# New Heavy Bosons (*W*′, *Z*′, leptoquarks, etc.), Searches for

We list here various limits on charged and neutral heavy vector bosons (other than W's and Z's), heavy scalar bosons (other than Higgs bosons), vector or scalar leptoquarks, and axigluons. The latest unpublished results are described in "W' Searches" and "Z'Searches" reviews. For recent searches on scalar bosons which could be identified as Higgs bosons, see the listings in the Higgs boson section.

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## MASS LIMITS for W' (Heavy Charged Vector Boson Other Than W) in Hadron Collider Experiments

Couplings of W' to quarks and leptons are taken to be identical with those of W. The following limits are obtained from  $p\overline{p}$  or  $pp \rightarrow W'X$  with W' decaying to the mode indicated in the comments. New decay channels (e.g.,  $W' \rightarrow WZ$ ) are assumed to be suppressed. The most recent preliminary results can be found in the "W'-boson searches" review above.

VALUE (GeV)	<u>CL%</u>	DOCUMENT ID TECN COMMENT
>2220	95	<sup>1</sup> AABOUD 17B ATLS $W' \rightarrow hW$
>2600	95	<sup>2</sup> AABOUD 16AE ATLS $W' \rightarrow WZ$
>4070	95	<sup>3</sup> AABOUD 16V ATLS $W' \rightarrow e\nu, \mu\nu$
>1810	95	<sup>4</sup> AAD 16R ATLS $W' \rightarrow WZ$
>2600	95	<sup>5</sup> AAD 16s ATLS $W' \rightarrow q \overline{q}$
>2150	95	<sup>6</sup> KHACHATRY16AO CMS $W' \rightarrow t b$
none 1000–1600	95	<sup>7</sup> KHACHATRY16AP CMS $W' \rightarrow hW$
none 800–1500	95	<sup>8</sup> KHACHATRY16BD CMS $W' \rightarrow hW \rightarrow b\overline{b}\ell\nu$
none 1500–2600	95	<sup>9</sup> KHACHATRY16κ CMS $W'  ightarrow q \overline{q}$
none 500–1600	95	$^{10}$ KHACHATRY16L CMS $W'  ightarrow q \overline{q}$
none 300–2700	95	$^{11}$ KHACHATRY160 CMS $W'  ightarrow  au  u$
none 400–1590	95	<sup>12</sup> AAD 15AU ATLS $W' \rightarrow WZ$
none 1500–1760	95	<sup>13</sup> AAD 15AV ATLS $W' \rightarrow t b$
none 300–1490	95	<sup>14</sup> AAD 15AZ ATLS $W' \rightarrow WZ$
none 1300–1500	95	<sup>15</sup> AAD 15CP ATLS $W' \rightarrow WZ$
none 500–1920	95	$^{16}$ AAD 15R ATLS $W'  ightarrow t b$
none 800–2450	95	$17$ AAD 15V ATLS $W' \rightarrow q \overline{q}$
>1470	95	<sup>18</sup> KHACHATRY15C CMS $W' \rightarrow WZ$
>3710	95	$^{19}$ KHACHATRY15T $$ CMS $$ $$ $W^\prime  ightarrow  e  u,  \mu  u$
none 1000–3010	95	<sup>20</sup> KHACHATRY140 CMS $W' \rightarrow N\ell \rightarrow \ell\ell j j$
• • • We do not use the	e followir	ng data for averages, fits, limits, etc. $ullet$ $ullet$
• • • We do not use the	e followir	and data for averages, fits, limits, etc. $\bullet \bullet \bullet$ <sup>21</sup> AAD 15BB ATLS $W' \to Wh$
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none 300–880 none 1200–1900 and	95	<sup>21</sup> AAD 15BB ATLS $W' \rightarrow Wh$ <sup>22</sup> AALTONEN 15C CDF $W' \rightarrow tb$ <sup>23</sup> KHACHATRY15V CMS $W' \rightarrow q\overline{q}$ AAD 14AI ATLS $W' \rightarrow e\nu, \mu\nu$
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none 200–1143	95	$\begin{array}{cccccccccccccccccccccccccccccccccccc$
>1120	95	AALTONEN 11C CDF $W' ightarrow e u$
none 180–690	95	<sup>41</sup> ABAZOV 11H D0 $W' \rightarrow WZ$
none 600–863	95	$^{42}$ ABAZOV 11L D0 $W' \rightarrow t b$
none 285–516	95	<sup>43</sup> AALTONEN 10N CDF $W' \rightarrow WZ$
none 280–840	95	<sup>44</sup> AALTONEN 09AC CDF $W' \rightarrow q \overline{q}$
>1000	95	ABAZOV 08C D0 $W'  ightarrow e  u$
none 300–800	95	ABAZOV 04C D0 $W'  ightarrow q \overline{q}$
none 225–536	95	<sup>45</sup> ACOSTA 03B CDF $W' \rightarrow tb$
none 200–480	95	<sup>46</sup> AFFOLDER 02C CDF $W' \rightarrow WZ$
> 786	95	$^{47}$ AFFOLDER 011 CDF $W^\prime  ightarrow  e  u$ , $\mu  u$
none 300–420	95	<sup>48</sup> ABE 97G CDF $W' \rightarrow q \overline{q}$
> 720	95	<sup>49</sup> ABACHI 96C D0 $W' \rightarrow e \nu$
> 610	95	<sup>50</sup> ABACHI 95E D0 $W'  ightarrow e  u$ , $ au  u$
none 260–600	95	<sup>51</sup> RIZZO 93 RVUE $W' \rightarrow q \overline{q}$
4		

<sup>1</sup>AABOUD 17B search for resonances decaying to hW ( $h \rightarrow b\overline{b}$ ,  $c\overline{c}$ ;  $W \rightarrow \ell\nu$ ) in pp collisions at  $\sqrt{s} = 13$  TeV. The quoted limit is for heavy-vector-triplet W' with  $g_V = 3$ . The limit becomes  $M_{W'} > 1750$  GeV for  $g_V = 1$ . If we assume  $M_{W'} = M_{Z'}$ , the limit increases  $M_{W'} > 2310$  GeV and  $M_{W'} > 1730$  GeV for  $g_V = 3$  and  $g_V = 1$ , respectively. See their Fig.3 for limits on  $\sigma \cdot B$ .

<sup>2</sup> AABOUD 16AE search for resonances decaying to VV (V = W or Z) in pp collisions at  $\sqrt{s} = 13$  TeV. Results from  $\nu \nu qq$ ,  $\nu \ell qq$ ,  $\ell \ell qq$  and qqqq final states are combined. The quoted limit is for a heavy-vector-triplet W' with  $g_V = 3$  and  $M_{W'} = M_{Z'}$ .

<sup>3</sup>AABOUD 16V limit is for W' with SM-like coupling using pp collisions at  $\sqrt{s} = 13$  TeV. The bosonic decays of W' and the interference with SM W process are neglected.

<sup>4</sup> AAD 16R search for  $W' \to WZ$  in pp collisions at  $\sqrt{s} = 8$  TeV.  $\ell \nu \ell' \ell'$ ,  $\ell \ell q \overline{q}$ ,  $\ell \nu q \overline{q}$ , and all hadronic channels are combined. The quoted limit assumes  $g_{W'WZ}/g_{WWZ} = (M_W/M_{W'})^2$ .

<sup>5</sup> AAD 16S search for a new resonance decaying to dijets in pp collisions at  $\sqrt{s} = 13$  TeV. The limit quoted above is for a W' having SM-like couplings to quarks.

<sup>6</sup> KHACHATRYAN 16A0 limit is for a SM-like right-handed W' using pp collisions at  $\sqrt{s} = 8$  TeV. The quoted limit combines  $t \rightarrow qqb$  and  $t \rightarrow \ell \nu b$  events.

<sup>7</sup> KHACHATRYAN 16AP search for a resonance decaying to hW in pp collisions at  $\sqrt{s} = 8$  TeV. Both h and W are assumed to decay to fat jets. The quoted limit is for heavy-vector-triplet W' with  $g_V = 3$ .

<sup>8</sup>KHACHATRYAN 16BD search for resonance decaying to hW in pp collisions at  $\sqrt{s} = 8$  TeV. The quoted limit is for heavy-vector-triplet (HVT) W' with  $g_V = 3$ . The HVT model  $m_{W'} = m_{Z'} > 1.8$  TeV is also obtained by combining  $W'/Z' \rightarrow Wh/Zh \rightarrow \ell\nu bb$ ,  $qq\tau\tau$ , qqbb, and qqqqqq channels.

<sup>9</sup>KHACHATRYAN 16K search for resonances decaying to dijets in pp collisions at  $\sqrt{s} = 2$  13 TeV.

<sup>10</sup> KHACHATRYAN 16L search for resonances decaying to dijets in pp collisions at  $\sqrt{s}$ = 8 TeV with the data scouting technique, increasing the sensitivity to the low mass resonances.

<sup>11</sup> KHACHATRYAN 160 limit is for W' having universal couplings. Interferences with the SM amplitudes are assumed to be absent.

- <sup>12</sup> AAD 15AU search for W' decaying into the WZ final state with  $W \rightarrow q\overline{q}', Z \rightarrow \ell^+ \ell^-$  using pp collisions at  $\sqrt{s} = 8$  TeV. The quoted limit assumes  $g_{W'WZ}/g_{WWZ} = (M_W/M_{W'})^2$ .
- <sup>13</sup>AAD 15AV limit is for a SM like right-handed W' using pp collisions at  $\sqrt{s} = 8$  TeV.  $W' \rightarrow \ell \nu$  decay is assumed to be forbidden.
- <sup>14</sup> AAD 15AZ search for W' decaying into the WZ final state with  $W \rightarrow \ell \nu$ ,  $Z \rightarrow q \overline{q}$  using pp collisions at  $\sqrt{s} = 8$  TeV. The quoted limit assumes  $g_{W'WZ}/g_{WWZ} = (M_W/M_{W'})^2$ .
- <sup>15</sup> AAD 15CP search for W' decaying into the WZ final state with  $W \rightarrow q\overline{q}, Z \rightarrow q\overline{q}$ using pp collisions at  $\sqrt{s} = 8$  TeV. The quoted limit assumes  $g_{W'WZ}/g_{WWZ} = (M_W/M_{W'})^2$ .
- <sup>16</sup> AAD 15R limit is for a SM like right-handed W' using pp collisions at  $\sqrt{s} = 8$  TeV.  $W' \rightarrow \ell \nu$  decay is assumed to be forbidden.
- <sup>17</sup> AAD 15V search for new resonance decaying to dijets in pp collisions at  $\sqrt{s} = 8$  TeV.
- <sup>18</sup> KHACHATRYAN 15C search for W' decaying via WZ to fully leptonic final states using pp collisions at  $\sqrt{s}=8$  TeV. The quoted limit assumes  $g_{W'WZ}/g_{WWZ} = M_W M_Z/M_{W'}^2$ .
- <sup>19</sup> KHACHATRYAN 15T limit is for W' with SM-like coupling which interferes the SM W boson constructively using pp collisions at  $\sqrt{s} = 8$  TeV. For W' without interference, the limit becomes > 3280 GeV.
- the limit becomes > 3280 GeV. <sup>20</sup> KHACHATRYAN 140 search for right-handed  $W_R$  in pp collisions at  $\sqrt{s} = 8$  TeV.  $W_R$ is assumed to decay into  $\ell$  and hypothetical heavy neutrino N, with N decaying into  $\ell jj$ . The quoted limit is for  $M_{\nu_{eR}} = M_{\nu_{\mu R}} = M_{W_R}/2$ . See their Fig. 3 and Fig. 5 for excluded regions in the  $M_{W_R} - M_{\nu}$  plane.
- <sup>21</sup> AAD 15BB search for W' decaying into Wh with  $W \rightarrow \ell \nu$ ,  $h \rightarrow b\overline{b}$ . See their Fig. 4 for the exclusion limits in the heavy vector triplet benchmark model parameter space.
- <sup>22</sup> AALTONEN 15C limit is for a SM-like right-handed W' assuming  $W' \rightarrow \ell \nu$  decays are forbidden, using  $p\overline{p}$  collisions at  $\sqrt{s}=1.96$  TeV. See their Fig. 3 for limit on  $g_{W'}/g_W$ .
- <sup>23</sup>KHACHATRYAN 15V search new resonance decaying to dijets in pp collisions at  $\sqrt{s} =$  . 8 TeV.
- <sup>24</sup> AAD 14AT search for a narrow charged vector boson decaying to  $W\gamma$ . See their Fig. 3a for the exclusion limit in  $m_{W'} \sigma B$  plane.
- <sup>25</sup> AAD 14S search for W' decaying into the WZ final state with  $W \rightarrow \ell \nu, Z \rightarrow \ell \ell$ using pp collisions at  $\sqrt{s}=8$  TeV. The quoted limit assumes  $g_{W'WZ}/g_{WWZ} = (M_W/M_{W'})^2$ .
- <sup>26</sup> KHACHATRYAN 14 search for W' decaying into WZ final state with  $W \rightarrow q \overline{q}, Z \rightarrow q \overline{q}$  using p p collisions at  $\sqrt{s}=8$  TeV. The quoted limit assumes  $g_{W'WZ}/g_{WWZ} = (M_W/M_{W'})^2$ .
- <sup>27</sup> KHACHATRYAN 14A search for W' decaying into the WZ final state with  $W \rightarrow \ell \nu$ ,  $Z \rightarrow q \overline{q}$ , or  $W \rightarrow q \overline{q}$ ,  $Z \rightarrow \ell \ell$ . pp collisions data at  $\sqrt{s}=8$  TeV are used for the search. See their Fig. 13 for the exclusion limit on the number of events in the mass-width plane.
- <sup>28</sup> AAD 13AO search for W' decaying into the WZ final state with  $W \rightarrow \ell \nu$ ,  $Z \rightarrow 2j$  using pp collisions at  $\sqrt{s}=7$  TeV. The quoted limit assumes  $g_{W'WZ}/g_{WWZ} = (M_W/M_{W'})^2$ .
- <sup>29</sup> CHATRCHYAN 13AJ search for resonances decaying to WZ pair, using the hadronic decay modes of W and Z, in pp collisions at  $\sqrt{s}=7$  TeV. See their Fig. 7 for the limit on the cross section.

- $^{30}$  CHATRCHYAN 13AQ limit is for W' with SM-like coupling which interferes with the SM W boson using pp collisions at  $\sqrt{s}=7$  TeV.
- <sup>31</sup>CHATRCHYAN 13E limit is for W' with SM-like coupling which intereferes with the SM W boson using pp collisions at  $\sqrt{s}=7$  TeV. For W' with right-handed coupling, the bound becomes >1850 GeV (>1910 GeV) if W' decays to both leptons and guarks (only to quarks). If both left- and right-handed couplings are present, the limit becomes >1640 GeV. <sup>32</sup> CHATRCHYAN 13U search for W' decaying to the WZ final state, with W decaying
- into jets, in pp collisions at  $\sqrt{s}=7$  TeV. The quoted limit assumes  $g_{W'WZ}/g_{WWZ}$  $= (M_W/M_{W'})^2.$
- <sup>33</sup> The AAD 12AV quoted limit is for a SM-like right-handed W' using pp collisions at  $\sqrt{s}{=}7$  TeV.  $W' \rightarrow \ell \nu$  decay is assumed to be forbidden. <sup>34</sup> AAD 12BB use pp collisions data at  $\sqrt{s}{=}7$  TeV. The quoted limit assumes
- $g_{W'WZ}/g_{WWZ} = (M_W/M_{W'})^2.$
- <sup>35</sup>AAD 12CK search for  $pp \rightarrow tW'$ ,  $W' \rightarrow \overline{t}q$  events in pp collisions. See their Fig. 5 for the limit on  $\sigma \cdot B$ . 36 AAD 12CR use *pp* collisions at  $\sqrt{s}=7$  TeV.
- $^{37}$  AAD 12M search for right-handed  $W_R$  in pp collisions at  $\sqrt{s}=$  7 TeV.  $W_R$  is assumed to decay into  $\ell$  and hypothetical heavy neutrino N, with N decaying into  $\ell jj$ . See their Fig. 4 for the limit in the  $m_N - m_{W'}$  plane.
- <sup>38</sup>AALTONEN 12N search for  $p\overline{p} \rightarrow tW'$ ,  $W' \rightarrow \overline{t}d$  events in  $p\overline{p}$  collisions. See their Fig. 3 for the limit on  $\sigma \cdot B$ .
- <sup>39</sup> CHATRCHYAN 12AR search for  $pp \rightarrow tW'$ ,  $W' \rightarrow \overline{t}d$  events in pp collisions. See their Fig. 2 for the limit on  $\sigma \cdot B$ .
- <sup>40</sup> CHATRCHYAN 12BG search for right-handed  $W_R$  in pp collisions  $\sqrt{s} = 7$  TeV.  $W_R$  is assumed to decay into  $\ell$  and hypothetical heavy neutrino N, with N decaying into  $\ell j j$ . See their Fig. 3 for the limit in the  $m_N - m_{W'}$  plane.
- <sup>41</sup>ABAZOV 11H use data from  $p \overline{p}$  collisions at  $\sqrt{s}$ =1.96 TeV. The quoted limit is obtained assuming W'WZ coupling strength is the same as the ordinary WWZ coupling strength in the Standard Model.  $^{42}\,\rm ABAZOV\,\,11L$  limit is for W' with SM-like coupling which interferes with the SM W
- boson, using  $p\overline{p}$  collisions at  $\sqrt{s}=1.96$  TeV. For W' with right-handed coupling, the bound becomes >885 GeV (>890 GeV) if W' decays to both leptons and quarks (only to quarks). If both left- and right-handed couplings present, the limit becomes >916 GeV.
- <sup>43</sup> AALTONEN 10N use  $p\overline{p}$  collision data at  $\sqrt{s}$ =1.96 TeV. The quoted limit assumes  $g_{W'WZ}/g_{WWZ} = (M_W/M_{W'})^2$ . See their Fig. 4 for limits in mass-coupling plane.
- <sup>44</sup>AALTONEN 09AC search for new particle decaying to dijets using  $p\overline{p}$  collisions at  $\sqrt{s}{=}1.96$  TeV.  $^{45}$  The ACOSTA 03B quoted limit is for  $M_{W'} \gg M_{\nu_R}$ , using  $p\overline{p}$  collisions at  $\sqrt{s}{=}1.8$  TeV.
- For  $M_{W'} < M_{\nu_P}$ ,  $M_{W'}$  between 225 and 566 GeV is excluded.
- $^{46}$  The quoted limit is obtained assuming W'WZ coupling strength is the same as the ordinary W W Z coupling strength in the Standard Model, using  $p\overline{p}$  collisions at  $\sqrt{s}=1.8$ TeV. See their Fig. 2 for the limits on the production cross sections as a function of the W' width.
- 47 AFFOLDER 011 combine a new bound on  $W' \rightarrow e\nu$  of 754 GeV, using  $p\overline{p}$  collisions at  $\sqrt{s}$ =1.8 TeV, with the bound of ABE 00 on  $W' \rightarrow \mu \nu$  to obtain quoted bound.
- <sup>48</sup>ÅBE 97G search for new particle decaying to dijets using  $p\overline{p}$  collisions at  $\sqrt{s}=1.8$  TeV.
- $^{49}$  For bounds on  $W_R$  with nonzero right-handed mass, see Fig. 5 from ABACHI 96C.
- $^{50}$  ABACHI 95E assume that the decay W' 
  ightarrow WZ is suppressed and that the neutrino from W' decay is stable and has a mass significantly less  $m_{W'}$ .
- <sup>51</sup> RIZZO 93 analyses CDF limit on possible two-jet resonances. The limit is sensitive to the inclusion of the assumed K factor.

#### W<sub>R</sub> (Right-Handed W Boson) MASS LIMITS

Assuming a light right-handed neutrino, except for BEALL 82, LANGACKER 89B, and COLANGELO 91.  $g_R = g_L$  assumed. [Limits in the section MASS LIMITS for W' below are also valid for  $W_R$  if  $m_{\nu_R} \ll m_{W_R}$ .] Some limits assume manifest left-right symmetry, *i.e.*, the equality of left- and right Cabibbo-Kobayashi-Maskawa matrices. For a comprehensive review, see LANGACKER 89B. Limits on the  $W_L$ - $W_R$  mixing angle  $\zeta$  are found in the next section. Values in brackets are from cosmological and astrophysical considerations and assume a light right-handed neutrino.

and astrophysic VALUE (GeV)	al consid	erations and assum <u>DOCUMENT ID</u>	ealı	ght right TECN	-handed neutrino.
> 592	90	<sup>1</sup> BUENO	11		$\mu$ decay
> 715	90 90	<sup>2</sup> CZAKON	99		$\mu$ decay Electroweak
• • • We do not use					
> 235	90	<sup>3</sup> PRIEELS	14	PIE3	$\mu$ decay
> 245	90	<sup>4</sup> WAUTERS	10	CNTR	<sup>60</sup> Co $\beta$ decay
>2500		<sup>5</sup> ZHANG	08	THEO	${}^{m}\kappa_{L}^{0}{}^{-m}\kappa_{S}^{0}$
> 180	90	<sup>6</sup> MELCONIAN	07	CNTR	$37 \text{ K}^{5} \beta^{+} \text{ decay}$
> 290.7	90	<sup>7</sup> SCHUMANN	07	CNTR	Polarized neutron decay
[> 3300]	95	<sup>8</sup> CYBURT	05	COSM	Nucleosynthesis; light $\nu_R$
> 310	90	<sup>9</sup> THOMAS	01	CNTR	
> 137	95	<sup>10</sup> ACKERSTAFF	<b>99</b> D	OPAL	au decay
>1400	68	<sup>11</sup> BARENBOIM	98	RVUE	Electroweak, $Z$ - $Z'$ mixing
> 549	68	<sup>12</sup> BARENBOIM	97	RVUE	$\mu$ decay
> 220	95	<sup>13</sup> STAHL	97	RVUE	au decay
> 220	90	<sup>14</sup> ALLET	96	CNTR	$\beta^+$ decay
> 281	90	<sup>15</sup> KUZNETSOV	95	CNTR	Polarized neutron decay
> 282	90	<sup>16</sup> KUZNETSOV	<b>94</b> B	CNTR	Polarized neutron decay
> 439	90	<sup>17</sup> BHATTACH	93	RVUE	Z-Z' mixing
> 250	90	<sup>18</sup> SEVERIJNS	93	CNTR	$\beta^+$ decay
		<sup>19</sup> IMAZATO	92	CNTR	$\kappa^+$ decay
> 475	90	<sup>20</sup> POLAK	<b>92</b> B	RVUE	$\mu$ decay
> 240	90	<sup>21</sup> AQUINO	91	RVUE	Neutron decay
> 496	90	<sup>21</sup> AQUINO	91	RVUE	Neutron and muon decay
> 700		<sup>22</sup> COLANGELO	91	THEO	${}^{m}\kappa^{0}_{L} - {}^{m}\kappa^{0}_{S}$
> 477	90	<sup>23</sup> POLAK	91	RVUE	$\mu$ decay
[none 540-23000]		<sup>24</sup> BARBIERI	<b>89</b> B	ASTR	SN 1987A; light $\nu_R$
> 300	90	<sup>25</sup> LANGACKER	<b>89</b> B	RVUE	General
> 160	90	<sup>26</sup> BALKE	88	CNTR	$\mu  ightarrow e  u \overline{ u}$
> 406	90	<sup>27</sup> Jodidio	86	ELEC	Any ζ
> 482	90	<sup>27</sup> JODIDIO	86	ELEC	$\zeta = 0$
> 800		MOHAPATRA	86	RVUE	$SU(2)_I \times SU(2)_R \times U(1)$
> 400	95	<sup>28</sup> STOKER	85	ELEC	Any $\zeta$
> 475	95	<sup>28</sup> STOKER	85	ELEC	ζ <0.041
		<sup>29</sup> BERGSMA	83	CHRM	$ u_{\mu} e \rightarrow \mu \nu_{e}$
> 380	90	<sup>30</sup> CARR	83	ELEC	$\mu^+$ decay
>1600		<sup>31</sup> BEALL	82	THEO	${}^{m}\kappa_{L}^{0}-{}^{m}\kappa_{S}^{0}$
					ni ns

<sup>1</sup> The quoted limit is for manifest left-right symmetric model.

