Abbiendi *et al.*	(OPAL Collab.)
ABDALLAH	05D	EPJ C44 147	J. Abdallah *et al.*	(DELPHI Collab.)
ACHARD	05	PL B609 35	P. Achard *et al.*	(L3 Collab.)
ACOSTA	05Q	PR D72 072004	D. Acosta *et al.*	(CDF Collab.)
PARK	05	PRL 94 021801	H.K. Park *et al.*	(FNAL HyperCP Collab.)
ABBIENDI	04K	PL B597 11	G. Abbiendi *et al.*	(OPAL Collab.)
ABBIENDI	04M	EPJ C37 49	G. Abbiendi *et al.*	(OPAL Collab.)
ABDALLAH	04	EPJ C32 145	J. Abdallah *et al.*	(DELPHI Collab.)
ABDALLAH	04B	EPJ C32 475	J. Abdallah *et al.*	(DELPHI Collab.)
ABDALLAH	04L	EPJ C35 313	J. Abdallah *et al.*	(DELPHI Collab.)
ABDALLAH	04O	EPJ C38 1	J. Abdallah *et al.*	(DELPHI Collab.)
ACHARD	04B	PL B583 14	P. Achard *et al.*	(L3 Collab.)
ACHARD	04F	PL B589 89	P. Achard *et al.*	(L3 Collab.)
ABBIENDI	03F	EPJ C27 311	G. Abbiendi *et al.*	(OPAL Collab.)
ABBIENDI	03G	EPJ C27 483	G. Abbiendi *et al.*	(OPAL Collab.)
ACHARD	03C	PL B568 191	P. Achard *et al.*	(L3 Collab.)
ABBIENDI	02D	EPJ C23 397	G. Abbiendi *et al.*	(OPAL Collab.)
ABBIENDI	02F	PL B544 44	G. Abbiendi *et al.*	(OPAL Collab.)
ACHARD	02C	PL B534 28	P. Achard *et al.*	(L3 Collab.)
ACHARD	02H	PL B545 30	P. Achard *et al.*	(L3 Collab.)
AKEROYD	02	PR D66 037702	A.G. Akeroyd *et al.*	
HEISTER	02	PL B526 191	A. Heister *et al.*	(ALEPH Collab.)
HEISTER	02L	PL B544 16	A. Heister *et al.*	(ALEPH Collab.)
HEISTER	02M	PL B544 25	A. Heister *et al.*	(ALEPH Collab.)
ABBIENDI	01E	EPJ C18 425	G. Abbiendi *et al.*	(OPAL Collab.)
ABREU	01F	PL B507 89	P. Abreu *et al.*	(DELPHI Collab.)
AFFOLDER	01H	PR D64 092002	T. Affolder *et al.*	(CDF Collab.)
BARATE	01C	PL B499 53	R. Barate *et al.*	(ALEPH Collab.)
ACCIARRI	00M	PL B485 85	M. Acciarri *et al.*	(L3 Collab.)
ACCIARRI	00R	PL B489 102	M. Acciarri *et al.*	(L3 Collab.)
ACCIARRI	00S	PL B489 115	M. Acciarri *et al.*	(L3 Collab.)
BARATE	00L	PL B487 241	R. Barate *et al.*	(ALEPH Collab.)
ABBIENDI	99E	EPJ C7 407	G. Abbiendi *et al.*	(OPAL Collab.)
ABBIENDI	99O	PL B464 311	G. Abbiendi *et al.*	(OPAL Collab.)
ABBOTT	99B	PRL 82 2244	B. Abbott *et al.*	(D0 Collab.)
ABREU	99P	PL B458 431	P. Abreu *et al.*	(DELPHI Collab.)
ABBOTT	98	PRL 80 442	B. Abbott *et al.*	(D0 Collab.)
ACKERSTAFF	98S	EPJ C5 19	K. Ackerstaff *et al.*	(OPAL Collab.)
ACKERSTAFF	98Y	PL B437 218	K. Ackerstaff *et al.*	(OPAL Collab.)
GONZALEZ...	98B	PR D57 7045	M.C. Gonzalez-Garcia, S.M. Lietti, S.F. Novaes	
KRAWCZYK	97	PR D55 6968	M. Krawczyk, J. Zochowski	(WARS)
ALEXANDER	96H	ZPHY C71 1	G. Alexander *et al.*	(OPAL Collab.)
ABREU	95H	ZPHY C67 69	P. Abreu *et al.*	(DELPHI Collab.)
BALEST	95	PR D51 2053	R. Balest *et al.*	(CLEO Collab.)
PICH	92	NP B388 31	A. Pich, J. Prades, P. Yepes	(CERN, CPPM)
ANTREASYAN	90C	PL B251 204	D. Antreasyan *et al.*	(Crystal Ball Collab.)