- <sup>2</sup>CZAKON 99 perform a simultaneous fit to charged and neutral sectors.
- <sup>3</sup>PRIEELS 14 limit is from  $\mu^+ \rightarrow e^+ \nu \overline{\nu}$  decay parameter  $\xi''$ , which is determined by the positron polarization measurement.
- <sup>4</sup>WAUTERS 10 limit is from a measurement of the asymmetry parameter of polarized  $^{60}$ Co  $\beta$  decays. The listed limit assumes no mixing.
- <sup>5</sup> ZHANG 08 limit uses a lattice QCD calculation of the relevant hadronic matrix elements, while BEALL 82 limit used the vacuum saturation approximation.
- <sup>6</sup> MELCONIAN 07 measure the neutrino angular asymmetry in  $\beta^+$ -decays of polarized <sup>37</sup>K, stored in a magneto-optical trap. Result is consistent with SM prediction and does not constrain the  $W_L W_R$  mixing angle appreciably.
- <sup>7</sup>SCHUMANN 07 limit is from measurements of the asymmetry  $\langle \vec{p}_{\nu} \cdot \sigma_n \rangle$  in the  $\beta$  decay of polarized neutrons. Zero mixing is assumed.
- <sup>8</sup> CYBURT 05 limit follows by requiring that three light  $\nu_R$ 's decouple when  $T_{dec} > 140$  MeV. For different  $T_{dec}$ , the bound becomes  $M_{W_R} > 3.3$  TeV  $(T_{dec} / 140 \text{ MeV})^{3/4}$ .
- <sup>9</sup>THOMAS 01 limit is from measurement of  $\beta^+$  polarization in decay of polarized <sup>12</sup>N. The listed limit assumes no mixing.
- $^{10}\,\rm ACKERSTAFF$  99D limit is from  $\tau$  decay parameters. Limit increase to 145 GeV for zero mixing.
- <sup>11</sup>BARENBOIM 98 assumes minimal left-right model with Higgs of SU(2)<sub>R</sub> in SU(2)<sub>L</sub> doublet. For Higgs in SU(2)<sub>L</sub> triplet,  $m_{W_R} > 1100$  GeV. Bound calculated from effect of corresponding  $Z_{LR}$  on electroweak data through  $Z Z_{LR}$  mixing.
- <sup>12</sup> The quoted limit is from  $\mu$  decay parameters. BARENBOIM 97 also evaluate limit from  $K_L$ - $K_S$  mass difference.
- <sup>13</sup>STAHL 97 limit is from fit to  $\tau$ -decay parameters.
- $^{14}$  ALLET 96 measured polarization-asymmetry correlation in  $^{12}{\rm N}\beta^+$  decay. The listed limit assumes zero *L-R* mixing.
- <sup>15</sup> KUZNETSOV 95 limit is from measurements of the asymmetry  $\langle \vec{p}_{\nu} \cdot \sigma_n \rangle$  in the  $\beta$  decay of polarized neutrons. Zero mixing assumed. See also KUZNETSOV 94B.
- <sup>16</sup> KUZNETSOV 94B limit is from measurements of the asymmetry  $\langle \vec{p}_{\nu} \cdot \sigma_n \rangle$  in the  $\beta$  decay of polarized neutrons. Zero mixing assumed.
- <sup>17</sup> BHATTACHARYYA 93 uses Z Z' mixing limit from LEP '90 data, assuming a specific Higgs sector of SU(2)<sub>L</sub>×SU(2)<sub>R</sub>×U(1) gauge model. The limit is for  $m_t$ =200 GeV and slightly improves for smaller  $m_t$ .

 $^{18}\,\text{SEVERIJNS}$  93 measured polarization-asymmetry correlation in  $^{107}\ln\beta^+$  decay. The listed limit assumes zero L-R mixing. Value quoted here is from SEVERIJNS 94 erratum.

 $^{19}$  IMAZATO 92 measure positron asymmetry in  ${\cal K}^+$   $\rightarrow$   $~\mu^+ \, \nu_{\mu}$  decay and obtain

 $\xi P_{\mu} > 0.990$  (90% CL). If  $W_R$  couples to  $u\overline{s}$  with full weak strength ( $V_{us}^R = 1$ ), the result corresponds to  $m_{W_R} > 653$  GeV. See their Fig. 4 for  $m_{W_R}$  limits for general  $|V_{us}^R|^2 = 1 - |V_{ud}^R|^2$ .

- <sup>20</sup> POLAK 92B limit is from fit to muon decay parameters and is essentially determined by JODIDIO 86 data assuming  $\zeta$ =0. Supersedes POLAK 91.
- <sup>21</sup> AQUINO 91 limits obtained from neutron lifetime and asymmetries together with unitarity of the CKM matrix. Manifest left-right symmetry assumed. Stronger of the two limits also includes muon decay results.
- <sup>22</sup> COLANGELO 91 limit uses hadronic matrix elements evaluated by QCD sum rule and is less restrictive than BEALL 82 limit which uses vacuum saturation approximation. Manifest left-right symmetry assumed.
- <sup>23</sup> POLAK 91 limit is from fit to muon decay parameters and is essentially determined by JODIDIO 86 data assuming  $\zeta$ =0. Superseded by POLAK 92B.
- $^{24}$  BARBIERI 89B limit holds for  $m_{
  u_R} \leq 10$  MeV.

- <sup>25</sup> LANGACKER 89B limit is for any  $\nu_R$  mass (either Dirac or Majorana) and for a general class of right-handed quark mixing matrices.
- <sup>26</sup> BALKE 88 limit is for  $m_{\nu_{eR}} = 0$  and  $m_{\nu_{\mu R}} \leq 50$  MeV. Limits come from precise measurements of the muon decay asymmetry as a function of the positron energy.
- <sup>27</sup> JODIDIO 86 is the same TRIUMF experiment as STOKER 85 (and CARR 83); however, it uses a different technique. The results given here are combined results of the two techniques. The technique here involves precise measurement of the end-point  $e^+$ spectrum in the decay of the highly polarized  $\mu^+$ .
- <sup>28</sup> STOKER 85 is same TRIUMF experiment as CARR 83. Here they measure the decay  $e^+$  spectrum asymmetry above 46 MeV/c using a muon-spin-rotation technique. Assumed a light right-handed neutrino. Quoted limits are from combining with CARR 83.
- <sup>29</sup> BERGSMA 83 set limit  $m_{W_2}/m_{W_1}$  >1.9 at CL = 90%.
- <sup>30</sup> CARR 83 is TRIUMF experiment with a highly polarized  $\mu^+$  beam. Looked for deviation from V-A at the high momentum end of the decay  $e^+$  energy spectrum. Limit from previous world-average muon polarization parameter is  $m_{W_R}$  >240 GeV. Assumes a light right-handed neutrino.
- <sup>31</sup> BEALL 82 limit is obtained assuming that  $W_R$  contribution to  $K_L^0 K_S^0$  mass difference is smaller than the standard one, neglecting the top quark contributions. Manifest left-right symmetry assumed.

#### Limit on $W_L$ - $W_R$ Mixing Angle $\zeta$

Lighter mass eigenstate  $W_1 = W_L \cos \zeta - W_R \sin \zeta$ . Light  $\nu_R$  assumed unless noted. Values in brackets are from cosmological and astrophysical considerations.

VALUE	<u>CL%</u>	ł	DOCUMENT ID		TECN	COMMENT
• • • We do not use the	following	da	ta for averages,	, fits,	limits, e	tc. ● ● ●
-0.020 to 0.017 < 0.022	90 90		BUENO MACDONALD			$\mu  ightarrow e  u \overline{ u} \ \mu  ightarrow e  u \overline{ u}$
< 0.022	95	1	ACKERSTAFF	<b>99</b> D	OPAL	au decay
< 0.013 < 0.0333	90		CZAKON BARENBOIM	99 07		Electroweak $\mu$ decay
< 0.04	90	4	MISHRA			$\nu N$ scattering
-0.0006 to 0.0028 [none 0.00001-0.02]	90			91 89в	RVUE ASTR	SN 1987A
< 0.040	90	7	JODIDIO	86	ELEC	$\mu$ decay
-0.056 to 0.040	90	•	JODIDIO	86	ELEC	$\mu$ decay

<sup>1</sup>ACKERSTAFF 99D limit is from au decay parameters.

 $^2$ CZAKON 99 perform a simultaneous fit to charged and neutral sectors.

- <sup>3</sup> The quoted limit is from  $\mu$  decay parameters. BARENBOIM 97 also evaluate limit from  $K_L$ - $K_S$  mass difference.
- <sup>4</sup> MISHRA 92 limit is from the absence of extra large-x, large-y  $\overline{\nu}_{\mu} N \rightarrow \overline{\nu}_{\mu} X$  events at Tevatron, assuming left-handed  $\nu$  and right-handed  $\overline{\nu}$  in the neutrino beam. The result gives  $\zeta^2(1-2m_{W_1}^2/m_{W_2}^2) < 0.0015$ . The limit is independent of  $\nu_R$  mass.
- <sup>5</sup> AQUINO 91 limits obtained from neutron lifetime and asymmetries together with unitarity of the CKM matrix. Manifest left-right asymmetry is assumed.

 $^{6}$  BARBIERI 89B limit holds for  $m_{
u_R} \leq$  10 MeV.

<sup>7</sup> First JODIDIO 86 result assumes  $m_{W_R} = \infty$ , second is for unconstrained  $m_{W_R}$ .

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#### MASS LIMITS for Z' (Heavy Neutral Vector Boson Other Than Z)

## Limits for $Z'_{SM}$

 $Z'_{SM}$  is assumed to have couplings with quarks and leptons which are identical to those of Z, and decays only to known fermions. The most recent preliminary results can be found in the "Z'-boson searches" review above.

VALUE (GeV)	<u>CL%</u>	DOCUMENT ID	s rev	TECN	COMMENT
>3360	95	<sup>1</sup> AABOUD	<b>16</b> ∪	ATLS	pp; Z'_ $SM$ $ ightarrow$ $e^+e^-$ , $\mu^+\mu^-$
>2020	95	<sup>2</sup> AAD	15AM	1ATLS	$pp; Z'_{SM} \rightarrow \tau^+ \tau^-$
>2900	95	<sup>3</sup> KHACHATRY	.15AE	CMS	$pp; Z'_{SM} \rightarrow \tau^+ \tau^-$ $pp; Z'_{SM} \rightarrow e^+ e^-, \mu^+ \mu^-$
none 1200-1700	95	<sup>4</sup> KHACHATRY	.15v	CMS	$pp; Z'_{SM} \rightarrow q\overline{q}$
>2900	95	<sup>5</sup> AAD	14V	ATLS	$pp; Z'_{SM}^{M} \rightarrow q\overline{q}$ $pp; Z'_{SM} \rightarrow e^+e^-, \mu^+\mu^-$
$\bullet \bullet \bullet$ We do not	use the	following data for a	iverag	es, fits,	limits, etc. $\bullet \bullet \bullet$
>1900	95	<sup>6</sup> AABOUD	16AA	ATLS	pp; $Z'_{SM} \rightarrow \tau^+ \tau^-$
>1400	95	<sup>7</sup> AAD	13S	ATLS	$pp; Z'_{SM} \rightarrow \tau^+ \tau^-$
>1470	95	<sup>8</sup> CHATRCHYAN	13A	CMS	$pp; Z'_{SM} \rightarrow q\overline{q}$
>2590	95	<sup>9</sup> CHATRCHYAN	13AF	CMS	pp; $Z'_{SM} \rightarrow e^+e^-, \mu^+\mu^-$
>2220	95	<sup>10</sup> AAD	12CC	ATLS	$pp; Z'_{SM} \rightarrow \tau^+ \tau^-$ $pp; Z'_{SM} \rightarrow e^+ e^-, \mu^+ \mu^-$ $pp; Z'_{SM} \rightarrow e^+ e^-, \mu^+ \mu^-$ $pp; Z'_{SM} \rightarrow e^+ e^-, \mu^+ \mu^-$
>1400	95	<sup>11</sup> CHATRCHYAN	120	CMS	$pp; Z'_{CM} \rightarrow \tau^{+}\tau^{-}$
>1071	95	<sup>12</sup> AALTONEN	11	CDF	$p\overline{p}; Z'_{SM} \rightarrow \mu^+ \mu^-$
>1023	95	<sup>13</sup> ABAZOV	11A	D0	$p\overline{p}; Z'_{SM} \rightarrow \mu^+ \mu^-$ $p\overline{p}; Z'_{SM} \rightarrow e^+ e^-$
none 247–544	95	<sup>14</sup> AALTONEN	10N	CDF	$Z' \rightarrow W W$
none 320–740	95	<sup>15</sup> AALTONEN	<b>09</b> AC	CDF	$Z' \rightarrow q \overline{q}$
> 963	95	<sup>13</sup> AALTONEN	09T	CDF	$p \overline{p}, Z'_{SM}  ightarrow e^+ e^-$
>1403	95	<sup>16</sup> ERLER	09	RVUE	
>1305	95	<sup>17</sup> ABDALLAH	<b>06</b> C	DLPH	e <sup>+</sup> e <sup>-</sup>
> 399	95	<sup>18</sup> ACOSTA	<b>05</b> R	CDF	$ \overline{p} p: Z'_{SM} \to \tau^+ \tau^- $ $ p \overline{p}: Z'_{SM} \to q \overline{q} $
none 400–640	95	ABAZOV	04C	D0	$p\overline{p}: Z'_{SM} \rightarrow q\overline{q}$
>1018	95	<sup>19</sup> ABBIENDI	<b>0</b> 4G	OPAL	e <sup>+</sup> e <sup>-</sup>
> 670	95	<sup>20</sup> ABAZOV	<b>01</b> B	D0	$p \overline{p}, Z'_{SM} \rightarrow e^+ e^-$
>1500	95	<sup>21</sup> CHEUNG	<b>01</b> B	RVUE	Electroweak
> 710	95	<sup>22</sup> ABREU	<b>00</b> S	DLPH	$e^+e^-$
> 898	95	<sup>23</sup> BARATE	001	ALEP	$e^+e^-$
> 809	95	<sup>24</sup> ERLER	99	RVUE	
> 690	95	<sup>25</sup> ABE	<b>97</b> S	CDF	$p \overline{p}; Z'_{SM}  ightarrow e^+ e^-, \mu^+ \mu^-$
> 398	95	<sup>26</sup> VILAIN	<b>94</b> B	CHM2	$ u_{\mu} e  ightarrow  u_{\mu} e$ and $\overline{ u}_{\mu} e  ightarrow  u_{\mu} e$
> 237	90	<sup>27</sup> ALITTI	93	UA2	$p\overline{p}; Z'_{SM} \rightarrow q\overline{q}$
none 260–600	95	<sup>28</sup> RIZZO	93	RVUE	
> 426	90	<sup>29</sup> ABE	90F	VNS	$e^+e^-$
-					

<sup>1</sup>AABOUD 16U search for resonances decaying to  $\ell^+ \ell^-$  in *pp* collisions at  $\sqrt{s} =$  13 TeV. <sup>2</sup>AAD 15AM search for resonances decaying to  $\tau^+ \tau^-$  in pp collisions at  $\sqrt{s} = 8$  TeV.

<sup>3</sup>KHACHATRYAN 15AE search for resonances decaying to  $e^+e^-$ ,  $\mu^+\mu^-$  in *pp* collisions at  $\sqrt{s} = 8$  TeV.

<sup>4</sup>KHACHATRYAN 15V search for resonances decaying to dijets in pp collisions at  $\sqrt{s}$  = 8 TeV.

- <sup>5</sup>AAD 14V search for resonances decaying to  $e^+e^-$ ,  $\mu^+\mu^-$  in pp collisions at  $\sqrt{s}=8$ TeV.
- <sup>6</sup>AABOUD 16AA search for resonances decaying to  $\tau^+ \tau^-$  in *pp* collisions at  $\sqrt{s} = 13$ TeV.
- <sup>7</sup> AAD 13S search for resonances decaying to  $\tau^+ \tau^-$  in pp collisions at  $\sqrt{s} = 7$  TeV.
- <sup>8</sup>CHATRCHYAN 13A use *pp* collisions at  $\sqrt{s}=7$  TeV.
- <sup>9</sup>CHATRCHYAN 13AF search for resonances decaying to  $e^+e^-$ ,  $\mu^+\mu^-$  in pp collisions at  $\sqrt{s} = 7$  TeV and 8 TeV.
- $^{10}$  AAD 12CC search for resonances decaying to  $e^+\,e^-$  ,  $\mu^+\,\mu^-$  in  $p\,p$  collisions at  $\sqrt{s}=7$ TeV.
- <sup>11</sup>CHATRCHYAN 120 search for resonances decaying to  $\tau^+ \tau^-$  in pp collisions at  $\sqrt{s} =$ 7 TeV.
- <sup>12</sup>AALTONEN 11I search for resonances decaying to  $\mu^+\mu^-$  in  $p\overline{p}$  collisions at  $\sqrt{s} = 1.96$ TeV.
- <sup>13</sup>ABAZOV 11A, AALTONEN 09T, AALTONEN 07H, and ABULENCIA 06L search for resonances decaying to  $e^+e^-$  in  $p\overline{p}$  collisions at  $\sqrt{s} = 1.96$  TeV.
- <sup>14</sup> The quoted limit assumes  $g_{W,W,Z'}/g_{W,W,Z} = (M_W/M_{Z'})^2$ . See their Fig. 4 for limits in mass-coupling plane.
- <sup>15</sup> AALTONEN 09AC search for new particle decaying to dijets.
- <sup>16</sup> ERLER 09 give 95% CL limit on the Z-Z' mixing  $-0.0026 < \theta < 0.0006$ .
- $^{17}$  ABDALLAH 06C use data  $\sqrt{s} = 130$ –207 GeV.
- <sup>18</sup>ACOSTA 05R search for resonances decaying to tau lepton pairs in  $\overline{p}p$  collisions at  $\sqrt{s}$ = 1.96 TeV.
- <sup>19</sup>ABBIENDI 04G give 95% CL limit on Z-Z' mixing  $-0.00422 < \theta < 0.00091$ .  $\sqrt{s} = 91$
- to 207 GeV. 20 ABAZOV 01B search for resonances in  $p\overline{p} \rightarrow e^+e^-$  at  $\sqrt{s}$ =1.8 TeV. They find  $\sigma$ .  $B(Z' \rightarrow ee) < 0.06 \text{ pb for } M_{7'} > 500 \text{ GeV}.$
- $^{21}$  CHEUNG 01B limit is derived from bounds on contact interactions in a global electroweak analysis.
- <sup>22</sup>ABREU 00S uses LEP data at  $\sqrt{s}$ =90 to 189 GeV.
- $^{23}$  BARATE 001 search for deviations in cross section and asymmetries in  $e^+e^- \rightarrow$  fermions at  $\sqrt{s}=90$  to 183 GeV. Assume  $\theta=0$ . Bounds in the mass-mixing plane are shown in their Figure 18.
- $^{24}$  ERLER 99 give 90%CL limit on the Z-Z' mixing -0.0041 < heta < 0.0003.  $ho_0=1$  is assumed. 25 ABE 97s find  $\sigma(Z') \times B(e^+e^-, \mu^+\mu^-) <$  40 fb for  $m_{Z'} > 600$  GeV at  $\sqrt{s} = 1.8$  TeV.

- <sup>26</sup> VILAIN 94B assume  $m_t = 150$  GeV.
- $^{27}$  ALITTI 93 search for resonances in the two-jet invariant mass. The limit assumes B(Z' 
  ightarrow $q \overline{q}$  =0.7. See their Fig. 5 for limits in the  $m_{\tau'}$  -B( $q \overline{q}$ ) plane.
- <sup>28</sup> RIZZO 93 analyses CDF limit on possible two-jet resonances.
- $^{29}$  ABE 90F use data for R,  $R_{\ell\ell}$  , and  $A_{\ell\ell}$  . They fix  $m_W=80.49\pm0.43\pm0.24$  GeV and  $m_7 = 91.13 \pm 0.03$  GeV.

#### Limits for $Z_{IR}$

 $Z_{LR}$  is the extra neutral boson in left-right symmetric models.  $g_L = g_R$  is assumed unless noted. Values in parentheses assume stronger constraint on the Higgs sector, usually motivated by specific left-right symmetric models (see the Note on the W'). Values in brackets are from cosmological and astrophysical considerations and assume a light right-handed neutrino. Direct search bounds assume decays to Standard Model fermions only, unless noted.

VALUE (GeV)	CL%	DOCUMENT ID		TECN	COMMENT
>1162	95	<sup>1</sup> DEL-AGUILA			
> 630	95	<sup>2</sup> ABE	<b>97</b> S	CDF	$ ho  \overline{ ho};  Z'_{LR}  ightarrow  e^+ e^-,  \mu^+ \mu^-$
HTTP://PDG	.LBL.GC	)V Pag	ge 10		Created: 5/30/2017 17:23

• • • We do not use the following data for averages, fits, limits, etc. • • •

		•	-		
> 998	95	<sup>3</sup> ERLER	09	RVUE	Electroweak
> 600	95	SCHAEL	<b>07</b> A	ALEP	e <sup>+</sup> e <sup>-</sup>
> 455	95	<sup>4</sup> ABDALLAH	<b>06</b> C	DLPH	e <sup>+</sup> e <sup>-</sup>
> 518	95	<sup>5</sup> ABBIENDI	<b>0</b> 4G	OPAL	e <sup>+</sup> e <sup>-</sup>
> 860	95	<sup>6</sup> CHEUNG	<b>01</b> B	RVUE	Electroweak
> 380	95	<sup>7</sup> ABREU	<b>00</b> S	DLPH	e <sup>+</sup> e <sup>-</sup>
> 436	95	<sup>8</sup> BARATE	001	ALEP	Repl. by SCHAEL 07A
> 550	95	<sup>9</sup> CHAY	00	RVUE	Electroweak
		<sup>10</sup> ERLER	00	RVUE	Cs
		<sup>11</sup> CASALBUONI	99	RVUE	Cs
(> 1205)	90	<sup>12</sup> CZAKON	99	RVUE	Electroweak
> 564	95	<sup>13</sup> ERLER	99	RVUE	Electroweak
(> 1673)	95	<sup>14</sup> ERLER	99	RVUE	Electroweak
(> 1700)	68	<sup>15</sup> BARENBOIM	98	RVUE	Electroweak
> 244	95	<sup>16</sup> CONRAD	98	RVUE	$ u_{\mu} N$ scattering
> 253	95	<sup>17</sup> VILAIN	<b>94</b> B	CHM2	$\dot{ u}_{\mu} e  ightarrow   u_{\mu} e$ and $\overline{ u}_{\mu} e  ightarrow  \overline{ u}_{\mu} e$
none 200–600	95	<sup>18</sup> RIZZO	93	RVUE	$p\overline{p}; Z_{LR} \rightarrow q\overline{q}$
[> 2000]		WALKER	91		Nucleosynthesis; light $\nu_R$
none 200–500		<sup>19</sup> GRIFOLS	90		SN 1987A; light $\nu_R$
none 350–2400		<sup>20</sup> BARBIERI	<b>89</b> B		SN 1987A; light $\nu_R$
_					

<sup>1</sup>DEL-AGUILA 10 give 95% CL limit on the Z-Z' mixing  $-0.0012 < \theta < 0.0004$ .

<sup>2</sup>ABE 97S find  $\sigma(Z') \times B(e^+e^-, \mu^+\mu^-) < 40$  fb for  $m_{Z'} > 600$  GeV at  $\sqrt{s} = 1.8$  TeV.

- <sup>3</sup>ERLER 09 give 95% CL limit on the Z-Z' mixing  $-0.0013 < \theta < 0.0006$ .
- <sup>4</sup>ABDALLAH 06C give 95% CL limit  $|\theta| < 0.0028$ . See their Fig. 14 for limit contours in the mass-mixing plane.
- <sup>5</sup>ABBIENDI 04G give 95% CL limit on Z-Z' mixing  $-0.00098 < \theta < 0.00190$ . See their Fig. 20 for the limit contour in the mass-mixing plane.  $\sqrt{s} = 91$  to 207 GeV.
- <sup>6</sup>CHEUNG 01B limit is derived from bounds on contact interactions in a global electroweak analysis.
- <sup>7</sup>ABREU 00S give 95% CL limit on Z-Z' mixing  $|\theta| <$  0.0018. See their Fig. 6 for the limit contour in the mass-mixing plane.  $\sqrt{s}=90$  to 189 GeV.
- $^8$  BARATE 001 search for deviations in cross section and asymmetries in  $e^+e^- 
  ightarrow$  fermions at  $\sqrt{s}$ =90 to 183 GeV. Assume  $\theta$ =0. Bounds in the mass-mixing plane are shown in their Figure 18.
- <sup>9</sup> CHAY 00 also find  $-0.0003 < \theta < 0.0019$ . For  $g_R$  free,  $m_{Z'} > 430$  GeV.

 $^{10}$  ERLER 00 discuss the possibility that a discrepancy between the observed and predicted values of  $Q_W(Cs)$  is due to the exchange of Z'. The data are better described in a certain class of the Z' models including  $Z_{LR}$  and  $Z_{\gamma}$ .

- $^{11}$ CASALBUONI 99 discuss the discrepancy between the observed and predicted values of  $Q_W(Cs)$ . It is shown that the data are better described in a class of models including the  $Z_{LR}$  model.
- $^{12}$  CZAKON 99 perform a simultaneous fit to charged and neutral sectors. Assumes manifest left-right symmetric model. Finds  $|\theta| < 0.0042$ .

 $^{13}$  ERLER 99 give 90% CL limit on the Z-Z' mixing  $-0.0009 < \theta < 0.0017.$ 

- <sup>14</sup> ERLER 99 assumes 2 Higgs doublets, transforming as 10 of SO(10), embedded in  $E_6$ .
- <sup>15</sup> BARENBOIM 98 also gives 68% CL limits on the Z-Z' mixing  $-0.0005 < \theta < 0.0033$ . Assumes Higgs sector of minimal left-right model.

<sup>16</sup>CONRAD 98 limit is from measurements at CCFR, assuming no Z-Z' mixing.

- $^{17}\,{\sf VILAIN}$  94B assume  $m_t$  = 150 GeV and  $\theta{=}0.\,$  See Fig.2 for limit contours in the mass-mixing plane.
- <sup>18</sup> RIZZO 93 analyses CDF limit on possible two-jet resonances. <sup>19</sup> GRIFOLS 90 limit holds for  $m_{\nu_R} \lesssim 1$  MeV. A specific Higgs sector is assumed. See also GRIFOLS 90D, RIZZO 91.
- $^{20}$  BARBIERI 89B limit holds for  $m_{\nu_R} \leq 10$  MeV. Bounds depend on assumed supernova core temperature.

**Limits for**  $Z_{\chi}$   $Z_{\chi}$  is the extra neutral boson in SO(10)  $\rightarrow$  SU(5)  $\times$  U(1) $_{\chi}$ .  $g_{\chi} = e/\cos\theta_{W}$  is assumed unless otherwise stated. We list limits with the assumption  $\rho = 1$  but with the assumption  $\rho = 1$  but with the assumption of  $\chi$  is the Higgs sector. Values in parentheses assume stronger constraint on the Higgs sector motivated by superstring models. Values in brackets are from cosmological and astrophysical considerations and assume a light right-handed neutrino.

VALUE (GeV)	CL%	DOCUMENT ID		TECN	COMMENT
>3050	95	$^1$ AABOUD	<b>16</b> ∪	ATLS	pp; $Z'_{\chi}  ightarrow e^+ e^-$ , $\mu^+ \mu^-$
>2620	95	<sup>2</sup> AAD	14V	ATLS	pp, $Z_{\gamma}^{\prime}  ightarrow e^+ e^-$ , $\mu^+ \mu^-$
• • • We do not ι	ise the fo	ollowing data for ave	rages	, fits, lim	nits, etc. $\bullet \bullet \bullet$
>1970	95	<sup>3</sup> AAD	1200	ATLS	pp, $Z'_{\chi}  ightarrow e^+ e^-$ , $\mu^+ \mu^-$
> 930	95	<sup>4</sup> AALTONEN	11	CDF	$ ho \overline{ ho}; Z_{\chi}^{\prime}  ightarrow \mu^+ \mu^-$
> 903	95	<sup>5</sup> ABAZOV	11A	D0	$p \overline{p}, Z'_{\chi}^{\Lambda} \rightarrow e^+ e^-$
>1022	95	<sup>6</sup> DEL-AGUILA	10	RVUE	Electroweak
> 862	95	<sup>5</sup> AALTONEN	09⊤	CDF	$p \overline{p}, Z'_{\chi} \rightarrow e^+ e^-$
> 892	95	<sup>7</sup> AALTONEN	09V	CDF	Repl. by AALTONEN 11
>1141	95	<sup>8</sup> ERLER	09	RVUE	Electroweak
> 822	95	<sup>5</sup> AALTONEN	07H		Repl. by AALTONEN 09T
> 680	95	SCHAEL		ALEP	e <sup>+</sup> e <sup>-</sup>
> 545	95	<sup>9</sup> ABDALLAH	06C		e <sup>+</sup> e <sup>-</sup>
> 740		<sup>5</sup> ABULENCIA	06L	CDF	Repl. by AALTONEN 07H
> 690	95	<sup>10</sup> ABULENCIA	05A	CDF	p $\overline{p};~Z_{\chi}^{\prime} ightarrow~e^{+}e^{-},~\mu^{+}\mu^{-}$
> 781	95	<sup>11</sup> ABBIENDI	<b>0</b> 4G	OPAL	e <sup>+</sup> e <sup>-</sup>
>2100		<sup>12</sup> BARGER	<b>03</b> B	COSM	Nucleosynthesis; light $ u_R$
> 680	95	<sup>13</sup> CHEUNG	<b>01</b> B	RVUE	Electroweak
> 440	95	<sup>14</sup> ABREU	00S	DLPH	e <sup>+</sup> e <sup>-</sup>
> 533	95	<sup>15</sup> BARATE	001	ALEP	Repl. by SCHAEL 07A
> 554	95	<sup>16</sup> CHO	00	RVUE	Electroweak
		<sup>17</sup> ERLER	00	RVUE	Cs
		<sup>18</sup> ROSNER	00	RVUE	Cs
> 545	95	<sup>19</sup> ERLER	99	RVUE	Electroweak
(> 1368)	95	<sup>20</sup> ERLER	99	RVUE	
> 215	95	<sup>21</sup> CONRAD	98	RVUE	$ u_{\mu}$ N scattering
> 595	95	<sup>22</sup> ABE	97s	CDF	p $\overline{ ho}; Z'_{\chi}  ightarrow  e^+  e^-$ , $\mu^+  \mu^-$
> 190	95	<sup>23</sup> ARIMA	97	VNS	Bhabha scattering
> 262	95	<sup>24</sup> VILAIN	<b>94</b> B	CHM2	$ u_{\mu} e  ightarrow \  u_{\mu} e;  \overline{ u}_{\mu} e  ightarrow \ \overline{ u}_{\mu} e$
[>1470]		<sup>25</sup> FARAGGI	91		Nucleosynthesis; light $\nu_R$
> 231	90	<sup>26</sup> ABE	90F	VNS	e <sup>+</sup> e <sup>-</sup>
[> 1140]		<sup>27</sup> GONZALEZ	<b>90</b> D	COSM	Nucleosynthesis; light $ u_{R}$
[> 2100]		<sup>28</sup> GRIFOLS	90	ASTR	SN 1987A; light $\nu_R$

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- <sup>1</sup>AABOUD 16U search for resonances decaying to  $\ell^+ \ell^-$  in *pp* collisions at  $\sqrt{s} = 13$  TeV.
- <sup>2</sup>AAD 14V search for resonances decaying to  $e^+e^-$ ,  $\mu^+\mu^-$  in *pp* collisions at  $\sqrt{s} = 8$  TeV.
- <sup>3</sup>AAD 12CC search for resonances decaying to  $e^+e^-$ ,  $\mu^+\mu^-$  in *pp* collisions at  $\sqrt{s} = 7$  TeV.
- <sup>4</sup> AALTONEN 111 search for resonances decaying to  $\mu^+ \mu^-$  in  $p \overline{p}$  collisions at  $\sqrt{s} = 1.96$  reV.
- <sup>5</sup>ABAZOV 11A, AALTONEN 09T, AALTONEN 07H, and ABULENCIA 06L search for resonances decaying to  $e^+e^-$  in  $p\overline{p}$  collisions at  $\sqrt{s} = 1.96$  TeV.
- <sup>6</sup>DEL-AGUILA 10 give 95% CL limit on the Z-Z' mixing  $-0.0011 < \theta < 0.0007$ .
- <sup>7</sup>AALTONEN 09V search for resonances decaying to  $\mu^+\mu^-$  in  $p\overline{p}$  collisions at  $\sqrt{s} = 1.96$  TeV.
- <sup>8</sup> ERLER 09 give 95% CL limit on the Z-Z' mixing  $-0.0016 < \theta < 0.0006$ .
- $^9$  ABDALLAH 06C give 95% CL limit  $|\theta|<$  0.0031. See their Fig. 14 for limit contours in the mass-mixing plane.
- <sup>10</sup> ABULENCIA 05A search for resonances decaying to electron or muon pairs in  $p\overline{p}$  collisions at  $\sqrt{s} = 1.96$  TeV.
- <sup>11</sup> ABBIENDI 04G give 95% CL limit on Z-Z' mixing  $-0.00099 < \theta < 0.00194$ . See their Fig. 20 for the limit contour in the mass-mixing plane.  $\sqrt{s} = 91$  to 207 GeV.
- $^{12}$  BARGER 03B limit is from the nucleosynthesis bound on the effective number of light neutrino  $\delta N_{\nu}$  <1. The quark-hadron transition temperature  $T_{c}{=}150$  MeV is assumed. The limit with  $T_{c}{=}400$  MeV is >4300 GeV.
- <sup>13</sup> CHEUNG 01B limit is derived from bounds on contact interactions in a global electroweak analysis.
- <sup>14</sup>ABREU 00S give 95% CL limit on Z-Z' mixing  $|\theta| < 0.0017$ . See their Fig. 6 for the limit contour in the mass-mixing plane.  $\sqrt{s}=90$  to 189 GeV.
- <sup>15</sup> BARATE 00I search for deviations in cross section and asymmetries in  $e^+e^- \rightarrow$  fermions at  $\sqrt{s}=90$  to 183 GeV. Assume  $\theta=0$ . Bounds in the mass-mixing plane are shown in their Figure 18.
- <sup>16</sup> CHO 00 use various electroweak data to constrain Z' models assuming  $m_H$ =100 GeV. See Fig. 3 for limits in the mass-mixing plane.
- <sup>17</sup> ERLER 00 discuss the possibility that a discrepancy between the observed and predicted values of  $Q_W(Cs)$  is due to the exchange of Z'. The data are better described in a certain class of the Z' models including  $Z_{LR}$  and  $Z_{\gamma}$ .
- <sup>18</sup> ROSNER 00 discusses the possibility that a discrepancy between the observed and predicted values of  $Q_W(Cs)$  is due to the exchange of Z'. The data are better described in a certain class of the Z' models including  $Z_{\chi}$ .
- <sup>19</sup> ERLER 99 give 90% CL limit on the Z-Z' mixing  $-0.0020 < \theta < 0.0015$ .
- $^{20}$  ERLER 99 assumes 2 Higgs doublets, transforming as 10 of SO(10), embedded in  $E_6$ .
- $^{21}$ CONRAD 98 limit is from measurements at CCFR, assuming no Z-Z $^{\prime}$  mixing.
- <sup>22</sup> ABE 97S find  $\sigma(Z') \times B(e^+e^-, \mu^+\mu^-) < 40$  fb for  $m_{Z'} > 600$  GeV at  $\sqrt{s} = 1.8$  TeV.
- $^{23}Z$ -Z' mixing is assumed to be zero.  $\sqrt{s}$ = 57.77 GeV.
- $^{24}$  VILAIN 94B assume  $m_t = 150$  GeV and  $\theta{=}0.$  See Fig. 2 for limit contours in the mass-mixing plane.
- $^{25}$  FARAGGI 91 limit assumes the nucleosynthesis bound on the effective number of neutrinos  $\Delta N_{\nu}~<~0.5$  and is valid for  $m_{\nu_{P}}~<~1$  MeV.
- <sup>26</sup> ABE 90F use data for R,  $R_{\ell\ell}$ , and  $A_{\ell\ell}$ . ABE 90F fix  $m_W = 80.49 \pm 0.43 \pm 0.24$  GeV and  $m_Z = 91.13 \pm 0.03$  GeV.
- <sup>27</sup> Assumes the nucleosynthesis bound on the effective number of light neutrinos ( $\delta N_{\nu} < 1$ ) and that  $\nu_R$  is light ( $\lesssim 1$  MeV).
- $^{28}\,{\rm GRIFOLS}$  90 limit holds for  $m_{\nu_R}\,\lesssim$  1 MeV. See also GRIFOLS 90D, RIZZO 91.

#### Limits for $Z_{\psi}$

 $Z_{\psi}$  is the extra neutral boson in  $E_6 \rightarrow SO(10) \times U(1)_{\psi}$ .  $g_{\psi} = e/\cos\theta_W$  is assumed unless otherwise stated. We list limits with the assumption  $\rho = 1$  but with no further constraints on the Higgs sector. Values in brackets are from cosmological and astrophysical considerations and assume a light right-handed neutrino.

VALUE (GeV)	<u>CL%</u>		DOCUMENT ID		TECN	COMMENT
>2740	95	1	AABOUD	<b>16</b> U	ATLS	pp; Z'_{\psi} \rightarrow e^+e^-, \mu^+\mu^-
>2570	95	2	KHACHATRY	.15ae	CMS	pp; $Z'_{\prime\prime}$ $\rightarrow e^+e^-$ , $\mu^+\mu^-$
>2510	95	3	AAD	14V	ATLS	pp, $Z_{\psi}^{ m p}  ightarrow  e^+ e^-$ , $\mu^+ \mu^-$
>1100	95	4	CHATRCHYAN	120	CMS	$pp, Z'_{\psi} \rightarrow \tau^+ \tau^-$
• • • We do not us	e the fol	lowi	ng data for avei	rages,	fits, lim	1
>2260	95	5	CHATRCHYAN	13AF	CMS	pp, Z'_{\psi} \rightarrow e^+e^-, \mu^+\mu^-
>1790	95	6	AAD	12CC	ATLS	pp, $Z_{\psi}^{\not r}  ightarrow e^+ e^-$ , $\mu^+ \mu^-$
>2000	95	7	CHATRCHYAN	12M	CMS	Repl. by CHA-
> 917	95	8	AALTONEN	11	CDF	TRCHYAN 13AF $p\overline{p}; Z'_{\psi} \rightarrow \mu^+ \mu^-$
> 891	95	9	ABAZOV	11A	D0	$p \overline{p}, Z'_{\psi} \rightarrow e^+ e^-$
> 476	95	10	DEL-AGUILA	10	RVUE	Electroweak
> 851	95	9	AALTONEN	09т	CDF	$p \overline{p}, Z'_{\psi}  ightarrow e^+ e^-$
> 878	95	11	AALTONEN	09v	CDF	Repl. by AALTONEN 11
> 147	95	12	ERLER	09	RVUE	Electroweak
> 822	95	9	AALTONEN	07н	CDF	Repl. by AALTONEN 09⊤
> 410	95		SCHAEL	07A	ALEP	e <sup>+</sup> e <sup>-</sup>
> 475	95	13	ABDALLAH	<b>06</b> C	DLPH	e <sup>+</sup> e <sup>-</sup>
> 725		9	ABULENCIA	06L	CDF	Repl. by AALTONEN 07H
> 675	95	14	ABULENCIA	05A	CDF	Repl. by AALTONEN 11
> 366	95	15	ABBIENDI	<b>0</b> 4G	OPAL	and AALTONEN 09T $e^+e^-$
> 600	50		BARGER	<b>03</b> B	COSM	Nucleosynthesis; light $\nu_R$
> 350	95		ABREU	00s	DLPH	$e^+e^-$
> 294	95	18	BARATE	001	ALEP	Repl. by SCHAEL 07A
> 137	95	19	СНО	00	RVUE	Electroweak
> 146	95	20	ERLER	99	RVUE	Electroweak
> 54	95	21	CONRAD	98	RVUE	$ u_{\mu} N$ scattering
> 590	95		ABE	<b>97</b> S	CDF	$p \overline{p}; Z'_{\psi} \rightarrow e^+ e^-, \mu^+ \mu^-$
> 135	95	23	VILAIN	<b>94</b> B	CHM2	$ u_{\mu} e \stackrel{\varphi}{\rightarrow} \nu_{\mu} e; \overline{\nu}_{\mu} e \rightarrow \overline{\nu}_{\mu} e$
> 105	90	24	ABE	90F	VNS	$e^{r} e^{-r}$
[> 160]			GONZALEZ			Nucleosynthesis; light $ u_R$
[> 2000]			GRIFOLS		ASTR	SN 1987A; light $\nu_R$
	oarch far					$r = collisions at \sqrt{c} = 12 \text{ TeV}$

<sup>1</sup>AABOUD 16U search for resonances decaying to  $\ell^+ \ell^-$  in *pp* collisions at  $\sqrt{s} = 13$  TeV. <sup>2</sup>KHACHATRYAN 15AE search for resonances decaying to  $e^+e^-$ ,  $\mu^+\mu^-$  in pp collisions at  $\sqrt{s} = 8$  TeV.

<sup>3</sup>AAD 14V search for resonances decaying to  $e^+e^-$ ,  $\mu^+\mu^-$  in *pp* collisions at  $\sqrt{s} = 8$ TeV.

<sup>4</sup>CHATRCHYAN 120 search for resonances decaying to  $\tau^+ \tau^-$  in pp collisions at  $\sqrt{s} =$ 7 TeV. <sup>5</sup> CHATRCHYAN 13AF search for resonances decaying to  $e^+e^-$ ,  $\mu^+\mu^-$  in pp collisions

at  $\sqrt{s} = 7$  TeV and 8 TeV.

- <sup>6</sup> AAD 12CC search for resonances decaying to  $e^+e^-$ ,  $\mu^+\mu^-$  in pp collisions at  $\sqrt{s} = 7$  TeV.
- <sup>7</sup> CHATRCHYAN 12M search for resonances decaying to  $e^+e^-$  or  $\mu^+\mu^-$  in pp collisions at  $\sqrt{s} = 7$  TeV.
- <sup>8</sup> AALTONEN 111 search for resonances decaying to  $\mu^+\mu^-$  in  $p\overline{p}$  collisions at  $\sqrt{s} = 1.96$  TeV.
- <sup>9</sup>ABAZOV 11A, AALTONEN 09T, AALTONEN 07H, and ABULENCIA 06L search for resonances decaying to  $e^+e^-$  in  $p\overline{p}$  collisions at  $\sqrt{s} = 1.96$  TeV.
- $^{10}$  DEL-AGUILA 10 give 95% CL limit on the Z-Z' mixing  $-0.0019 < \theta < 0.0007.$
- <sup>11</sup>AALTONEN 09V search for resonances decaying to  $\mu^+\mu^-$  in  $p\overline{p}$  collisions at  $\sqrt{s} = 1.96$  TeV.
- <sup>12</sup> ERLER 09 give 95% CL limit on the Z-Z' mixing  $-0.0018 < \theta < 0.0009$ .
- $^{13}$  ABDALLAH 06C give 95% CL limit  $\left|\theta\right|<$  0.0027. See their Fig. 14 for limit contours in the mass-mixing plane.
- <sup>14</sup> ABULENCIA 05A search for resonances decaying to electron or muon pairs in  $p\overline{p}$  collisions at  $\sqrt{s} = 1.96$  TeV.
- <sup>15</sup> ABBIENDI 04G give 95% CL limit on Z-Z' mixing  $-0.00129 < \theta < 0.00258$ . See their Fig. 20 for the limit contour in the mass-mixing plane.  $\sqrt{s} = 91$  to 207 GeV.
- $^{16}$  BARGER 03B limit is from the nucleosynthesis bound on the effective number of light neutrino  $\delta N_{\nu}$  <1. The quark-hadron transition temperature  $T_{c}{=}150$  MeV is assumed. The limit with  $T_{c}{=}400$  MeV is  ${>}1100$  GeV.
- <sup>17</sup> ABREU 00S give 95% CL limit on Z-Z' mixing  $|\theta| < 0.0018$ . See their Fig. 6 for the limit contour in the mass-mixing plane.  $\sqrt{s}=90$  to 189 GeV.
- <sup>18</sup> BARATE 00I search for deviations in cross section and asymmetries in  $e^+e^- \rightarrow$  fermions at  $\sqrt{s}=90$  to 183 GeV. Assume  $\theta=0$ . Bounds in the mass-mixing plane are shown in their Figure 18.
- <sup>19</sup> CHO 00 use various electroweak data to constrain Z' models assuming  $m_H$ =100 GeV. See Fig. 3 for limits in the mass-mixing plane.
- $^{20}$  ERLER 99 give 90% CL limit on the Z-Z' mixing  $-0.0013 < \theta < 0.0024$ .
- $^{21}$ CONRAD 98 limit is from measurements at CCFR, assuming no Z-Z' mixing.
- <sup>22</sup> ABE 97S find  $\sigma(Z') \times B(e^+e^-, \mu^+\mu^-) <$  40 fb for  $m_{Z'} >$  600 GeV at  $\sqrt{s}$ = 1.8 TeV.
- <sup>23</sup> VILAIN 94B assume  $m_t = 150$  GeV and  $\theta = 0$ . See Fig. 2 for limit contours in the mass-mixing plane.
- <sup>24</sup> ABE 90F use data for R,  $R_{\ell\ell}$ , and  $A_{\ell\ell}$ . ABE 90F fix  $m_W = 80.49 \pm 0.43 \pm 0.24$  GeV and  $m_Z = 91.13 \pm 0.03$  GeV.
- $^{25}$  Assumes the nucleosynthesis bound on the effective number of light neutrinos ( $\delta N_{\nu} < 1$ ) and that  $\nu_R$  is light ( $\lesssim$  1 MeV).

 $^{26}\,{\rm GRIFOLS}$  90D limit holds for  $m_{\nu_R}\,\lesssim\,1$  MeV. See also RIZZO 91.

#### Limits for $Z_{\eta}$

 $Z_\eta$  is the extra neutral boson in E<sub>6</sub> models, corresponding to  $Q_\eta = \sqrt{3/8} \ Q_\chi - \sqrt{5/8} \ Q_\psi$ .  $g_\eta = e/\cos\theta_W$  is assumed unless otherwise stated. We list limits with the assumption  $\rho = 1$  but with no further constraints on the Higgs sector. Values in parentheses assume stronger constraint on the Higgs sector motivated by superstring models. Values in brackets are from cosmological and astrophysical considerations and assume a light right-handed neutrino.

VALUE (GeV)	CL%	DOCUMENT ID		TECN	COMMENT	
>2810	95	<sup>1</sup> AABOUD	<b>16</b> U	ATLS	pp; Z'_\eta $ ightarrow  { m e^+  e^-}$ , $\mu^+ \mu^-$	
• • • We do not us	se the foll	owing data for ave	erages,	, fits, lin	nits, etc. • • •	
>1870	95	<sup>2</sup> AAD	1200	ATLS	pp, Z'_\eta $ ightarrow$ $e^+e^-$ , $\mu^+\mu^-$	

> 938	95	<sup>3</sup> AALTONEN	11	CDF	$ ho \overline{ ho}; Z'_{\eta}  ightarrow \mu^+ \mu^-$
> 923	95	<sup>4</sup> ABAZOV	11A	D0	$p \overline{p}, Z'_{n} \rightarrow e^{+} e^{-}$
> 488	95	<sup>5</sup> DEL-AGUILA	10	RVUE	Electroweak
> 877	95	<sup>4</sup> AALTONEN	<b>09</b> T	CDF	p $\overline{p}$ , $Z'_\eta  ightarrow e^+e^-$
> 904	95	<sup>6</sup> AALTONEN	09v	CDF	Repl. by AALTONEN 11
> 427	95	<sup>7</sup> ERLER	09	RVUE	Electroweak
> 891	95	<sup>4</sup> AALTONEN	07н	CDF	Repl. by AALTONEN 09T
> 350	95	SCHAEL	07A	ALEP	e <sup>+</sup> e <sup>-</sup>
> 360	95	<sup>8</sup> ABDALLAH	<b>06</b> C	DLPH	e <sup>+</sup> e <sup>-</sup>
> 745		<sup>4</sup> ABULENCIA	06L	CDF	Repl. by AALTONEN 07H
> 720	95	<sup>9</sup> ABULENCIA	05A	CDF	Repl. by AALTONEN 111 and AALTONEN 09T
> 515	95	<sup>10</sup> ABBIENDI	<b>0</b> 4G	OPAL	$e^+e^-$
>1600		<sup>11</sup> BARGER	<b>03</b> B	COSM	Nucleosynthesis; light $ u_R$
> 310	95	<sup>12</sup> ABREU	<b>00</b> S	DLPH	e <sup>+</sup> e <sup>-</sup>
> 329	95	<sup>13</sup> BARATE	001	ALEP	Repl. by SCHAEL 07A
> 619	95	<sup>14</sup> CHO	00	RVUE	Electroweak
> 365	95	<sup>15</sup> ERLER	99	RVUE	Electroweak
> 87	95	<sup>16</sup> CONRAD	98	RVUE	$ u_{\mu} N$ scattering
> 620	95	<sup>17</sup> ABE	97s	CDF	$p \overline{p}; Z'_{\eta} \rightarrow e^+ e^-, \mu^+ \mu^-$
> 100	95	<sup>18</sup> VILAIN	<b>94</b> B	CHM2	$\nu_{\mu} e \rightarrow \nu_{\mu} e; \overline{\nu}_{\mu} e \rightarrow \overline{\nu}_{\mu} e$
> 125	90	<sup>19</sup> ABE	90F	VNS	e <sup>+</sup> e <sup>-</sup>
[> 820]		<sup>20</sup> GONZALEZ	<b>90</b> D	COSM	Nucleosynthesis; light $ u_R$
[> 3300]		<sup>21</sup> GRIFOLS	90	ASTR	SN 1987A; light $\nu_R$
[> 1040]		<sup>20</sup> LOPEZ	90		Nucleosynthesis; light $\nu_R$
[· =- · · ]					R

<sup>1</sup>AABOUD 16U search for resonances decaying to  $\ell^+ \ell^-$  in *pp* collisions at  $\sqrt{s} = 13$  TeV.

<sup>2</sup>AAD 12CC search for resonances decaying to  $e^+e^-$ ,  $\mu^+\mu^-$  in pp collisions at  $\sqrt{s}=7$ 

TeV. <sup>3</sup>AALTONEN 111 search for resonances decaying to  $\mu^+\mu^-$  in  $p\overline{p}$  collisions at  $\sqrt{s}=1.96$ 

<sup>4</sup>ABAZOV 11A, AALTONEN 09T, AALTONEN 07H, and ABULENCIA 06L search for resonances decaying to  $e^+e^-$  in  $p\overline{p}$  collisions at  $\sqrt{s} = 1.96$  TeV.

<sup>5</sup> DEL-AGUILA 10 give 95% CL limit on the Z-Z' mixing  $-0.0023 < \theta < 0.0027$ .

<sup>6</sup>AALTONEN 09V search for resonances decaying to  $\mu^+\mu^-$  in  $p\overline{p}$  collisions at  $\sqrt{s} =$ 1.96 TeV.

<sup>7</sup> ERLER 09 give 95% CL limit on the Z-Z' mixing  $-0.0047 < \theta < 0.0021$ .

<sup>8</sup>ABDALLAH 06C give 95% CL limit  $|\theta| <$  0.0092. See their Fig. 14 for limit contours in the mass-mixing plane.

<sup>9</sup>ABULENCIA 05A search for resonances decaying to electron or muon pairs in  $p \overline{p}$  collisions at  $\sqrt{s} = 1.96$  TeV.

 $^{10}\,{\sf ABBIENDI}$  04G give 95% CL limit on Z-Z' mixing  $-0.00447 < \theta < 0.00331.$  See their Fig. 20 for the limit contour in the mass-mixing plane.  $\sqrt{s} = 91$  to 207 GeV.

<sup>11</sup>BARGER 03B limit is from the nucleosynthesis bound on the effective number of light neutrino  $\delta N_{\nu}$  <1. The quark-hadron transition temperature  $T_c$ =150 MeV is assumed. The limit with  $T_c$ =400 MeV is >3300 GeV.

 $^{12}\text{ABREU}$  005 give 95% CL limit on Z-Z' mixing  $|\theta|$  < 0.0024. See their Fig.6 for the limit contour in the mass-mixing plane.  $\sqrt{s}$ =90 to 189 GeV.

 $^{13}$  BARATE 001 search for deviations in cross section and asymmetries in  $e^+\,e^ightarrow$  fermions at  $\sqrt{s}$ =90 to 183 GeV. Assume  $\theta$ =0. Bounds in the mass-mixing plane are shown in their Figure 18.

- <sup>14</sup> CHO 00 use various electroweak data to constrain Z' models assuming  $m_H$ =100 GeV. See Fig. 3 for limits in the mass-mixing plane.
- <sup>15</sup> ERLER 99 give 90% CL limit on the Z-Z' mixing  $-0.0062 < \theta < 0.0011$ .
- <sup>16</sup> CONRAD 98 limit is from measurements at CCFR, assuming no Z-Z' mixing.
- <sup>17</sup>ABE 97S find  $\sigma(Z') \times B(e^+e^-, \mu^+\mu^-) <$  40 fb for  $m_{Z'} > 600$  GeV at  $\sqrt{s} = 1.8$  TeV.
- $^{18}\,{\rm VILAIN}$  94B assume  $m_t$  = 150 GeV and  $\theta{=}0.~$  See Fig. 2 for limit contours in the mass-mixing plane.
- <sup>19</sup> ABE 90F use data for *R*,  $R_{\ell\ell}$ , and  $A_{\ell\ell}$ . ABE 90F fix  $m_W = 80.49 \pm 0.43 \pm 0.24$  GeV and  $m_Z = 91.13 \pm 0.03$  GeV.
- $^{20}$  These authors claim that the nucleosynthesis bound on the effective number of light neutrinos ( $\delta N_{\nu}$  < 1) constrains Z' masses if  $\nu_R$  is light ( $\lesssim$  1 MeV). <sup>21</sup> GRIFOLS 90 limit holds for  $m_{\nu_R} \lesssim$  1 MeV. See also GRIFOLS 90D, RIZZO 91.

Limits for other Z'	<u>CL%</u>	DOCUMENT ID TECN COMMENT
$\bullet \bullet \bullet$ We do not use the	ne follow	ing data for averages, fits, limits, etc. $ullet$ $ullet$ $ullet$
>1580	95	<sup>1</sup> AABOUD 17B ATLS $Z' \rightarrow hZ$
none 1100–1500	95	<sup>2</sup> AABOUD 16 ATLS $Z' \rightarrow b\overline{b}$
		<sup>3</sup> AAD 16L ATLS $Z' \rightarrow a\gamma$ , $a \rightarrow \gamma\gamma$
none 1500–2600	95	<sup>4</sup> AAD 16S ATLS $Z' \rightarrow q \overline{q}$
none 1000–1100, none 1300–1500	95	<sup>5</sup> KHACHATRY16AP CMS $Z' \rightarrow hZ$
>2400	95	<sup>6</sup> KHACHATRY16E CMS $Z' \rightarrow t \overline{t}$
		<sup>7</sup> AAD 15AO ATLS $Z' \rightarrow t \overline{t}$
		<sup>8</sup> AAD 15AT ATLS monotop
		<sup>9</sup> AAD 15CD ATLS $h \rightarrow ZZ', Z'Z'; Z' \rightarrow \ell^+ \ell^-$
		<sup>10</sup> KHACHATRY15F CMS monotop
		<sup>11</sup> KHACHATRY150 CMS $Z' \rightarrow hZ$
		$\begin{array}{ccc} 12 \\ 12 \\ 12 \end{array} \text{ AAD} \qquad 14 \text{AT LS}  Z' \to Z\gamma$
		<sup>13</sup> KHACHATRY14A CMS $Z' \rightarrow VV$
		14 MARTINEZ 14 RVUE Electroweak
none 500–1740	95	$\begin{array}{ccc} 15 \text{ AAD} & 13 \text{AQ ATLS} & Z' \rightarrow t \overline{t} \\ 16 & \cdots & & & & \\ \end{array}$
>1320 or 1000-1280	95	$\begin{array}{ccc} 16 \text{ AAD} & 13 \text{G} \text{ ATLS} & Z' \rightarrow t \overline{t} \\ 16 & z = z & z' \rightarrow z \end{array}$
> 915	95	$\begin{array}{cccc} 16 \text{ AALTONEN} & 13 \text{ A CDF} & Z' \rightarrow t \overline{t} \\ 17 \text{ ABLTONEN} & 13 \text{ A CDF} & Z' \rightarrow t \overline{t} \end{array}$
>1300	95	<sup>17</sup> CHATRCHYAN 13AP CMS $Z' \rightarrow t\overline{t}$
>2100	95	<sup>16</sup> CHATRCHYAN 13BMCMS $Z' \rightarrow t\overline{t}$
		$\begin{array}{ccc} 18 \text{ AAD} & 12 \text{ BV ATLS} & Z' \to t \overline{t} \\ 10 & & & \\ \end{array}$
		$\begin{array}{ccc} 19 \text{ AAD} & 12 \text{K} \text{ ATLS} & Z' \rightarrow t \overline{t} \\ 20 \text{ AAD} & 12 \text{K} \text{ ATLS} & Z' \rightarrow t \overline{t} \end{array}$
		<sup>20</sup> AALTONEN 12AR CDF Chromophilic
	~-	<sup>21</sup> AALTONEN 12N CDF $Z' \rightarrow \overline{t}u$
> 835	95	$\begin{array}{cccc} 22 \text{ ABAZOV} & 12 \text{ R} & \text{D0} & Z' \rightarrow t \overline{t} \\ 23 \text{ and } \overline{z} & \overline{z} & \overline{z} \end{array}$
		<sup>23</sup> CHATRCHYAN 12ai CMS $Z' \rightarrow t\overline{u}$
. 1400	<u>-</u>	<sup>24</sup> CHATRCHYAN 12AQ CMS $Z' \rightarrow t\overline{t}$
>1490	95	<sup>16</sup> CHATRCHYAN 12BL CMS $Z' \rightarrow t\overline{t}$
		<sup>25</sup> AALTONEN 11AD CDF $Z' \rightarrow t\overline{t}$
		<sup>26</sup> AALTONEN 11AE CDF $Z' \rightarrow t\overline{t}$
		<sup>27</sup> CHATRCHYAN 110 CMS $pp \rightarrow tt$

<sup>28</sup> AALTONEN	<b>08</b> D	CDF	$Z' \rightarrow t \overline{t}$
<sup>28</sup> AALTONEN	08Y	CDF	$Z' \rightarrow t \overline{t}$
<sup>28</sup> ABAZOV	08AA	D0	$Z' \rightarrow t \overline{t}$
<sup>29</sup> ABAZOV	04A	D0	Repl. by ABAZOV 08AA
<sup>30</sup> BARGER			Nucleosynthesis; light $\nu_R$
<sup>31</sup> CHO			E <sub>6</sub> -motivated
<sup>32</sup> CHO	98	RVUE	$E_6$ -motivated
<sup>33</sup> ABE	<b>97</b> G	CDF	$Z' \rightarrow \overline{q}q$

- <sup>1</sup>AABOUD 17B search for resonances decaying to hZ ( $h \rightarrow b\overline{b}$ ,  $c\overline{c}$ ;  $Z \rightarrow \ell^+\ell^-$ ,  $\nu\overline{\nu}$ ) in pp collisions at  $\sqrt{s} = 13$  TeV. The quoted limit is for heavy-vector-triplet Z' with  $g_V = 3$ . The limit becomes  $M_{Z'} > 1490$  GeV for  $g_V = 1$ . If we assume  $M_{Z'} = M_{W'}$ , the limit increases  $M_{Z'} > 2310$  GeV and  $M_{Z'} > 1730$  GeV for  $g_V = 3$  and  $g_V = 1$ , respectively. See their Fig.3 for limits on  $\sigma \cdot \overline{B}$ .
- <sup>2</sup> AABOUD 16 search for a narrow resonance decaying into  $b\overline{b}$  in pp collisions at  $\sqrt{s} = 13$ TeV. The limit quoted above is for a leptophobic Z' with SM-like couplings to quarks. See their Fig.6 for limits on  $\sigma \cdot B$ .
- <sup>3</sup>AAD 16L search for  $Z' \rightarrow a\gamma$ ,  $a \rightarrow \gamma\gamma$  in pp collisions at  $\sqrt{s} = 8$  TeV. See their Table 6 for limits on  $\sigma \cdot B$ .
- <sup>4</sup> AAD 16S search for a new resonance decaying to dijets in pp collisions at  $\sqrt{s} = 13$  TeV. The limit quoted above is for a leptophobic Z' having coupling strength with quark  $g_q = 0.3$  and is taken from their Figure 3.
- <sup>5</sup> KHACHATRYAN 16AP search for a resonance decaying to hZ in pp collisions at  $\sqrt{s} = 8$  TeV. Both h and Z are assumed to decay to fat jets. The quoted limit is for heavy-vector-triplet Z' with  $g_V = 3$ .
- <sup>6</sup> KHACHATRYAN 16E search for a leptophobic top-color Z' decaying to  $t\bar{t}$  using pp collisions at  $\sqrt{s} = 8$  TeV. The quoted limit assumes that  $\Gamma_{Z'}/m_{Z'} = 0.012$ . Also  $m_{Z'} < 2.9$  TeV is excluded for wider topcolor Z' with  $\Gamma_{Z'}/m_{Z'} = 0.1$ .
- <sup>7</sup>AAD 15AO search for narrow resonance decaying to  $t\bar{t}$  using pp collisions at  $\sqrt{s} = 8$  TeV. See Fig. 11 for limit on  $\sigma B$ .

<sup>8</sup> AAD 15AT search for monotop production plus large missing  $E_T$  events in pp collisions at  $\sqrt{s} = 8$  TeV and give constraints on a Z' model having Z'  $u\bar{t}$  coupling. Z' is assumed to decay invisibly. See their Fig. 6 for limits on  $\sigma \cdot B$ .

- <sup>9</sup> AAD 15CD search for decays of Higgs bosons to  $4 \ \ell$  states via Z' bosons,  $h \to ZZ' \to 4\ell$  or  $h \to Z'Z' \to 4\ell$ . See Fig. 5 for the limit on the signal strength of the  $h \to ZZ' \to 4\ell$  process and Fig. 16 for the limit on  $h \to Z'Z' \to 4\ell$ .
- <sup>10</sup> KHACHATRYAN 15F search for monotop production plus large missing  $E_T$  events in pp collisions at  $\sqrt{s} = 8$  TeV and give constraints on a Z' model having Z'  $u\bar{t}$  coupling. Z' is assumed to decay invisibly. See Fig. 3 for limits on  $\sigma B$ .
- <sup>11</sup> KHACHATRYAN 150 search for narrow Z' resonance decaying to Z h in pp collisions at  $\sqrt{s} = 8$  TeV. See their Fig. 6 for limit on  $\sigma B$ .
- <sup>12</sup> AAD 14AT search for a narrow neutral vector boson decaying to  $Z\gamma$ . See their Fig. 3b for the exclusion limit in  $m_{\gamma'} \sigma B$  plane.
- <sup>13</sup> KHACHATRYAN 14A search for new resonance in the  $WW(\ell \nu q \overline{q})$  and the  $ZZ(\ell \ell q \overline{q})$  channels using pp collisions at  $\sqrt{s}=8$  TeV. See their Fig.13 for the exclusion limit on the number of events in the mass-width plane.
- <sup>14</sup> MARTINEZ 14 use various electroweak data to constrain the Z' boson in the 3-3-1 \_\_\_\_ models.
- <sup>15</sup> AAD 13AQ search for a leptophobic top-color Z' decaying to  $t\overline{t}$ . The quoted limit assumes that  $\Gamma_{Z'}/m_{Z'} = 0.012$ .
- <sup>16</sup> CHATRCHYAN 13BM search for top-color Z' decaying to  $t\bar{t}$  using pp collisions at  $\sqrt{s}=8$  TeV. The quoted limit is for  $\Gamma_{Z'}/m_{Z'} = 0.012$ .

- <sup>17</sup> CHATRCHYAN 13AP search for top-color leptophobic Z' decaying to  $t\bar{t}$  using pp collisions at  $\sqrt{s}=7$  TeV. The quoted limit is for  $\Gamma_{Z'}/m_{Z'}=0.012$ .
- <sup>18</sup> AAD 12BV search for narrow resonance decaying to  $t\bar{t}$  using pp collisions at  $\sqrt{s}=7$  TeV. See their Fig. 7 for limit on  $\sigma \cdot B$ .
- <sup>19</sup>AAD 12K search for narrow resonance decaying to  $t\bar{t}$  using pp collisions at  $\sqrt{s}=7$  TeV. See their Fig. 5 for limit on  $\sigma \cdot B$ .
- <sup>20</sup> AALTONEN 12AR search for chromophilic Z' in  $p\overline{p}$  collisions at  $\sqrt{s} = 1.96$  TeV. See their Fig. 5 for limit on  $\sigma \cdot B$ .
- <sup>21</sup>AALTONEN 12N search for  $p\overline{p} \rightarrow tZ'$ ,  $Z' \rightarrow \overline{t}u$  events in  $p\overline{p}$  collisions. See their Fig. 3 for the limit on  $\sigma \cdot B$ .
- <sup>22</sup> ABAZOV 12R search for top-color Z' boson decaying exclusively to  $t\bar{t}$ . The quoted limit is for  $\Gamma_{Z'}/m_{Z'}=0.012$ .
- <sup>23</sup> CHATRCHYAN 12AI search for  $pp \rightarrow tt$  events and give constraints on a Z' model having  $Z'\overline{u}t$  coupling. See their Fig. 4 for the limit in mass-coupling plane.
- <sup>24</sup> Search for resonance decaying to  $t \bar{t}$ . See their Fig. 6 for limit on  $\sigma \cdot B$ .
- <sup>25</sup>Search for narrow resonance decaying to  $t\bar{t}$ . See their Fig. 4 for limit on  $\sigma \cdot B$ .
- <sup>26</sup>Search for narrow resonance decaying to  $t \bar{t}$ . See their Fig. 3 for limit on  $\sigma \cdot B$ .
- <sup>27</sup> CHATRCHYAN 110 search for same-sign top production in pp collisions induced by a hypothetical FCNC Z' at  $\sqrt{s} = 7$  TeV. See their Fig. 3 for limit in mass-coupling plane.
- <sup>28</sup>Search for narrow resonance decaying to  $t\bar{t}$ . See their Fig. 3 for limit on  $\sigma \cdot B$ .
- <sup>29</sup>Search for narrow resonance decaying to  $t \bar{t}$ . See their Fig. 2 for limit on  $\sigma \cdot B$ .
- <sup>30</sup> BARGER 03B use the nucleosynthesis bound on the effective number of light neutrino  $\delta N_{\nu}$ . See their Figs. 4–5 for limits in general  $E_6$  motivated models.
- <sup>31</sup> CHO 00 use various electroweak data to constrain Z' models assuming  $m_H$ =100 GeV. See Fig. 2 for limits in general  $E_6$ -motivated models.
- $^{32}$  CHO 98 study constraints on four-Fermi contact interactions obtained from low-energy electroweak experiments, assuming no Z-Z' mixing.
- <sup>33</sup>Search for Z' decaying to dijets at  $\sqrt{s}=1.8$  TeV. For Z' with electromagnetic strength coupling, no bound is obtained.

#### Searches for Z' with Lepton-Flavor-Violating decays

The following limits are obtained from  $p\overline{p}$  or  $pp \rightarrow Z'X$  with Z' decaying to the mode indicated in the comments. *DOCUMENT ID* TECN COMMENT

•	•	We d	o not	use	the	following	data	for	averages.	fits.	limits.	etc.	•	•	•

 $\begin{array}{cccccccc} 1 & \text{AABOUD} & 16 \text{P} & \text{ATLS} & Z' \rightarrow e\mu, e\tau, \mu\tau \\ 2 & \text{KHACHATRY...16BE CMS} & Z' \rightarrow e\mu \\ 3 & \text{AAD} & 150 & \text{ATLS} & Z' \rightarrow e\mu, e\tau, \mu\tau \\ 4 & \text{AAD} & 11 \text{H} & \text{ATLS} & Z' \rightarrow e\mu \\ 5 & \text{AAD} & 11 \text{Z} & \text{ATLS} & Z' \rightarrow e\mu \\ 6 & \text{ABULENCIA} & 06 \text{M} & \text{CDF} & Z' \rightarrow e\mu \\ \end{array}$ 

- <sup>1</sup>AABOUD 16P search for new particle with lepton flavor violating decay in pp collisions at  $\sqrt{s} = 13$  TeV. See their Figs.2, 3, and 4 for limits on  $\sigma \cdot B$ .
- <sup>2</sup> KHACHATRYAN 16BE search for new particle Z' with lepton flavor violating decay in pp collisions at  $\sqrt{s} = 8$  TeV in the range of 200 GeV  $< M_{Z'} < 2000$  GeV. See their Fig.4 for limits on  $\sigma \cdot B$  and their Table 5 for bounds on various masses.
- <sup>3</sup> AAD 150 search for new particle Z' with lepton flavor violating decay in pp collisions at  $\sqrt{s} = 8$  TeV in the range of 500 GeV  $< M_{Z'} < 3000$  GeV. See their Fig. 2 for limits on  $\sigma B$ .

<sup>4</sup> AAD 11H search for new particle Z' with lepton flavor violating decay in pp collisions at  $\sqrt{s} = 7$  TeV in the range of 700 GeV  $< M_{Z'} < 1000$  GeV. See their Fig. 3 for limits on  $\sigma \cdot B$ .

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VALUE

<sup>5</sup> AAD 11Z search for new particle Z' with lepton flavor violating decay in pp collisions at  $\sqrt{s} = 7$  TeV in the range 700 GeV  $< M_{Z'} < 2000$  GeV. See their Fig. 3 for limits on  $c \sigma \cdot B$ .

<sup>6</sup>ABULENCIA 06M search for new particle Z' with lepton flavor violating decay in  $p\overline{p}$  collisions at  $\sqrt{s} = 1.96$  TeV in the range of 100 GeV  $< M_{Z'} < 800$  GeV. See their Fig. 4 for limits in the mass-coupling plane.

#### Indirect Constraints on Kaluza-Klein Gauge Bosons

Bounds on a Kaluza-Klein excitation of the Z boson or photon in d=1 extra dimension. These bounds can also be interpreted as a lower bound on 1/R, the size of the extra dimension. Unless otherwise stated, bounds assume all fermions live on a single brane and all gauge fields occupy the 4+d-dimensional bulk. See also the section on "Extra Dimensions" in the "Searches" Listings in this *Review*.

VALUE (TeV)	CL%	DOCUMENT ID		TECN	COMMENT
• • • We do not use the	following	data for averages	, fits,	limits, e	etc. • • •
> 4.7		<sup>1</sup> MUECK	02	RVUE	Electroweak
> 3.3	95	<sup>2</sup> CORNET	00	RVUE	evqq′
>5000		<sup>3</sup> DELGADO	00	RVUE	€K
> 2.6	95	<sup>4</sup> DELGADO	00	RVUE	Electroweak
> 3.3	95	<sup>5</sup> RIZZO	00	RVUE	Electroweak
> 2.9	95	<sup>6</sup> MARCIANO	99	RVUE	Electroweak
> 2.5	95	<sup>7</sup> MASIP	99	RVUE	Electroweak
> 1.6	90	<sup>8</sup> NATH	99	RVUE	Electroweak
> 3.4	95	<sup>9</sup> STRUMIA	99	RVUE	Electroweak

<sup>1</sup> MUECK 02 limit is  $2\sigma$  and is from global electroweak fit ignoring correlations among observables. Higgs is assumed to be confined on the brane and its mass is fixed. For scenarios of bulk Higgs, of brane-SU(2)<sub>L</sub>, bulk-U(1)<sub>Y</sub>, and of bulk-SU(2)<sub>L</sub>, brane-U(1)<sub>Y</sub>, the corresponding limits are > 4.6 TeV, > 4.3 TeV and > 3.0 TeV, respectively.

<sup>2</sup>Bound is derived from limits on  $e\nu q q'$  contact interaction, using data from HERA and the Tevatron.

<sup>3</sup>Bound holds only if first two generations of quarks lives on separate branes. If quark mixing is not complex, then bound lowers to 400 TeV from  $\Delta m_K$ .

<sup>4</sup> See Figs. 1 and 2 of DELGADO 00 for several model variations. Special boundary conditions can be found which permit KK states down to 950 GeV and that agree with the measurement of  $Q_W$ (Cs). Quoted bound assumes all Higgs bosons confined to brane; placing one Higgs doublet in the bulk lowers bound to 2.3 TeV.

<sup>5</sup> Bound is derived from global electroweak analysis assuming the Higgs field is trapped on the matter brane. If the Higgs propagates in the bulk, the bound increases to 3.8 TeV.

- $^{6}$ Bound is derived from global electroweak analysis but considering only presence of the \_KK W bosons.
- <sup>7</sup>Global electroweak analysis used to obtain bound independent of position of Higgs on brane or in bulk.

<sup>8</sup> Bounds from effect of KK states on  $G_F$ ,  $\alpha$ ,  $M_W$ , and  $M_Z$ . Hard cutoff at string scale determined using gauge coupling unification. Limits for d=2,3,4 rise to 3.5, 5.7, and 7.8 TeV.

<sup>9</sup> Bound obtained for Higgs confined to the matter brane with  $m_H$ =500 GeV. For Higgs in the bulk, the bound increases to 3.5 TeV.

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#### MASS LIMITS for Leptoquarks from Pair Production

These limits rely only on the color or electroweak charge of the leptoquark.

VALUE (GeV)	CL%	DOCUMENT ID		TECN	COMMENT
>1050	95	<sup>1</sup> AAD	16G	ATLS	First generation
>1000	95	<sup>2</sup> AAD		ATLS	Second generation
> 625	95	<sup>3</sup> AAD	16G	ATLS	Third generation
none 200-640	95	<sup>4</sup> AAD		ATLS	Third generation
>1010	95	<sup>5</sup> KHACHATRY	.16AF	CMS	First generation
>1080	95	<sup>6</sup> KHACHATRY			Second generation
> 685	95	<sup>7</sup> KHACHATRY			Third generation
> 740	95	<sup>8</sup> KHACHATRY			Third generation
-		ollowing data for av			0
> 534	95	<sup>9</sup> AAD	13AE	ATLS	Third generation
> 525	95	<sup>10</sup> CHATRCHYAN	13M	CMS	Third generation
> 660	95	<sup>11</sup> AAD		ATLS	First generation
> 685	95	<sup>12</sup> AAD		ATLS	Second generation
> 830	95	<sup>13</sup> CHATRCHYAN			First generation
> 840	95	<sup>14</sup> CHATRCHYAN	12AG	CMS	Second generation
> 450	95	<sup>15</sup> CHATRCHYAN	12BO	CMS	Third generation
> 376	95	<sup>16</sup> AAD		ATLS	Superseded by AAD 12H
> 422	95	<sup>17</sup> AAD		ATLS	Superseded by AAD 120
> 326	95	<sup>18</sup> ABAZOV		D0	First generation
> 339	95	<sup>19</sup> CHATRCHYAN			Superseded by CHA-
> 384	95	<sup>20</sup> KHACHATRY	. <b>11</b> D	CMS	TRCHYAN 12AG Superseded by CHA-
> 394	95	<sup>21</sup> KHACHATRY	.11E	CMS	TRCHYAN 12AG Superseded by CHA-
> 247	95	<sup>22</sup> ABAZOV	10L	D0	TRCHYAN 12AG Third generation
> 316	95	<sup>23</sup> ABAZOV	09	D0	Second generation
> 299	95	<sup>24</sup> ABAZOV	09AF		Superseded by ABAZOV 11v
		<sup>25</sup> AALTONEN			Third generation
> 153	95	<sup>26</sup> AALTONEN	08Z		Third generation
> 205	95	<sup>27</sup> ABAZOV	08AD		All generations
> 210	95	<sup>26</sup> ABAZOV	08AN		Third generation
> 229	95	<sup>28</sup> ABAZOV		D0	Superseded by ABAZOV 10L
> 251	95	<sup>29</sup> ABAZOV		D0	Superseded by ABAZOV 09
> 136	95	<sup>30</sup> ABAZOV		D0	Superseded by ABAZOV 08AD
> 226	95	<sup>31</sup> ABULENCIA	06T		Second generation
> 256	95	<sup>32</sup> ABAZOV	05H		First generation
> 117	95	<sup>27</sup> ACOSTA	051	CDF	First generation
> 236	95 95	<sup>33</sup> ACOSTA			First generation
> 99	95 95	<sup>34</sup> ABBIENDI		OPAL	First generation
> 100	95 95	<sup>34</sup> ABBIENDI		OPAL	Second generation
> 98	95 95	<sup>34</sup> ABBIENDI		OPAL	Third generation
	95 95	<sup>35</sup> ABAZOV		D0	All generations
~~-	95 95	<sup>36</sup> ABAZOV		D0 D0	-
		<sup>37</sup> ABBIENDI		OPAL	First generation
> 85.8	95 05	<sup>37</sup> ABBIENDI			Superseded by ABBIENDI 03R
> 85.5	95 05	<sup>37</sup> ABBIENDI		OPAL	Superseded by ABBIENDI 03R
> 82.7	95 05	<sup>38</sup> ABBOTT		OPAL	Superseded by ABBIENDI 03R
> 200	95	ARROTT	00C	D0	Second generation

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> 123	95	<sup>39</sup> AFFOLDER	00K	CDF	Second generation
> 148	95	<sup>40</sup> AFFOLDER	00K	CDF	Third generation
> 160	95	<sup>41</sup> АВВОТТ	99J	D0	Second generation
> 225	95	<sup>42</sup> АВВОТТ	98E	D0	First generation
> 94	95	<sup>43</sup> АВВОТТ	98J	D0	Third generation
> 202	95	<sup>44</sup> ABE	98s	CDF	Second generation
> 242	95	<sup>45</sup> GROSS-PILCH	.98		First generation
> 99	95	<sup>46</sup> ABE	97F	CDF	Third generation
> 213	95	<sup>47</sup> ABE	97X	CDF	First generation
> 45.5	95	<sup>48,49</sup> ABREU	93J	DLPH	$First + second \ generation$
> 44.4	95	<sup>50</sup> ADRIANI	<b>9</b> 3M	L3	First generation
> 44.5	95	<sup>50</sup> ADRIANI	<b>9</b> 3M	L3	Second generation
> 45	95	<sup>50</sup> DECAMP	92	ALEP	Third generation
none 8.9–22.6	95	<sup>51</sup> KIM	90	AMY	First generation
none 10.2-23.2	95	<sup>51</sup> KIM	90	AMY	Second generation
none 5–20.8	95	<sup>52</sup> BARTEL	<b>87</b> B	JADE	
none 7–20.5	95	<sup>53</sup> BEHREND	<b>86</b> B	CELL	

<sup>1</sup>AAD 16G search for scalar leptoquarks using e e j j events in collisions at  $\sqrt{s} = 8$  TeV. The limit above assumes B(eq) = 1.

<sup>2</sup>AAD 16G search for scalar leptoquarks using  $\mu \mu j j$  events in collisions at  $\sqrt{s} = 8$  TeV. The limit above assumes  $B(\mu q) = 1$ .

<sup>3</sup>AAD 16G search for scalar leptoquarks decaying to  $b\nu$ . The limit above assumes  $B(b\nu) = 1$ .

<sup>4</sup> AAD 16G search for scalar leptoquarks decaying to  $t\nu$ . The limit above assumes  $B(t\nu) = 1$ .

<sup>5</sup> KHACHATRYAN 16AF search for scalar leptoquarks using eejj and  $e\nu jj$  events in pp collisions at  $\sqrt{s} = 8$  TeV. The limit above assumes B(eq) = 1. For B(eq) = 0.5, the limit becomes 850 GeV.

<sup>6</sup> KHACHATRYAN 16AF search for scalar leptoquarks using  $\mu\mu jj$  and  $\mu\nu jj$  events in pp collisions at  $\sqrt{s} = 8$  TeV. The limit above assumes  $B(\mu q) = 1$ . For  $B(\mu q) = 0.5$ , the \_limit becomes 760 GeV.

<sup>7</sup> KHACHATRYAN 15AJ search for scalar leptoquarks using  $\tau \tau t t$  events in pp collisions at  $\sqrt{s} = 8$  TeV. The limit above assumes  $B(\tau t) = 1$ .

<sup>8</sup> KHACHATRYAN 14T search for scalar leptoquarks decaying to  $\tau b$  using pp collisions at  $\sqrt{s}$ =8 TeV. The limit above assumes B( $\tau b$ ) = 1. See their Fig. 5 for exclusion limit as function of B( $\tau b$ ).

<sup>9</sup>AAD 13AE search for scalar leptoquarks using  $\tau \tau b b$  events in pp collisions at  $E_{cm} = 7$  TeV. The limit above assumes  $B(\tau b) = 1$ .

<sup>10</sup> CHATRCHYAN 13M search for scalar and vector leptoquarks decaying to  $\tau b$  in pp collisions at  $E_{cm} = 7$  TeV. The limit above is for scalar leptoquarks with  $B(\tau b) = 1$ .

<sup>11</sup> AAD 12H search for scalar leptoquarks using eejj and  $e\nu jj$  events in pp collisions at  $E_{cm} = 7$  TeV. The limit above assumes B(eq) = 1. For B(eq) = 0.5, the limit becomes 607 GeV.

<sup>12</sup> AAD 120 search for scalar leptoquarks using  $\mu \mu j j$  and  $\mu \nu j j$  events in p p collisions at  $E_{cm} = 7$  TeV. The limit above assumes  $B(\mu q) = 1$ . For  $B(\mu q) = 0.5$ , the limit becomes 594 GeV.

<sup>13</sup> CHATRCHYAN 12AG search for scalar leptoquarks using eejj and  $e\nu jj$  events in pp collisions at  $E_{cm} = 7$  TeV. The limit above assumes B(eq) = 1. For B(eq) = 0.5, the limit becomes 640 GeV.

<sup>14</sup> CHATRCHYAN 12AG search for scalar leptoquarks using  $\mu \mu j j$  and  $\mu \nu j j$  events in p p collisions at  $E_{\rm cm} = 7$  TeV. The limit above assumes  $B(\mu q) = 1$ . For  $B(\mu q) = 0.5$ , the limit becomes 650 GeV.

<sup>15</sup> CHATRCHYAN 12BO search for scalar leptoquarks decaying to  $\nu b$  in pp collisions at  $\sqrt{s} = 7$  TeV. The limit above assumes  $B(\nu b) = 1$ .

- <sup>16</sup> AAD 11D search for scalar leptoquarks using eejj and  $e\nu jj$  events in pp collisions at  $E_{cm} = 7$  TeV.The limit above assumes B(eq) = 1. For B(eq) = 0.5, the limit becomes 319 GeV.
- <sup>17</sup> AAD 11D search for scalar leptoquarks using  $\mu\mu jj$  and  $\mu\nu jj$  events in pp collisions at  $E_{\rm cm} = 7$  TeV. The limit above assumes  $B(\mu q) = 1$ . For  $B(\mu q) = 0.5$ , the limit becomes 362 GeV.
- <sup>18</sup>ABAZOV 11V search for scalar leptoquarks using  $e\nu jj$  events in  $p\overline{p}$  collisions at  $E_{\rm cm}$  = 1.96 TeV. The limit above assumes B(eq) = 0.5.
- <sup>19</sup> CHATRCHYAN 11N search for scalar leptoquarks using  $e\nu jj$  events in pp collisions at  $E_{\rm cm} = 7$  TeV. The limit above assumes B(eq) = 0.5.
- <sup>20</sup> KHACHATRYAN 11D search for scalar leptoquarks using e e j j events in p p collisions at  $E_{cm} = 7$  TeV. The limit above assumes B(e q) = 1.
- <sup>21</sup> KHACHATRYAN 11E search for scalar leptoquarks using  $\mu \mu j j$  events in pp collisions at  $E_{cm} = 7$  TeV. The limit above assumes  $B(\mu q) = 1$ .
- <sup>22</sup> ABAZOV 10L search for pair productions of scalar leptoquark state decaying to  $\nu b$  in  $p\overline{p}$  collisions at  $E_{\rm cm} = 1.96$  TeV. The limit above assumes B( $\nu b$ ) = 1.
- <sup>23</sup>ABAZOV 09 search for scalar leptoquarks using  $\mu\mu jj$  and  $\mu\nu jj$  events in  $p\overline{p}$  collisions at  $E_{\rm cm} = 1.96$  TeV. The limit above assumes  $B(\mu q) = 1$ . For  $B(\mu q) = 0.5$ , the limit becomes 270 GeV.
- <sup>24</sup> ABAZOV 09AF search for scalar leptoquarks using e e j j and  $e \nu j j$  events in  $p \overline{p}$  collisions at  $E_{\rm cm} = 1.96$  TeV. The limit above assumes B(eq) = 1. For B(eq) = 0.5 the bound becomes 284 GeV.
- <sup>25</sup> AALTONEN 08P search for vector leptoquarks using  $\tau^+ \tau^- b\overline{b}$  events in  $p\overline{p}$  collisions at  $E_{\rm cm} = 1.96$  TeV. Assuming Yang-Mills (minimal) couplings, the mass limit is >317 GeV (251 GeV) at 95% CL for B( $\tau b$ ) = 1.
- <sup>26</sup>Search for pair production of scalar leptoquark state decaying to  $\tau b$  in  $p\overline{p}$  collisions at  $E_{cm} = 1.96$  TeV. The limit above assumes  $B(\tau b) = 1$ .
- <sup>27</sup> Search for scalar leptoquarks using  $\nu \nu j j$  events in  $\overline{p}p$  collisions at  $E_{\rm cm} = 1.96$  TeV. The limit above assumes  $B(\nu q) = 1$ .
- <sup>28</sup>ABAZOV 07J search for pair productions of scalar leptoquark state decaying to  $\nu b$  in  $p\overline{p}$  collisions at  $E_{\rm cm} = 1.96$  TeV. The limit above assumes  $B(\nu b) = 1$ .
- <sup>29</sup> ABAZOV 06A search for scalar leptoquarks using  $\mu \mu j j$  events in  $p \overline{p}$  collisions at  $E_{cm} = 1.8$  TeV and 1.96 TeV. The limit above assumes  $B(\mu q) = 1$ . For  $B(\mu q) = 0.5$ , the limit becomes 204 GeV.
- <sup>30</sup>ABAZOV 06L search for scalar leptoquarks using  $\nu \nu j j$  events in  $p \overline{p}$  collisions at  $E_{cm} = 1.8$  TeV and at 1.96 TeV. The limit above assumes B( $\nu q$ ) = 1.
- <sup>31</sup> ABULENCIA 06T search for scalar leptoquarks using  $\mu\mu jj$ ,  $\mu\nu jj$ , and  $\nu\nu jj$  events in  $p\overline{p}$  collisions at  $E_{\rm cm} = 1.96$  TeV. The quoted limit assumes  $B(\mu q) = 1$ . For  $B(\mu q) = 0.5$  or 0.1, the bound becomes 208 GeV or 143 GeV, respectively. See their Fig. 4 for the exclusion limit as a function of  $B(\mu q)$ .
- <sup>32</sup>ABAZOV 05H search for scalar leptoquarks using eejj and  $e\nu jj$  events in  $\overline{p}p$  collisions at  $E_{cm} = 1.8$  TeV and 1.96 TeV. The limit above assumes B(eq) = 1. For B(eq) = 0.5 the bound becomes 234 GeV.
- <sup>33</sup> ACOSTA 05P search for scalar leptoquarks using eejj,  $e\nu jj$  events in  $\overline{p}p$  collisions at  $E_{\rm cm} = 1.96$ TeV. The limit above assumes B(eq) = 1. For B(eq) = 0.5 and 0.1, the bound becomes 205 GeV and 145 GeV, respectively.
- <sup>34</sup> ABBIENDI 03R search for scalar/vector leptoquarks in  $e^+e^-$  collisions at  $\sqrt{s} = 189-209$  GeV. The quoted limits are for charge -4/3 isospin 0 scalar-leptoquark with  $B(\ell q) = 1$ . See their table 12 for other cases.
- <sup>35</sup> ABAZOV 02 search for scalar leptoquarks using  $\nu \nu j j$  events in  $\overline{p} p$  collisions at  $E_{cm} = 1.8$  TeV. The bound holds for all leptoquark generations. Vector leptoquarks are likewise constrained to lie above 200 GeV.
- <sup>36</sup> ABAZOV 01D search for scalar leptoquarks using  $e\nu jj$ , eejj, and  $\nu\nu jj$  events in  $p\overline{p}$  collisions at  $E_{\rm cm}$ =1.8 TeV. The limit above assumes B(eq)=1. For B(eq)=0.5 and 0, the bound becomes 204 and 79 GeV, respectively. Bounds for vector leptoquarks are also given. Supersedes ABBOTT 98E.

- <sup>37</sup> ABBIENDI 00M search for scalar/vector leptoquarks in  $e^+e^-$  collisions at  $\sqrt{s}=183$  GeV. The quoted limits are for charge -4/3 isospin 0 scalar-leptoquarks with B( $\ell q$ )=1. See their Table 8 and Figs. 6–9 for other cases.
- <sup>38</sup>ABBOTT 00C search for scalar leptoquarks using  $\mu\mu jj$ ,  $\mu\nu jj$ , and  $\nu\nu jj$  events in  $p\overline{p}$  collisions at  $E_{\rm cm}$ =1.8 TeV. The limit above assumes B( $\mu q$ )=1. For B( $\mu q$ )=0.5 and 0, the bound becomes 180 and 79 GeV respectively. Bounds for vector leptoquarks are also given.
- <sup>39</sup> AFFOLDER 00K search for scalar leptoquark using  $\nu\nu cc$  events in  $p\overline{p}$  collisions at  $E_{\rm cm}$ =1.8 TeV. The quoted limit assumes B( $\nu c$ )=1. Bounds for vector leptoquarks are also given.
- <sup>40</sup> AFFOLDER 00K search for scalar leptoquark using  $\nu\nu bb$  events in  $p\overline{p}$  collisions at  $E_{\rm cm}$ =1.8 TeV. The quoted limit assumes B( $\nu b$ )=1. Bounds for vector leptoquarks are also given.
- <sup>41</sup> ABBOTT 99J search for leptoquarks using  $\mu \nu j j$  events in  $p \overline{p}$  collisions at  $E_{cm} = 1.8$ TeV. The quoted limit is for a scalar leptoquark with  $B(\mu q) = B(\nu q) = 0.5$ . Limits on vector leptoquarks range from 240 to 290 GeV.
- <sup>42</sup> ABBOTT 98E search for scalar leptoquarks using  $e\nu jj$ , eejj, and  $\nu\nu jj$  events in  $p\overline{p}$  collisions at  $E_{\rm cm}$ =1.8 TeV. The limit above assumes B(eq)=1. For B(eq)=0.5 and 0, the bound becomes 204 and 79 GeV, respectively.
- <sup>43</sup>ABBOTT 98J search for charge -1/3 third generation scalar and vector leptoquarks in  $p\overline{p}$  collisions at  $E_{cm} = 1.8$  TeV. The quoted limit is for scalar leptoquark with B( $\nu b$ )=1.
- <sup>44</sup>ABE 98S search for scalar leptoquarks using  $\mu\mu jj$  events in  $p\overline{p}$  collisions at  $E_{cm}$ = 1.8 TeV. The limit is for B( $\mu q$ )= 1. For B( $\mu q$ )=B( $\nu q$ )=0.5, the limit is > 160 GeV.
- <sup>45</sup> GROSS-PILCHER 98 is the combined limit of the CDF and DØ Collaborations as determined by a joint CDF/DØ working group and reported in this FNAL Technical Memo. Original data published in ABE 97X and ABBOTT 98E.
- <sup>46</sup> ABE 97F search for third generation scalar and vector leptoquarks in  $p\overline{p}$  collisions at  $E_{\rm cm} = 1.8$  TeV. The quoted limit is for scalar leptoquark with B( $\tau b$ ) = 1.
- <sup>47</sup> ABE 97X search for scalar leptoquarks using eejj events in  $p\overline{p}$  collisions at  $E_{cm}=1.8$  TeV. The limit is for B(eq)=1.
- <sup>48</sup> Limit is for charge -1/3 isospin-0 leptoquark with B( $\ell q$ ) = 2/3.
- <sup>49</sup> First and second generation leptoquarks are assumed to be degenerate. The limit is slightly lower for each generation.
- <sup>50</sup> Limits are for charge -1/3, isospin-0 scalar leptoquarks decaying to  $\ell^- q$  or  $\nu q$  with any branching ratio. See paper for limits for other charge-isospin assignments of leptoquarks.
- <sup>51</sup> KIM 90 assume pair production of charge 2/3 scalar-leptoquark via photon exchange. The decay of the first (second) generation leptoquark is assumed to be any mixture of  $de^+$  and  $u\overline{\nu}$  ( $s\mu^+$  and  $c\overline{\nu}$ ). See paper for limits for specific branching ratios.
- <sup>52</sup> BARTEL 87B limit is valid when a pair of charge 2/3 spinless leptoquarks X is produced with point coupling, and when they decay under the constraint  $B(X \rightarrow c\overline{\nu}_{\mu}) + B(X \rightarrow s\mu^+) = 1$ .
- <sup>53</sup>BEHREND 86B assumed that a charge 2/3 spinless leptoquark,  $\chi$ , decays either into  $s\mu^+$  or  $c\overline{\nu}$ : B( $\chi \rightarrow s\mu^+$ ) + B( $\chi \rightarrow c\overline{\nu}$ ) = 1.

#### MASS LIMITS for Leptoquarks from Single Production

These limits depend on the *q*- $\ell$ -leptoquark coupling  $g_{LQ}$ . It is often assumed that  $g_{LQ}^2/4\pi = 1/137$ . Limits shown are for a scalar, weak isoscalar, charge -1/3 leptoquark.

VALUE (GeV)	<u>CL%</u>	DOCUMENT ID	TECN	COMMENT	
>1755	95 05	<sup>1</sup> KHACHATRY1	6AG CMS	First generation Second generation	
> <b>660</b> > 304	95 95	<sup>3</sup> ABRAMOWICZ1			
> 73	95	<sup>4</sup> ABREU 9	3J DLPH	Second generation	
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• • We do not use the following data for averages, fits, limits, etc. • • •

		<sup>5</sup> DEY	16	ICCB	$\nu q \rightarrow LQ \rightarrow \nu q$
		<sup>6</sup> AARON	11A	H1	Lepton-flavor violation
> 300	95	<sup>7</sup> AARON	11B	H1	First generation
		<sup>8</sup> ABAZOV	07E	D0	Second generation
> 295	95	<sup>9</sup> AKTAS	<b>05</b> B	H1	First generation
		<sup>10</sup> CHEKANOV	05A	ZEUS	Lepton-flavor violation
> 298	95	<sup>11</sup> CHEKANOV	<b>03</b> B	ZEUS	First generation
> 197	95	<sup>12</sup> ABBIENDI	<b>0</b> 2B	OPAL	First generation
		<sup>13</sup> CHEKANOV	02	ZEUS	Repl. by CHEKANOV 05A
> 290	95	<sup>14</sup> ADLOFF	<b>01</b> C	H1	First generation
> 204	95	<sup>15</sup> BREITWEG	01	ZEUS	First generation
		<sup>16</sup> BREITWEG	00e	ZEUS	First generation
> 161	95	<sup>17</sup> ABREU	<b>99</b> G	DLPH	First generation
> 200	95	<sup>18</sup> ADLOFF	99	H1	First generation
		<sup>19</sup> DERRICK	97	ZEUS	Lepton-flavor violation
> 168	95	<sup>20</sup> DERRICK	93	ZEUS	First generation

<sup>1</sup> KHACHATRYAN 16AG search for single production of charge  $\pm 1/3$  scalar leptoquarks using e e j events in p p collisions at  $\sqrt{s} = 8$  TeV. The limit above assumes B(eq) = 1 and the leptoquark coupling strength  $\lambda = 1$ .

<sup>2</sup>KHACHATRYAN 16AG search for single production of charge  $\pm 1/3$  scalar leptoquarks using  $\mu \mu j$  events in pp collisions at  $\sqrt{s} = 8$  TeV. The limit above assumes  $B(\mu q) = 1$  and the leptoquark coupling strength  $\lambda = 1$ .

- <sup>3</sup>ABRAMOWICZ 12A limit is for a scalar, weak isoscalar, charge -1/3 leptoquark coupled with  $e_R$ . See their Figs. 12–17 and Table 4 for states with different quantum numbers.
- <sup>4</sup>Limit from single production in Z decay. The limit is for a leptoquark coupling of electromagnetic strength and assumes  $B(\ell q) = 2/3$ . The limit is 77 GeV if first and second leptoquarks are degenerate.
- <sup>5</sup>DEY 16 use the 2010-2012 IceCube PeV energy data set to constrain the leptoquark production cross section through the  $\nu q \rightarrow LQ \rightarrow \nu q$  process. See their Figure 4 for the exclusion limit in the mass-coupling plane.
- <sup>6</sup>AARON 11A search for various leptoquarks with lepton-flavor violating couplings. See their Figs. 2–3 and Tables 1–4 for detailed limits.
- <sup>7</sup> The quoted limit is for a scalar, weak isoscalar, charge -1/3 leptoquark coupled with  $e_R$ . See their Figs. 3–5 for limits on states with different quantum numbers.
- <sup>8</sup>ABAZOV 07E search for leptoquark single production through qg fusion process in  $p\overline{p}$  collisions. See their Fig. 4 for exclusion plot in mass-coupling plane.
- <sup>9</sup>AKTAS 05B limit is for a scalar, weak isoscalar, charge -1/3 leptoquark coupled with  $e_R$ . See their Fig. 3 for limits on states with different quantum numbers.
- <sup>10</sup> CHEKANOV 05 search for various leptoquarks with lepton-flavor violating couplings. See their Figs.6–10 and Tables 1–8 for detailed limits.
- <sup>11</sup> CHEKANOV 03B limit is for a scalar, weak isoscalar, charge -1/3 leptoquark coupled with  $e_R$ . See their Figs. 11–12 and Table 5 for limits on states with different quantum numbers.
- $^{12}$  For limits on states with different quantum numbers and the limits in the mass-coupling plane, see their Fig. 4 and Fig. 5.
- <sup>13</sup> CHEKANOV 02 search for various leptoquarks with lepton-flavor violating couplings. See their Figs. 6–7 and Tables 5–6 for detailed limits.
- <sup>14</sup> For limits on states with different quantum numbers and the limits in the mass-coupling plane, see their Fig. 3.
- <sup>15</sup>See their Fig. 14 for limits in the mass-coupling plane.
- <sup>16</sup> BREITWEG 00E search for F=0 leptoquarks in  $e^+ p$  collisions. For limits in masscoupling plane, see their Fig. 11.

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- <sup>17</sup>ABREU 99G limit obtained from process  $e\gamma \rightarrow LQ+q$ . For limits on vector and scalar states with different quantum numbers and the limits in the coupling-mass plane, see their Fig. 4 and Table 2.
- <sup>18</sup> For limits on states with different quantum numbers and the limits in the mass-coupling plane, see their Fig. 13 and Fig. 14. ADLOFF 99 also search for leptoquarks with lepton-flavor violating couplings. ADLOFF 99 supersedes AID 96B.
- $^{19}\,\text{DERRICK}$  97 search for various leptoquarks with lepton-flavor violating couplings. See their Figs. 5–8 and Table 1 for detailed limits.
- <sup>20</sup> DERRICK 93 search for single leptoquark production in ep collisions with the decay eq and  $\nu q$ . The limit is for leptoquark coupling of electromagnetic strength and assumes  $B(eq) = B(\nu q) = 1/2$ . The limit for B(eq) = 1 is 176 GeV. For limits on states with different quantum numbers, see their Table 3.

#### Indirect Limits for Leptoquarks

VAL	UE (TeV)	<u>CL%</u>	DOCUMENT ID		TECN	COMMENT
• •	• We do	not use	the following data f	for av	erages, f	its, limits, etc. • • •
			<sup>1</sup> BARRANCO	16	RVUE	D decays
			<sup>2</sup> KUMAR	16	RVUE	neutral $K$ mixing, rare $K$ decays
			<sup>3</sup> BESSAA	15	RVUE	$q \overline{q} \rightarrow e^+ e^-$
>	• 14	95	<sup>4</sup> SAHOO	15A	RVUE	$B_{s,d} \rightarrow \mu^+ \mu^-$
			<sup>5</sup> SAKAKI	13	RVUE	$B \rightarrow D^{(*)} \tau \overline{\nu}, B \rightarrow X_{s} \nu \overline{\nu}$
			<sup>6</sup> KOSNIK	12	RVUE	$b \rightarrow s \ell^+ \ell^-$
>	2.5	95	<sup>7</sup> AARON	11C	H1	First generation
			<sup>8</sup> DORSNER	11	RVUE	scalar, weak singlet, charge 4/3
			<sup>9</sup> AKTAS	07A	H1	Lepton-flavor violation
>	0.49	95	<sup>10</sup> SCHAEL	07A	ALEP	$e^+e^- \rightarrow q \overline{q}$
			<sup>11</sup> SMIRNOV	07	RVUE	$K  ightarrow e \mu, \ B  ightarrow e  au$
			<sup>12</sup> CHEKANOV	05A	ZEUS	Lepton-flavor violation
>	1.7	96	<sup>13</sup> ADLOFF	03	H1	First generation
>	46	90	<sup>14</sup> CHANG	03	BELL	Pati-Salam type
			<sup>15</sup> CHEKANOV	02	ZEUS	Repl. by CHEKANOV 05A
>	1.7	95	<sup>16</sup> CHEUNG	<b>01</b> B	RVUE	First generation
>	0.39	95	<sup>17</sup> ACCIARRI	<b>00</b> P	L3	$e^+e^- \rightarrow q q$
>	1.5	95	<sup>18</sup> ADLOFF	00	H1	First generation
>	0.2	95	<sup>19</sup> BARATE	001	ALEP	Repl. by SCHAEL 07A
			<sup>20</sup> BARGER	00	RVUE	Cs
			<sup>21</sup> GABRIELLI	00	RVUE	Lepton flavor violation
>	0.74	95	<sup>22</sup> ZARNECKI	00	RVUE	$S_1$ leptoquark
			<sup>23</sup> ABBIENDI	99	OPAL	
>	19.3	95	<sup>24</sup> ABE	98v	CDF	$B_{m{s}}  o ~e^{\pm} \mu^{\mp}$ , Pati-Salam type
			<sup>25</sup> ACCIARRI	98J	L3	$e^+e^- \rightarrow q \overline{q}$
			<sup>26</sup> ACKERSTAFF	98v	OPAL	~
>	0.76	95	<sup>27</sup> DEANDREA	97	RVUE	$\widetilde{R}_2$ leptoquark
			<sup>28</sup> DERRICK	97	ZEUS	Lepton-flavor violation
			<sup>29</sup> GROSSMAN	97	RVUE	$B \rightarrow \tau^+ \tau^-(X)$
			<sup>30</sup> JADACH	97	RVUE	$e^+e^-  ightarrow q \overline{q}$
>1	200		<sup>31</sup> KUZNETSOV	<b>95</b> B	RVUE	Pati-Salam type
			<sup>32</sup> MIZUKOSHI	95	RVUE	Third generation scalar leptoquark
>	0.3	95	<sup>33</sup> BHATTACH	94	RVUE	Spin-0 leptoquark coupled to $\overline{e}_R t_L$

			<sup>34</sup> DAVIDSON	94	RVUE	
>	18		<sup>35</sup> KUZNETSOV	94	RVUE	Pati-Salam type
>	0.43	95	<sup>36</sup> LEURER	94	RVUE	First generation spin-1 leptoquark
>	0.44	95	<sup>36</sup> LEURER	<b>94</b> B	RVUE	First generation spin-0 leptoquark
			<sup>37</sup> MAHANTA	94	RVUE	P and T violation
>	1		<sup>38</sup> SHANKER	82	RVUE	Nonchiral spin-0 leptoquark
> 3	125		<sup>38</sup> SHANKER	82	RVUE	Nonchiral spin-1 leptoquark

<sup>1</sup>BARRANCO 16 give bounds on leptoquark induced four-fermion interactions from  $D \rightarrow$  $K\ell\nu$  and  $D_{s} \rightarrow \ell\nu$ .

<sup>2</sup> KUMAR 16 gives bound on SU(2) singlet scalar leptoquark with chrge -1/3 from  $K^0 - \overline{K}^0$  mixing,  $K \to \pi \nu \overline{\nu}$ ,  $K^0_L \to \mu^+ \mu^-$ , and  $K^0_L \to \mu^\pm e^\mp$  decays.

 $^3$  BESSAA 15 obtain limit on leptoquark induced four-fermion interactions from the ATLAS and CMS limit on the  $\overline{q} q \overline{e} e$  contact interactions.

<sup>4</sup>SAHOO 15A obtain limit on leptoquark induced four-fermion interactions from  $B_{s,d} \rightarrow$  $\mu^+\mu^-$  for  $\lambda \simeq O(1)$ .

<sup>5</sup>SAKAKI 13 explain the  $B \rightarrow D^{(*)} \tau \overline{\nu}$  anomaly using Wilson coefficients of leptoquarkinduced four-fermion operators.

 $^{6}$ KOSNIK 12 obtains limits on leptoquark induced four-fermion interactions from  $b \rightarrow$  $s\ell^+\ell^-$  decays.

<sup>7</sup>AARON 11C limit is for weak isotriplet spin-0 leptoquark at strong coupling  $\lambda = \sqrt{4\pi}$ . For the limits of leptoquarks with different quantum numbers, see their Table 3. Limits are derived from bounds of *eq* contact intereractions.

- $^8$  DORSNER 11 give bounds on scalar, weak singlet, charge 4/3 leptoquark from K, B, audecays, meson mixings, LFV, g-2 and  $Z \rightarrow b\overline{b}$ .
- $^9$ AKTAS 07A search for lepton-flavor violation in ep collision. See their Tables 4–7 for limits on lepton-flavor violating four-fermion interactions induced by various leptoquarks.
- $^{10}$  SCHAEL 07A limit is for the weak-isoscalar spin-0 left-handed leptoquark with the coupling of electromagnetic strength. For the limits of leptoquarks with different quantum numbers, see their Table 35.
- $^{11}$ SMIRNOV 07 obtains mass limits for the vector and scalar chiral leptoquark states from  $K \rightarrow e \mu, B \rightarrow e \tau$  decays.

 $^{12}$  CHEKANOV 05 search for various leptoquarks with lepton-flavor violating couplings. See their Figs.6-10 and Tables 1-8 for detailed limits.

- <sup>13</sup> ADLOFF 03 limit is for the weak isotriplet spin-0 leptoquark at strong coupling  $\lambda = \sqrt{4\pi}$ . For the limits of leptoquarks with different quantum numbers, see their Table 3. Limits are derived from bounds on  $e^{\pm}q$  contact interactions.
- <sup>14</sup> The bound is derived from B( $B^0 \rightarrow e^{\pm} \mu^{\mp}$ ) <  $1.7 \times 10^{-7}$ .
- $^{15}$  CHEKANOV 02 search for lepton-flavor violation in *ep* collisions. See their Tables 1–4 for limits on lepton-flavor violating and four-fermion interactions induced by various leptoquarks.
- $^{16}$  CHEUNG 01B quoted limit is for a scalar, weak isoscalar, charge -1/3 leptoquark with a coupling of electromagnetic strength. The limit is derived from bounds on contact interactions in a global electroweak analysis. For the limits of leptoquarks with different quantum numbers, see Table 5.
- <sup>17</sup> ACCIARRI 00P limit is for the weak isoscalar spin-0 leptoquark with the coupling of electromagnetic strength. For the limits of leptoquarks with different quantum numbers, see their Table 4.
- <sup>18</sup> ADLOFF 00 limit is for the weak isotriplet spin-0 leptoquark at strong coupling,  $\lambda = \sqrt{4\pi}$ . For the limits of leptoquarks with different quantum numbers, see their Table 2. ADLOFF 00 limits are from the  $Q^2$  spectrum measurement of  $e^+ p \rightarrow e^+ X$ .
- $^{19}$  BARATE 001 search for deviations in cross section and jet-charge asymmetry in  $e^+\,e^ightarrow$  $\overline{q}q$  due to *t*-channel exchange of a leptoquark at  $\sqrt{s}=130$  to 183 GeV. Limits for other scalar and vector leptoquarks are also given in their Table 22.

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- <sup>20</sup> BARGER 00 explain the deviation of atomic parity violation in cesium atoms from prediction is explained by scalar leptoquark exchange.
- $^{21}$  GABRIELLI 00 calculate various process with lepton flavor violation in leptoquark models.
- <sup>22</sup> ZARNECKI 00 limit is derived from data of HERA, LEP, and Tevatron and from various low-energy data including atomic parity violation. Leptoquark coupling with electromagnetic strength is assumed.
- <sup>23</sup> ABBIENDI 99 limits are from  $e^+e^- \rightarrow q\overline{q}$  cross section at 130–136, 161–172, 183 GeV. See their Fig. 8 and Fig. 9 for limits in mass-coupling plane. <sup>24</sup> ABE 98V quoted limit is from B( $B_s \rightarrow e^{\pm}\mu^{\mp}$ )< 8.2 × 10<sup>-6</sup>. ABE 98V also obtain
- <sup>24</sup> ABE 98V quoted limit is from  $B(B_s \rightarrow e^{\pm}\mu^{+}) < 8.2 \times 10^{-0}$ . ABE 98V also obtain a similar limit on  $M_{LQ} > 20.4$  TeV from  $B(B_d \rightarrow e^{\pm}\mu^{\mp}) < 4.5 \times 10^{-6}$ . Both bounds assume the non-canonical association of the *b* quark with electrons or muons under SU(4).
- <sup>25</sup> ACCIARRI 98J limit is from  $e^+e^- \rightarrow q\bar{q}$  cross section at  $\sqrt{s} = 130-172$  GeV which can be affected by the *t* and *u*-channel exchanges of leptoquarks. See their Fig. 4 and Fig. 5 for limits in the mass-coupling plane.
- <sup>26</sup> ACKERSTAFF 98V limits are from  $e^+e^- \rightarrow q \overline{q}$  and  $e^+e^- \rightarrow b \overline{b}$  cross sections at  $\sqrt{s} = 130-172$  GeV, which can be affected by the *t* and *u*-channel exchanges of leptoquarks. See their Fig. 21 and Fig. 22 for limits of leptoquarks in mass-coupling plane.
- <sup>27</sup> DEANDREA 97 limit is for  $\tilde{R}_2$  leptoquark obtained from atomic parity violation (APV). The coupling of leptoquark is assumed to be electromagnetic strength. See Table 2 for limits of the four-fermion interactions induced by various scalar leptoquark exchange. DEANDREA 97 combines APV limit and limits from Tevatron and HERA. See Fig. 1–4 for combined limits of leptoquark in mass-coupling plane.
- <sup>28</sup> DERRICK 97 search for lepton-flavor violation in *ep* collision. See their Tables 2–5 for limits on lepton-flavor violating four-fermion interactions induced by various leptoquarks.
- <sup>29</sup> GROSSMAN 97 estimate the upper bounds on the branching fraction  $B \rightarrow \tau^+ \tau^-(X)$  from the absence of the *B* decay with large missing energy. These bounds can be used to constrain leptoquark induced four-fermion interactions.
- <sup>30</sup> JADACH 97 limit is from  $e^+e^- \rightarrow q \overline{q}$  cross section at  $\sqrt{s}$ =172.3 GeV which can be affected by the *t* and *u*-channel exchanges of leptoquarks. See their Fig. 1 for limits on vector leptoquarks in mass-coupling plane.
- <sup>31</sup> KUZNETSOV 95B use  $\pi$ , K, B,  $\tau$  decays and  $\mu e$  conversion and give a list of bounds on the leptoquark mass and the fermion mixing matrix in the Pati-Salam model. The quoted limit is from  $K_I \rightarrow \mu e$  decay assuming zero mixing.
- <sup>32</sup> MIZUKOSHI 95 calculate the one-loop radiative correction to the *Z*-physics parameters in various scalar leptoquark models. See their Fig. 4 for the exclusion plot of third generation leptoquark models in mass-coupling plane.
- <sup>33</sup> BHATTACHARYYA 94 limit is from one-loop radiative correction to the leptonic decay width of the Z.  $m_H$ =250 GeV,  $\alpha_s(m_Z)$ =0.12,  $m_t$ =180 GeV, and the electroweak strength of leptoquark coupling are assumed. For leptoquark coupled to  $\overline{e}_L t_R$ ,  $\overline{\mu} t$ , and  $\overline{\tau} t$ , see Fig. 2 in BHATTACHARYYA 94B erratum and Fig. 3.
- <sup>34</sup> DAVIDSON 94 gives an extensive list of the bounds on leptoquark-induced four-fermion interactions from  $\pi$ , K, D, B,  $\mu$ ,  $\tau$  decays and meson mixings, *etc.* See Table 15 of DAVIDSON 94 for detail.
- <sup>35</sup> KUZNETSOV 94 gives mixing independent bound of the Pati-Salam leptoquark from the cosmological limit on  $\pi^0 \rightarrow \overline{\nu}\nu$ .
- <sup>36</sup> LEURER 94, LEURER 94B limits are obtained from atomic parity violation and apply to any chiral leptoquark which couples to the first generation with electromagnetic strength. For a nonchiral leptoquark, universality in  $\pi_{\ell 2}$  decay provides a much more stringent bound.
- $^{37}$  MAHANTA 94 gives bounds of *P* and *T*-violating scalar-leptoquark couplings from atomic and molecular experiments.
- <sup>38</sup> From  $(\pi \rightarrow e\nu)/(\pi \rightarrow \mu\nu)$  ratio. SHANKER 82 assumes the leptoquark induced four-fermion coupling  $4g^2/M^2$  ( $\overline{\nu}_{eL} u_R$ ) ( $\overline{d}_L e_R$ )with g=0.004 for spin-0 leptoquark and  $g^2/M^2$  ( $\overline{\nu}_{eL} \gamma_{\mu} u_L$ ) ( $\overline{d}_R \gamma^{\mu} e_R$ ) with  $g\simeq 0.6$  for spin-1 leptoquark.

MASS LIMITS f	MASS LIMITS for Diquarks					
VALUE (GeV)	CL%	DOCUMENT ID		TECN	COMMENT	
>6000 (CL = 95%	) OUR LII	TIN				
none 1500–6000	95	<sup>1</sup> KHACHATRY.	<b>16</b> K	CMS	E <sub>6</sub> diquark	
none 500–1600	95	<sup>2</sup> KHACHATRY.			E <sub>6</sub> diquark	
none 1200–4700	95	<sup>3</sup> KHACHATRY.	15V	CMS	E <sub>6</sub> diquark	
• • • We do not u	se the follo	owing data for average	ges, fits	s, limits	s, etc. ● ● ●	
>3750	95	<sup>4</sup> CHATRCHYAN			E <sub>6</sub> diquark	
none 1000–4280	95	<sup>5</sup> CHATRCHYAN			Superseded by KHACHA- TRYAN 15∨	
>3520	95	<sup>6</sup> CHATRCHYAN		CMS	Superseded by CHA- TRCHYAN 13A	
none 970–1080, 1450–1600	95	<sup>7</sup> KHACHATRY.	10	CMS	Superseded by CHA- TRCHYAN 13A	
none 290-630	95	<sup>8</sup> AALTONEN	<b>09</b> AC	CDF	E <sub>6</sub> diquark	
none 290-420	95	<sup>9</sup> ABE	<b>97</b> G	CDF	E <sub>6</sub> diquark	
none 15-31.7	95	<sup>10</sup> ABREU	940	DLPH	SUSY E <sub>6</sub> diquark	
<sup>1</sup> KHACHATRYAN 16K search for resonances decaying to dijets in $pp$ collisions at $\sqrt{s} = 13$ TeV. <sup>2</sup> KHACHATRYAN 16L search for resonances decaying to dijets in $pp$ collisions at $\sqrt{s} = 8$ TeV with the data scouting technique, increasing the sensitivity to the low mass resonances.						

- <sup>3</sup>KHACHATRYAN 15V search for resonances decaying to dijets in pp collisions at  $\sqrt{s} =$ 8 TeV.
- <sup>4</sup> CHATRCHYAN 13A search for new resonance decaying to dijets in pp collisions at  $\sqrt{s}$ = 7 TeV.
- $^5$  CHATRCHYAN 13AS search for new resonance decaying to dijets in pp collisions at  $\sqrt{s}$ = 8 TeV.
- $^{6}$ CHATRCHYAN 11Y search for new resonance decaying to dijets in pp collisions at  $\sqrt{s} = 7 \text{ TeV}.$
- <sup>7</sup>KHACHATRYAN 10 search for new resonance decaying to dijets in pp collisions at  $\sqrt{s} = 7$  TeV. <sup>8</sup> AALTONEN 09AC search for new narrow resonance decaying to dijets.
- $^9\,{\rm ABE}$  97G search for new particle decaying to dijets.
- <sup>10</sup>ABREU 940 limit is from  $e^+e^- \rightarrow \overline{cs}cs$ . Range extends up to 43 GeV if diquarks are degenerate in mass.

#### MASS LIMITS for $g_A$ (axigluon) and Other Color-Octet Gauge Bosons

Axigluons are massive color-octet gauge bosons in chiral color models and have axialvector coupling to quarks with the same coupling strength as gluons.

VALUE (GeV)	CL%	DOCUMENT ID	TECN	COMMENT
>5100 (CL = 95%)	our limi <sup>.</sup>	Г		
none 1500–5100	95	<sup>1</sup> KHACHATRY16K		$pp  ightarrow g_A X, g_A  ightarrow 2j$
none 500–1600	95	<sup>2</sup> KHACHATRY16L	CMS	$pp \rightarrow g_A X, g_A \rightarrow 2j$
none 1300–3600	95	<sup>3</sup> KHACHATRY15v	CMS	$pp  ightarrow g_{oldsymbol{A}} X$ , $g_{oldsymbol{A}}  ightarrow 2j$
$\bullet \bullet \bullet$ We do not use	the followi	ng data for averages,	fits, limit	s, etc. ● ● ●
				$pp \rightarrow g_A X, g_A \rightarrow b\overline{b}b\overline{b}$
>2800	95	<sup>5</sup> KHACHATRY16E	CMS	$pp \xrightarrow{\rightarrow} g_{KK} X, g_{KK} \rightarrow t\overline{t}$

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		<sup>6</sup> KHACHATRY15AV CMS $pp \rightarrow \Theta^0 \Theta^0 \rightarrow b\overline{b}Zg$
		<sup>7</sup> AALTONEN 13R CDF $p\overline{p} \rightarrow g_A X, g_A \rightarrow \sigma \sigma, \sigma \rightarrow 2j$
>3360	95	<sup>8</sup> CHATRCHYAN 13A CMS $pp \rightarrow g_A X, g_A \rightarrow 2j$
none 1000-3270	95	<sup>9</sup> CHATRCHYAN 13AS CMS Superseded by KHACHA- TRYAN 15V
none 250–740	95	<sup>10</sup> CHATRCHYAN 13AU CMS $pp \rightarrow 2g_A X, g_A \rightarrow 2j$
> 775	95	<sup>11</sup> ABAZOV 12R D0 $p \overline{p} \rightarrow g_A X, g_A \rightarrow t \overline{t}$
>2470	95	<sup>12</sup> CHATRCHYAN 11Y CMS Superseded by CHA- TRCHYAN 13A
		<sup>13</sup> AALTONEN 10L CDF $p\overline{p} \rightarrow g_A X, g_A \rightarrow t\overline{t}$
none 1470–1520	95	<sup>14</sup> KHACHATRY10 CMS Superseded by CHA- TRCHYAN 13a
none 260–1250	95	<sup>15</sup> AALTONEN 09AC CDF $p \overline{p} \rightarrow g_A X, g_A \rightarrow 2j$
> 910	95	<sup>16</sup> CHOUDHURY 07 RVUE $p \overline{p} \rightarrow t \overline{t} X$
> 365	95	<sup>17</sup> DONCHESKI 98 RVUE $\Gamma(Z \rightarrow hadron)$
none 200–980	95	<sup>18</sup> ABE 97G CDF $p \overline{p} \rightarrow g_A X, g_A \rightarrow 2j$
none 200–870	95	<sup>19</sup> ABE 95N CDF $p \overline{p} \rightarrow g_A X, g_A \rightarrow q \overline{q}$
none 240–640	95	<sup>20</sup> ABE 93G CDF $p \overline{p} \rightarrow g_A X, g_A \rightarrow 2j$
> 50	95	<sup>21</sup> CUYPERS 91 RVUE $\sigma(e^+e^- \rightarrow \text{hadrons})$
none 120–210	95	<sup>22</sup> ABE 90H CDF $p \overline{p} \rightarrow g_A X, g_A \rightarrow 2j$
> 29		<sup>23</sup> ROBINETT 89 THEO Partial-wave unitarity
none 150–310	95	<sup>24</sup> ALBAJAR 88B UA1 $p \overline{p} \rightarrow g_A X, g_A \rightarrow 2j$
> 20		BERGSTROM 88 RVUE $p \overline{p} \rightarrow \Upsilon X$ via $g_A g$
> 9		$^{25}$ CUYPERS 88 RVUE $\Upsilon$ decay
> 25		$^{26}$ DONCHESKI 88B RVUE $\Upsilon$ decay

<sup>1</sup>KHACHATRYAN 16K search for resonances decaying to dijets in pp collisions at  $\sqrt{s} =$  13 TeV.

<sup>2</sup> KHACHATRYAN 16L search for resonances decaying to dijets in pp collisions at  $\sqrt{s}$ = 8 TeV with the data scouting technique, increasing the sensitivity to the low mass resonances.

<sup>3</sup>KHACHATRYAN 15V search for resonances decaying to dijets in pp collisions at  $\sqrt{s} = 8$  TeV.

<sup>4</sup> AAD 16W search for a new resonance decaying to a pair of *b* and  $B_H$  in *pp* collisions at  $\sqrt{s} = 8$  TeV. The vector-like quark  $B_H$  is assumed to decay to bH. See their Fig. 3 and Fig. 4 for limits on  $\sigma \cdot B$ .

<sup>5</sup> KHACHATRYAN 16E search for KK gluon decaying to  $t\overline{t}$  in pp collisions at  $\sqrt{s} = 8$  reV.

<sup>6</sup> KHACHATRYAN 15AV search for pair productions of neutral color-octet weak-triplet scalar particles ( $\Theta^0$ ), decaying to  $b\overline{b}$ , Zg or  $\gamma g$ , in pp collisions at  $\sqrt{s} = 8$  TeV. The  $\Theta^0$  particle is often predicted in coloron (G', color-octet gauge boson) models and appear in the pp collisions through  $G' \rightarrow \Theta^0 \Theta^0$  decays. Assuming  $B(\Theta^0 \rightarrow b\overline{b}) = 0.5$ , they give limits  $m_{\Theta^0} > 623$  GeV (426 GeV) for  $m_{G'} = 2.3$   $m_{\Theta^0}$  ( $m_{G'} = 5$   $m_{\Theta^0}$ ).

<sup>7</sup> AALTONEN 13R search for new resonance decaying to  $\sigma\sigma$ , with hypothetical strongly interacting  $\sigma$  particle subsequently decaying to 2 jets, in  $p\overline{p}$  collisions at  $\sqrt{s} = 1.96$  TeV, using data corresponding to an integrated luminosity of 6.6 fb<sup>-1</sup>. For 50 GeV  $< m_{\sigma} < m_{g_A}/2$ , axigluons in mass range 150–400 GeV are excluded.

<sup>8</sup>CHATRCHYAN 13A search for new resonance decaying to dijets in pp collisions at  $\sqrt{s}$  $_{o} = 7$  TeV.

<sup>9</sup>CHATRCHYAN 13AS search for new resonance decaying to dijets in pp collisions at  $\sqrt{s}$  $_{a} = 8$  TeV.

<sup>10</sup> CHATRCHYAN 13AU search for the pair produced color-octet vector bosons decaying to  $q\overline{q}$  pairs in pp collisions. The quoted limit is for  $B(g_A \rightarrow q\overline{q}) = 1$ .

- <sup>11</sup>ABAZOV 12R search for massive color octet vector particle decaying to  $t\bar{t}$ . The quoted limit assumes  $g_A$  couplings with light quarks are suppressed by 0.2.
- <sup>12</sup> CHATRCHYAN 11Y search for new resonance decaying to dijets in pp collisions at  $\sqrt{s} = 7$  TeV.
- <sup>13</sup>ÅALTONEN 10L search for massive color octet non-chiral vector particle decaying into  $t\bar{t}$  pair with mass in the range 400 GeV < M < 800 GeV. See their Fig. 6 for limit in the mass-coupling plane.
- <sup>14</sup> KHACHATRYAN 10 search for new resonance decaying to dijets in pp collisions at  $\sqrt{s} = 7$  TeV.
- <sup>15</sup> AALTONEN 09AC search for new narrow resonance decaying to dijets.
- <sup>16</sup> CHOUDHURY 07 limit is from the  $t\bar{t}$  production cross section measured at CDF.
- <sup>17</sup> DONCHESKI 98 compare  $\alpha_s$  derived from low-energy data and that from  $\Gamma(Z \rightarrow hadrons)/\Gamma(Z \rightarrow leptons)$ .
- $^{18}$  ABE 97G search for new particle decaying to dijets.
- $^{19}$ ABE 95N assume axigluons decaying to quarks in the Standard Model only.
- <sup>20</sup>ABE 93G assume  $\Gamma(g_A) = N \alpha_s m_{g_A}/6$  with N = 10.
- $^{21}\,\rm CUYPERS$  91 compare  $\alpha_s$  measured in  $\,\Upsilon$  decay and that from R at PEP/PETRA energies.
- <sup>22</sup> ABE 90H assumes  $\Gamma(g_A) = N\alpha_s m_{g_A}/6$  with N = 5 ( $\Gamma(g_A) = 0.09 m_{g_A}$ ). For N = 10, the excluded region is reduced to 120–150 GeV.
- <sup>23</sup> ROBINETT 89 result demands partial-wave unitarity of J = 0  $t\bar{t} \rightarrow t\bar{t}$  scattering amplitude and derives a limit  $m_{g_A} > 0.5 m_t$ . Assumes  $m_t > 56$  GeV.
- <sup>24</sup> ALBAJAR 88B result is from the nonobservation of a peak in two-jet invariant mass distribution.  $\Gamma(g_A) < 0.4 m_{g_A}$  assumed. See also BAGGER 88.
- <sup>25</sup> CUYPERS 88 requires  $\Gamma(\Upsilon \rightarrow gg_A) < \Gamma(\Upsilon \rightarrow ggg)$ . A similar result is obtained by CONCHESKI 88.
- <sup>26</sup> DONCHESKI 88B requires  $\Gamma(\Upsilon \rightarrow g q \overline{q})/\Gamma(\Upsilon \rightarrow g g g) < 0.25$ , where the former decay proceeds via axigluon exchange. A more conservative estimate of < 0.5 leads to  $m_{g_A} > 21$  GeV.

#### MASS LIMITS for Color-Octet Scalar Bosons

VALUE (GeV)	CL%	DOCUMENT ID	TECN	COMMENT			
• • • We do not use	ullet $ullet$ $ullet$ We do not use the following data for averages, fits, limits, etc. $ullet$ $ullet$ $ullet$						
none 150–287	95	<sup>1</sup> KHACHATRY. <sup>2</sup> AAD		$\begin{array}{rcl} p  p  \to & \Theta^0  \Theta^0  \to & b  \overline{b}  Z  g \\ p  p  \to & S_8  S_8  X, S_8  \to  2 \text{ jets} \end{array}$			
<sup>1</sup> KHACHATRYAN 15AV search for pair productions of neutral color-octet weak-triplet scalar particles ( $\Theta^0$ ), decaying to $b\overline{b}$ , $Zg$ or $\gamma g$ , in $pp$ collisions at $\sqrt{s} = 8$ TeV. The							
$\Theta^0$ particle is often predicted in coloron ( $G'$ , color-octet gauge boson) models and appear in the $pp$ collisions through $G' \to \Theta^0 \Theta^0$ decays. Assuming $B(\Theta^0 \to b\overline{b}) = 0.5$ , they give limits $m_{\Theta^0} > 623$ GeV (426 GeV) for $m_{G'} = 2.3 m_{\Theta^0} (m_{G'} = 5 m_{\Theta^0})$ .							
<sup>2</sup> AAD 13K search for pair production of color-octet scalar particles in $pp$ collisions at $\sqrt{s}$ = 7 TeV. Cross section limits are interpreted as mass limits on scalar partners of a Dirac							

gluino.

#### $X^0$ (Heavy Boson) Searches in Z Decays

Searches for radiative transition of Z to a lighter spin-0 state  $X^0$  decaying to hadrons, a lepton pair, a photon pair, or invisible particles as shown in the comments. The limits are for the product of branching ratios.

VALUE	<u>CL%</u>	DOCUMENT ID		TECN	COMMENT
• • • We do not u	use the fol	lowing data for ave	rages,	fits, lim	its, etc. ● ● ●
		<sup>1</sup> BARATE	<b>98</b> U	ALEP	$X^0 \rightarrow \ell \overline{\ell},  q \overline{q},  g g,  \gamma \gamma,  \nu \overline{\nu}$
		<sup>2</sup> ACCIARRI	97Q	L3	$X^0  ightarrow$ invisible particle(s)
		<sup>3</sup> ACTON	93E	OPAL	$X^0 \rightarrow \gamma \gamma$
		<sup>4</sup> ABREU	<b>92</b> D	DLPH	$X^0  ightarrow$ hadrons
		<sup>5</sup> ADRIANI	92F	L3	$X^0  ightarrow$ hadrons
		<sup>6</sup> ACTON	91	OPAL	$X^0  ightarrow$ anything
$< 1.1  imes 10^{-4}$	95	<sup>7</sup> ACTON	<b>91</b> B	OPAL	$X^0 \rightarrow e^+ e^-$
$< 9 \times 10^{-5}$	95	<sup>7</sup> ACTON	<b>91</b> B	OPAL	$X^0 \rightarrow \mu^+ \mu^-$
$< 1.1  imes 10^{-4}$	95	<sup>7</sup> ACTON	<b>91</b> B	OPAL	$X^0 \rightarrow \tau^+ \tau^-$
$< 2.8  imes 10^{-4}$	95	<sup>8</sup> ADEVA	<b>91</b> D	L3	$X^0 \rightarrow e^+ e^-$
$<\!\!2.3  imes 10^{-4}$	95	<sup>8</sup> ADEVA	<b>91</b> D	L3	$X^0 \rightarrow \mu^+ \mu^-$
$< 4.7 \times 10^{-4}$	95	<sup>9</sup> ADEVA	<b>91</b> D	L3	$X^0 \rightarrow$ hadrons
$< 8 \times 10^{-4}$	95	<sup>10</sup> AKRAWY	<b>90</b> J	OPAL	$X^0  ightarrow$ hadrons

<sup>1</sup>BARATE 980 obtain limits on  $B(Z \rightarrow \gamma X^0)B(X^0 \rightarrow \ell \overline{\ell}, q \overline{q}, g g, \gamma \gamma, \nu \overline{\nu})$ . See their Fig. 17.

<sup>2</sup>See Fig. 4 of ACCIARRI 97Q for the upper limit on B( $Z \rightarrow \gamma X^0$ ;  $E_{\gamma} > E_{min}$ ) as a function of Emin.

<sup>3</sup>ACTON 93E give  $\sigma(e^+e^- \rightarrow X^0\gamma) \cdot B(X^0 \rightarrow \gamma\gamma) < 0.4 \text{ pb } (95\%\text{CL}) \text{ for } m_{\chi^0} = 60 \pm 10^{-10} \text{ m}$ 2.5 GeV. If the process occurs via s-channel  $\gamma$  exchange, the limit translates to  $\Gamma(X^0) \cdot B(X^0 \rightarrow \gamma \gamma)^2 < 20 \text{ MeV for } m_{X^0} = 60 \pm 1 \text{ GeV.}$ <sup>4</sup>ABREU 92D give  $\sigma_Z \cdot B(Z \rightarrow \gamma X^0) \cdot B(X^0 \rightarrow \text{ hadrons}) < (3-10) \text{ pb for } m_{X^0} =$ 

10–78 GeV. A very similar limit is obtained for spin-1  $X^0$ .

<sup>5</sup> ADRIANI 92F search for isolated  $\gamma$  in hadronic Z decays. The limit  $\sigma_{Z} \cdot B(Z \rightarrow \gamma X^{0})$  $\cdot$  B(X^0 \rightarrow \, hadrons) <(2–10) pb (95%CL) is given for  $m_{X^0} =$  25–85 GeV.

<sup>6</sup> ACTON 91 searches for  $Z \rightarrow Z^* X^0$ ,  $Z^* \rightarrow e^+ e^-$ ,  $\mu^+ \mu^-$ , or  $\nu \overline{\nu}$ . Excludes any new scalar  $X^0$  with  $m_{\chi^0} <$  9.5 GeV/c if it has the same coupling to  $ZZ^*$  as the MSM Higgs boson.

<sup>7</sup> ACTON 91B limits are for  $m_{\chi 0} = 60-85$  GeV.

<sup>8</sup> ADEVA 91D limits are for  $m_{\chi 0} = 30-89$  GeV.

<sup>9</sup>ADEVA 91D limits are for  $m_{\chi 0} = 30-86$  GeV.

<sup>10</sup> AKRAWY 90J give  $\Gamma(Z \rightarrow \gamma X^0) \cdot B(X^0 \rightarrow \text{hadrons}) < 1.9 \text{ MeV}$  (95%CL) for  $m_{\chi 0}$ = 32-80 GeV. We divide by  $\Gamma(Z) = 2.5$  GeV to get product of branching ratios. For nonresonant transitions, the limit is  $B(Z \rightarrow \gamma q \overline{q}) < 8.2$  MeV assuming three-body phase space distribution.

MASS LIMITS	<b>5 for a l</b>	leavy Neutral Be DOCUMENT ID	oson	Couplin TECN	р <b>g to e<sup>+</sup>e<sup>-</sup></b> соммент
• • • We do not	use the	following data for a	average	es, fits, l	imits, etc. ● ● ●
none 55–61		<sup>1</sup> ODAKA	89	VNS	$\Gamma(X^0 \rightarrow e^+ e^-)$
					$B(X^0 \rightarrow had.) \gtrsim 0.2 \text{ MeV}$
>45	95	<sup>2</sup> DERRICK	86	HRS	$\Gamma(X^0 \rightarrow e^+e^-) = 6 \text{ MeV}$
>46.6	95	<sup>3</sup> ADEVA	85	MRKJ	$\Gamma(X^0 \rightarrow e^+e^-)=10 \text{ keV}$
>48	95	<sup>3</sup> ADEVA	85	MRKJ	$\Gamma(X^0 \rightarrow e^+ e^-) = 4 \text{ MeV}$
		<sup>4</sup> BERGER	<b>85</b> B	PLUT	
none 39.8–45.5		<sup>5</sup> ADEVA	84	MRKJ	$\Gamma(X^0  ightarrow e^+ e^-) = 10 \; { m keV}$
>47.8	95	<sup>5</sup> ADEVA	84	MRKJ	$\Gamma(X^0 \rightarrow e^+ e^-) = 4 \text{ MeV}$
none 39.8–45.2		<sup>5</sup> BEHREND	84C	CELL	
>47	95	<sup>5</sup> BEHREND	84C	CELL	$\Gamma(X^0 \rightarrow e^+e^-)=4 \text{ MeV}$

<sup>1</sup>ODAKA 89 looked for a narrow or wide scalar resonance in  $e^+e^- \rightarrow$  hadrons at  $E_{cm}$ = 55.0–60.8 GeV.  $^2$  DERRICK 86 found no deviation from the Standard Model Bhabha scattering at  $E_{\rm cm}=$ 

29 GeV and set limits on the possible scalar boson  $e^+e^-$  coupling. See their figure 4 for excluded region in the  $\Gamma(X^0 \rightarrow e^+e^-) \cdot m_{\chi 0}$  plane. Electronic chiral invariance requires a parity doublet of  $X^0$ , in which case the limit applies for  $\Gamma(X^0 \rightarrow e^+ e^-) =$ 

<sup>3</sup> MeV. <sup>3</sup> ADEVA 85 first limit is from  $2\gamma$ ,  $\mu^+\mu^-$ , hadrons assuming  $X^0$  is a scalar. Second limit is from  $e^+e^-$  channel.  $E_{cm} = 40-47$  GeV. Supersedes ADEVA 84.

<sup>4</sup> BERGER 85B looked for effect of spin-0 boson exchange in  $e^+e^- \rightarrow e^+e^-$  and  $\mu^+\mu^$ at  $E_{\rm cm} = 34.7$  GeV. See Fig. 5 for excluded region in the  $m_{\chi^0} - \Gamma(X^0)$  plane.

 $^5$  ADEVA 84 and BEHREND 84C have  $E_{
m cm}=$  39.8–45.5 GeV. MARK-J searched  $X^0$  in  $e^+e^- \rightarrow$  hadrons,  $2\gamma$ ,  $\mu^+\mu^-$ ,  $e^+e^-$  and CELLO in the same channels plus  $\tau$  pair. No narrow or broad  $X^0$  is found in the energy range. They also searched for the effect of  $X^0$  with  $m_X > E_{cm}$ . The second limits are from Bhabha data and for spin-0 singlet. The same limits apply for  $\Gamma(X^0 \rightarrow e^+e^-) = 2$  MeV if  $X^0$  is a spin-0 doublet. The second limit of BEHREND 84C was read off from their figure 2. The original papers also list limits in other channels.

Search for  $X^0$  Resonance in  $e^+e^-$  Collisions The limit is for  $\Gamma(X^0 \rightarrow e^+e^-) \cdot B(X^0 \rightarrow f)$ , where f is the specified final state. Spin 0 is assumed for  $X^0$ .

VALUE (keV)	CL%	DOCUMENT ID		TECN	COMMENT
• • • We do not use t	he follow	ing data for averages	s, fits,	limits, e	etc. • • •
<10 <sup>3</sup>	95	<sup>1</sup> ABE	<b>93</b> C	VNS	Γ(ee)
<(0.4–10)	95	<sup>2</sup> ABE	<b>93</b> C	VNS	$f = \gamma \gamma$
<(0.3–5)	95	<sup>3,4</sup> ABE	<b>93</b> D	TOPZ	$f = \gamma \gamma$
<(2–12)	95	<sup>3,4</sup> ABE	<b>93</b> D	TOPZ	f = hadrons
<(4–200)	95	<sup>4,5</sup> ABE	<b>93</b> D	TOPZ	f = e e
<(0.1–6)	95	<sup>4,5</sup> ABE	<b>93</b> D	TOPZ	$f = \mu \mu$
<(0.5–8)	90	<sup>6</sup> STERNER	93	AMY	$f = \gamma \gamma$

<sup>1</sup>Limit is for  $\Gamma(X^0 \rightarrow e^+e^-) m_{\chi^0} =$  56–63.5 GeV for  $\Gamma(X^0) =$  0.5 GeV.

<sup>2</sup>Limit is for  $m_{\chi 0} = 56-61.5$  GeV and is valid for  $\Gamma(X^0) \ll 100$  MeV. See their Fig. 5 for limits for  $\Gamma = 1,2$  GeV. <sup>3</sup>Limit is for  $m_{\chi^0} = 57.2-60$  GeV.

<sup>4</sup>Limit is valid for  $\Gamma(X^0) \ll 100$  MeV. See paper for limits for  $\Gamma = 1$  GeV and those for  $5 \text{ Limit is for } m_{\chi^0} = 56.6-60 \text{ GeV}.$ 

 $^6$  STERNER 93 limit is for  $m_{\chi 0}$  = 57–59.6 GeV and is valid for  $\Gamma(X^0){<}100$  MeV. See their Fig. 2 for limits for  $\Gamma = 1,3$  GeV.

#### Search for $X^0$ Resonance in ep Collisions

VALUE	DOCUMENT ID	TECN	<u>COMMENT</u>	
• • • We do not use the followin	g data for average	s, fits, limits,	etc. • • •	
1	$^1$ CHEKANOV	02B ZEUS	$X \rightarrow jj$	

<sup>1</sup>CHEKANOV 02B search for photoproduction of X decaying into dijets in *ep* collisions. See their Fig. 5 for the limit on the photoproduction cross section.

#### Search for $X^0$ Resonance in $e^+e^- \rightarrow X^0\gamma$

VALUE (GeV)	DOCUMENT ID	·	TECN	COMMENT
$\bullet \bullet \bullet$ We do not use the follow	ing data for average	es, fits,	limits, e	etc. • • •
	<sup>1</sup> ABBIENDI <sup>2</sup> ABREU <sup>3</sup> ADAM	00Z	DLPH	$egin{array}{lll} X^{m 0} &  ightarrow ~ \gamma \gamma \ X^{m 0} & { m decaying invisibly} \ X^{m 0} & { m decaying invisibly} \end{array}$
$^1$ ABBIENDI 03D measure th	$e e^+e^- \rightarrow \gamma\gamma\gamma\gamma$	cross s	ection a	t $\sqrt{s}$ =181–209 GeV. The

upper bound on the production cross section,  $\sigma(e^+e^- \rightarrow X^0\gamma)$  times the branching ratio for  $X^0 \rightarrow \gamma \gamma$ , is less than 0.03 pb at 95%CL for  $X^0$  masses between 20 and 180 GeV. See their Fig. 9b for the limits in the mass-cross section plane.

<sup>2</sup>ABREU 00Z is from the single photon cross section at  $\sqrt{s}$ =183, 189 GeV. The production cross section upper limit is less than 0.3 pb for  $X^0$  mass between 40 and 160 GeV. See their Fig. 4 for the limit in mass-cross section plane.

<sup>3</sup>ADAM 96C is from the single photon production cross at  $\sqrt{s}$ =130, 136 GeV. The upper bound is less than 3 pb for  $X^0$  masses between 60 and 130 GeV. See their Fig. 5 for the exact bound on the cross section  $\sigma(e^+e^- \rightarrow \gamma X^0)$ .

#### Search for $X^0$ Resonance in $Z \rightarrow f \overline{f} X_0^0$

The limit is for ${\sf B}(Z  o$	$f \overline{f} X^0$ · B( $X^0 \rightarrow$	F) where f	is a fermio	n and F	is the
specified final state. Spin	0 is assumed for $X^0$				

CL%	DOCUMENT ID		TECN	COMMENT
following	data for averages	, fits,	limits, e	etc. • • •
	<sup>1</sup> ABREU	96⊤	DLPH	f=e, $\mu$ , $\tau$ ; F= $\gamma\gamma$
95	<sup>2</sup> ABREU	<b>96</b> ⊤	DLPH	$f=\nu; F=\gamma\gamma$
		96T	DLPH	$f=q; F=\gamma\gamma$
95		93E	OPAL	$f = e, \mu, \tau; F = \gamma \gamma$
95		93E	OPAL	$f = q; F = \gamma \gamma$
95		93E	OPAL	$f = \nu; F = \gamma \gamma$
95		93E	OPAL	$f = e, \mu; F = \ell \overline{\ell}, q \overline{q}, \nu \overline{\nu}$
95		93F	ALEP	$f = e, \mu; F = \ell \overline{\ell}, q \overline{q}, \nu \overline{\nu}$
	<sup>4</sup> ADRIANI	92F	L3	$f = q; F = \gamma \gamma$
	following 95 95 95 95 95	following data for averages 1 ABREU 95 2 ABREU 3 ABREU 95 2 ACTON 95 2 ACTON 95 2 ACTON 95 2 ACTON 95 2 ACTON	following         data for averages, fits,           1         ABREU         96T           95         2         ABREU         96T           3         ABREU         96T           95         2         ACTON         93E           95         2         BUSKULIC         93F	following data for averages, fits, limits, e 1 ABREU 96T DLPH 95 2 ABREU 96T DLPH 3 ABREU 96T DLPH 95 2 ACTON 93E OPAL 95 2 ACTON 93E OPAL 95 2 ACTON 93E OPAL 95 2 ACTON 93E OPAL 95 2 BUSKULIC 93F ALEP

- <sup>1</sup>ABREU 96T obtain limit as a function of  $m_{\chi 0}$ . See their Fig. 6.
- <sup>2</sup>Limit is for  $m_{\chi^0}$  around 60 GeV.
- <sup>3</sup>ABREU 96T obtain limit as a function of  $m_{\chi^0}$ . See their Fig. 15.

<sup>4</sup> ADRIANI 92F give  $\sigma_Z \cdot B(Z \rightarrow q \overline{q} X^0) \cdot B(X^0 \rightarrow \gamma \gamma) < (0.75-1.5) \text{ pb } (95\% \text{CL}) \text{ for } m_{\chi 0} = 10-70 \text{ GeV}$ . The limit is 1 pb at 60 GeV.

#### Search for $X^0$ Resonance in $WX^0$ final state

VALUE (MeV)	DOCUMENT ID	TECN	COMMENT	
$\bullet$ $\bullet$ We do not use the follo	wing data for average	es, fits, limits,	etc. • • •	
	<sup>1</sup> AALTONEN			
	<sup>2</sup> CHATRCHYA			
	<sup>3</sup> ABAZOV			
	<sup>4</sup> ABE	97W CDF	$X^0 \rightarrow b\overline{b}$	
<sup>1</sup> AALTONEN 13AA search	for $X^0$ production as	sociated with	W (or Z) in $p\overline{p}$ collis	sions

at  $E_{\rm cm} = 1.96$  TeV. The upper limit on the cross section  $\sigma(p\overline{p} \rightarrow WX^0)$  is 2.2 pb for  $M_{\chi^0} = 145$  GeV.

- <sup>2</sup> CHATRCHYAN 12BR search for  $X^0$  production associated with W in pp collisions at  $E_{\rm cm} = 7$  TeV. The upper limit on the cross section is 5.0 pb at 95% CL for  $m_{\chi^0} = 150$  GeV.
- <sup>3</sup>ABAZOV 111 search for  $X^0$  production associated with W in  $p\overline{p}$  collisions at  $E_{\rm cm} = 1.96$  TeV. The 95% CL upper limit on the cross section ranges from 2.57 to 1.28 pb for  $X^0$  mass between 110 and 170 GeV.
- $X^0$  mass between 110 and 170 GeV. <sup>4</sup>ABE 97W search for  $X^0$  production associated with W in  $p\overline{p}$  collisions at  $E_{\rm cm}$ =1.8 TeV. The 95%CL upper limit on the production cross section times the branching ratio for  $X^0 \rightarrow b\overline{b}$  ranges from 14 to 19 pb for  $X^0$  mass between 70 and 120 GeV. See their Fig. 3 for upper limits of the production cross section as a function of  $m_{\chi 0}$ .

#### Search for $X^0$ Resonance in Quarkonium Decays

<sup>1</sup> BALEST 95 three-body limit is for phase-space photon energy distribution and angular distribution same as for  $\Upsilon \rightarrow gg\gamma$ .

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ABAZOV ABAZOV	08AA 08AD	PL B668 357	V.M. Abazov <i>et al.</i>	(D0 Collab.)

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MACDONALD	08	PR D78 032010	R.P. MacDonald <i>et al.</i>	(TWIST Collab.)
ZHANG	08	NP B802 247	Y. Zhang et al.	(PKGU, UMD)
AALTONEN	07H	PRL 99 171802	T. Aaltonen <i>et al.</i>	(CDF Collab.)
ABAZOV	07E	PL B647 74	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	07J	PRL 99 061801	V.M. Abazov <i>et al.</i>	(D0 Collab.)
AKTAS	07A	EPJ C52 833	A. Aktas <i>et al.</i>	(H1 Collab.)
CHOUDHURY	07	PL B657 69	D. Choudhury <i>et al.</i>	( )
MELCONIAN	07	PL B649 370	D. Melconian <i>et al.</i>	(TRIUMF)
SCHAEL	07A	EPJ C49 411	S. Schael <i>et al.</i>	(ALEPH Collab.)
SCHUMANN	07	PRL 99 191803	M. Schumann <i>et al.</i>	(HEID, ILLG, KARL+)
SMIRNOV	07	MPL A22 2353	A.D. Smirnov	
ABAZOV	06A	PL B636 183	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	06L	PL B640 230	V.M. Abazov et al.	(D0 Collab.)
ABDALLAH	06C	EPJ C45 589	J. Abdallah <i>et al.</i>	(DELPHI Collab.)
ABULENCIA	06L	PRL 96 211801	A. Abulencia <i>et al.</i>	(CDF Collab.)
ABULENCIA	06M	PRL 96 211802	A. Abulencia <i>et al.</i>	(CDF Collab.)
ABULENCIA	06T	PR D73 051102	A. Abulencia <i>et al.</i>	(CDF Collab.)
ABAZOV	05H	PR D71 071104	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABULENCIA	05A	PRL 95 252001	A. Abulencia <i>et al.</i>	(CDF Collab.)
ACOSTA	051	PR D71 112001	D. Acosta <i>et al.</i>	(CDF Collab.)
ACOSTA	05P	PR D72 051107	D. Acosta <i>et al.</i>	(CDF Collab.)
ACOSTA	05R	PRL 95 131801	D. Acosta <i>et al.</i>	(CDF Collab.)
AKTAS	05B	PL B629 9	A. Aktas <i>et al.</i>	(H1 Collab.)
CHEKANOV	05	PL B610 212	S. Chekanov <i>et al.</i>	(HERA ZEUS Collab.)
CHEKANOV	05A	EPJ C44 463	S. Chekanov <i>et al.</i>	(ZEUS Collab.)
CYBURT	05	ASP 23 313	R.H. Cyburt <i>et al.</i>	(2200 00000)
ABAZOV	04A	PRL 92 221801	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	04C	PR D69 111101	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABBIENDI	04G	EPJ C33 173	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBIENDI	03D	EPJ C26 331	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBIENDI	03R	EPJ C31 281	G. Abbiendi <i>et al.</i>	(OPAL)
ACOSTA	03B	PRL 90 081802	D. Acosta <i>et al.</i>	(CDF Collab.)
ADLOFF	03	PL B568 35	C. Adloff <i>et al.</i>	(H1 Collab.)
BARGER	03B	PR D67 075009	V. Barger, P. Langacker, H. Lee	(
CHANG	03	PR D68 111101	MC. Chang <i>et al.</i>	(BELLE Collab.)
CHEKANOV	03B	PR D68 052004	S. Chekanov <i>et al.</i>	(ZEUS Collab.)
ABAZOV	02	PRL 88 191801	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABBIENDI	02B	PL B526 233	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
AFFOLDER	02C	PRL 88 071806	T. Affolder <i>et al.</i>	(CDF Collab.)
CHEKANOV	02	PR D65 092004	S. Chekanov <i>et al.</i>	(ZEUS Collab.)
CHEKANOV	02B	PL B531 9	S. Chekanov <i>et al.</i>	(ZEUS Collab.)
MUECK	02	PR D65 085037	A. Mueck, A. Pilaftsis, R. Rueckl	
ABAZOV	01B	PRL 87 061802	V.M. Abazov et al.	(D0 Collab.)
ABAZOV	01D	PR D64 092004	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ADLOFF	01C	PL B523 234	C. Adloff <i>et al.</i>	(H1 Collab.)
AFFOLDER	011	PRL 87 231803	T. Affolder <i>et al.</i>	(CDF Collab.)
BREITWEG	01	PR D63 052002	J. Breitweg <i>et al.</i>	(ŻEUS Collab.)
CHEUNG	01B	PL B517 167	K. Cheung	, , , , , , , , , , , , , , , , , , ,
THOMAS	01	NP A694 559	E. Thomas <i>et al.</i>	
ABBIENDI	00M	EPJ C13 15	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBOTT	00C	PRL 84 2088	B. Abbott <i>et al.</i>	(D0 Collab.)
ABE	00	PRL 84 5716	F. Abe <i>et al.</i>	(CDF Collab.)
ABREU	00S	PL B485 45	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	00Z	EPJ C17 53	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ACCIARRI	00P	PL B489 81	M. Acciarri <i>et al.</i>	(L3 Collab.)
ADLOFF	00	PL B479 358	C. Adloff <i>et al.</i>	(H1 Collab.)
AFFOLDER	00K	PRL 85 2056	T. Affolder <i>et al.</i>	(CDF Collab.)
BARATE	001	EPJ C12 183	R. Barate <i>et al.</i>	(ALEPH Collab.)
BARGER	00	PL B480 149	V. Barger, K. Cheung	
BREITWEG	00E	EPJ C16 253	J. Breitweg <i>et al.</i>	(ZEUS Collab.)
CHAY	00	PR D61 035002	J. Chay, K.Y. Lee, S. Nam	
СНО	00	MPL A15 311	G. Cho	
CORNET	00	PR D61 037701	F. Cornet, M. Relano, J. Rico	
DELGADO	00	JHEP 0001 030	A. Delgado, A. Pomarol, M. Quir	OS
ERLER	00	PRL 84 212	J. Erler, P. Langacker	
GABRIELLI	00	PR D62 055009	E. Gabrielli	
RIZZO	00	PR D61 016007	T.G. Rizzo, J.D. Wells	
ROSNER	00	PR D61 016006	J.L. Rosner	
ZARNECKI	00	EPJ C17 695	A. Zarnecki	
ABBIENDI	99	EPJ C6 1	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBOTT	99J	PRL 83 2896	B. Abbott <i>et al.</i>	(D0 Collab.)

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ABREU 99G PL B446 62 P. Abreu <i>et al.</i> (DELPH ACKERSTAFF 99D EPJ C8 3 K. Ackerstaff <i>et al.</i> (OPA	
	II Collab )
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CASALBUONI 99 PL B460 135 R. Casalbuoni <i>et al.</i>	I Collab.)
ERLER 99 PL B456 68 J. Erler, P. Langacker	
MARCIANO 99 PR D60 093006 W. Marciano	
MASIP 99 PR D60 096005 M. Masip, A. Pomarol	
NATH 99 PR D60 116004 P. Nath, M. Yamaguchi	
STRUMIA 99 PL B466 107 A. Strumia	
	0 Collab.)
	0 Collab.)
	F Collab.)
	F Collab.)
ACCIARRI 98J PL B433 163 M. Acciarri <i>et al.</i> (L	.3 Collab.)
ACKERSTAFF 98V EPJ C2 441 K. Ackerstaff <i>et al.</i> (OPA	L Collab.)
BARATE 98U EPJ C4 571 R. Barate <i>et al.</i> (ALEPI	H Collab.)
BARENBOIM 98 EPJ C1 369 G. Barenboim	
CHO 98 EPJ C5 155 G. Cho, K. Hagiwara, S. Matsumoto	
CONRAD 98 RMP 70 1341 J.M. Conrad, M.H. Shaevitz, T. Bolton	
DONCHESKI 98 PR D58 097702 M.A. Doncheski, R.W. Robinett	
GROSS-PILCH98 hep-ex/9810015 C. Grosso-Pilcher, G. Landsberg, M. Paterno	
	F Collab.)
ABE 97G PR D55 R5263 F. Abe <i>et al.</i> (CD	F Collab.)
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DEANDREA 97 PL B409 277 A. Deandrea	(MARS)
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STAHL 97 ZPHY C74 73 A. Stahl, H. Voss	< · · · · · · · · · · · · · · · ·
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ALLET 96 PL B383 139 M. Allet <i>et al.</i> (VILL, LEUV, LOU	
	0 Collab.)
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	O Collab.)
KUZNETSOV 95 PRL 75 794 I.A. Kuznetsov <i>et al.</i> (PNPI, KIAE,	
KUZNETSOV 95B PAN 58 2113 A.V. Kuznetsov, N.V. Mikheev	(YARO)
Translated from YAF 58 2228. MIZUKOSHI 95 NP B443 20 J.K. Mizukoshi, O.J.P. Eboli, M.C. Gonzalez-Ga	rcia
APPELL 040 7DUV C64 192 D Abrou at al (DELDL	(6==)
ABREU 940 ZPHY C64 183 P. Abreu et al. (DELPH	(CERN)
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DERRICK	93	PL B306 173	M. Derrick et al. (ZEUS Collab.)
RIZZO	93	PR D48 4470	T.G. Rizzo (ANL)
SEVERIJNS	93	PRL 70 4047	N. Severijns <i>et al.</i> (LOUV, WISC, LEUV+)
Also	50	PRL 73 611 (erratum)	N. Severijns <i>et al.</i> (LOUV, WISC, LEUV+)
STERNER	93	PL B303 385	K.L. Sterner <i>et al.</i> (AMY Collab.)
ABREU	92D	ZPHY C53 555	P. Abreu <i>et al.</i> (DELPHI Collab.)
ADRIANI	92D 92F	PL B292 472	O. Adriani <i>et al.</i> (L3 Collab.)
DECAMP	921	PRPL 216 253	D. Decamp <i>et al.</i> (ALEPH Collab.)
IMAZATO	92 92	PRL 69 877	
MISHRA	92 92		
-	-	PRL 68 3499	
POLAK	92B	PR D46 3871	J. Polak, M. Zralek (SILES)
ACTON	91 01 D	PL B268 122	D.P. Acton <i>et al.</i> (OPAL Collab.)
ACTON	91B	PL B273 338	D.P. Acton <i>et al.</i> (OPAL Collab.)
ADEVA	91D	PL B262 155	B. Adeva <i>et al.</i> (L3 Collab.)
AQUINO	91	PL B261 280	M. Aquino, A. Fernandez, A. Garcia (CINV, PUEB)
COLANGELO	91	PL B253 154	P. Colangelo, G. Nardulli (BARI)
CUYPERS	91	PL B259 173	F. Cuypers, A.F. Falk, P.H. Frampton (DURH, HARV+)
FARAGGI	91	MPL A6 61	A.E. Faraggi, D.V. Nanopoulos (TAMU)
POLAK	91	NP B363 385	J. Polak, M. Zralek (SILES)
RIZZO	91	PR D44 202	T.G. Rizzo (WISC, ISU)
WALKER	91	APJ 376 51	T.P. Walker <i>et al.</i> (HSCA, OSU, CHIC+)
ABE	90F	PL B246 297	K. Abe <i>et al.</i> (VENUS Collab.)
ABE	90H	PR D41 1722	F. Abe <i>et al.</i> (CDF Collab.)
AKRAWY	90 J	PL B246 285	M.Z. Akrawy et al. (OPAL Collab.)
GONZALEZ	90D	PL B240 163	M.C. Gonzalez-Garcia, J.W.F. Valle (VALE)
GRIFOLS	90	NP B331 244	J.A. Grifols, E. Masso (BARC)
GRIFOLS	90D	PR D42 3293	J.A. Grifols, E. Masso, T.G. Rizzo (BARC, CERN+)
KIM	90	PL B240 243	G.N. Kim <i>et al.</i> (AMY Collab.)
LOPEZ	90	PL B241 392	J.L. Lopez, D.V. Nanopoulos (TAMU)
BARBIERI	89B	PR D39 1229	R. Barbieri, R.N. Mohapatra (PISA, UMD)
LANGACKER	89B	PR D40 1569	P. Langacker, S. Uma Sankar (PENN)
ODAKA	89	JPSJ 58 3037	S. Odaka <i>et al.</i> (VENUS Collab.)
ROBINETT	89	PR D39 834	R.W. Robinett (PSU)
ALBAJAR	88B	PL B209 127	C. Albajar <i>et al.</i> (UA1 Collab.)
BAGGER	88	PR D37 1188	J. Bagger, C. Schmidt, S. King (HARV, BOST)
BALKE	88	PR D37 587	B. Balke et al. (LBL, UCB, COLO, NWES+)
BERGSTROM	88	PL B212 386	L. Bergstrom (STOH)
CUYPERS	88	PRL 60 1237	F. Cuypers, P.H. Frampton (UNCCH)
DONCHESKI	88	PL B206 137	M.A. Doncheski, H. Grotch, R. Robinett (PSU)
DONCHESKI	88B	PR D38 412	M.A. Doncheski, H. Grotch, R.W. Robinett (PSU)
BARTEL	87B	ZPHY C36 15	W. Bartel et al. (JADE Collab.)
BEHREND	86B	PL B178 452	H.J. Behrend <i>et al.</i> (CELLO Collab.)
DERRICK	86	PL 166B 463	M. Derrick <i>et al.</i> (HRS Collab.)
Also		PR D34 3286	M. Derrick <i>et al.</i> (HRS Collab.)
JODIDIO	86	PR D34 1967	A. Jodidio <i>et al.</i> (LBL, NWES, TRIU)
Also	00	PR D37 237 (erratum)	A. Jodidio <i>et al.</i> (LBL, NWES, TRIU)
MOHAPATRA	86	PR D34 909	R.N. Mohapatra (UMD)
ADEVA	85	PL 152B 439	B. Adeva <i>et al.</i> (Mark-J Collab.)
BERGER	85B	ZPHY C27 341	C. Berger <i>et al.</i> (PLUTO Collab.)
STOKER	85	PRL 54 1887	D.P. Stoker <i>et al.</i> (LBL, NWES, TRIU)
ADEVA	84	PRL 53 134	B. Adeva <i>et al.</i> (Mark-J Collab.)
BEHREND	84C	PL 140B 130	H.J. Behrend <i>et al.</i> (CELLO Collab.)
BERGSMA	83	PL 122B 465	F. Bergsma <i>et al.</i> (CHARM Collab.)
CARR	83	PRL 51 627	J. Carr <i>et al.</i> (LBL, NWES, TRIU)
BEALL	82	PRL 48 848	G. Beall, M. Bander, A. Soni (UCI, UCLA)
SHANKER	82	NP B204 375	O. Shanker (TRIU)
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